



Originally published as:

Dobslaw, H., Bergmann-Wolf, I., Forootan, E., Dahle, C., Mayer-Gürr, T., Kusche, J., Flechtner, F. (2016): Modeling of present-day atmosphere and ocean non-tidal de-aliasing errors for future gravity mission simulations. - *Journal of Geodesy*, 90, 5, pp. 423–436.

DOI: <http://doi.org/10.1007/s00190-015-0884-3>

1 **Modeling of Present-Day Atmosphere and Ocean**
2 **Non-Tidal De-Aliasing Errors for Future Gravity**
3 **Mission Simulations**

4 **Henryk Dobsław · Inga Bergmann-Wolf ·**
5 **Ehsan Forootan · Christoph Dahle ·**
6 **Torsten Mayer-Gürr · Jürgen Kusche ·**
7 **Frank Flechtner**

8
9 Received: date / Accepted: date

10 **Abstract** A realistically perturbed synthetic de-aliasing model consistent with
11 the updated Earth System Model of the European Space Agency is now available
12 over the period 1995 – 2006. The data-set contains realizations of (i) errors at large
13 spatial scales assessed individually for periods between 10 – 30, 3 – 10, and 1 – 3

H. Dobsław
Deutsches GeoForschungsZentrum GFZ, Department 1: Geodesy and Remote Sensing,
Telegrafenberg, 14473 Potsdam, Germany. Tel.: +49-331-288 1974, E-mail: dobslaw@gfz-
potsdam.de

I. Bergmann-Wolf
Deutsches GeoForschungsZentrum GFZ Department 1: Geodesy and Remote Sensing, Tele-
grafenberg, 14473 Potsdam, Germany. Tel.: +49-331-288 1770 E-mail: ingab@gfz-potsdam.de

E. Forootan
Bonn University, Institute of Geodesy and Geoinformation, Nussallee 17, 53115 Bonn, Ger-
many. Tel.: +49-228-73-6423, E-mail: forootan@geod.uni-bonn.de

Ch. Dahle
Deutsches GeoForschungsZentrum GFZ, Department 1: Geodesy and Remote Sensing, c/o
DLR Oberpfaffenhofen, 82234 Wessling, Germany. Tel.: +49-8153-28-1589, E-mail: dahle@gfz-
potsdam.de

T. Mayer-Gürr
Graz University of Technology, Institute for Geodesy, Steyrergasse 30/III, 8010 Graz, Austria.
Tel.: +43-316-873-6359, E-mail: mayer-guerr@tugraz.at

J. Kusche
Bonn University, Institute of Geodesy and Geoinformation, Nussallee 17, 53115 Bonn, Ger-
many. Tel.: +49-228-73-2629, E-mail: jkusche@geod.uni-bonn.de

F. Flechtner
Deutsches GeoForschungsZentrum GFZ, Department 1: Geodesy and Remote Sensing,
c/o DLR Oberpfaffenhofen, 82234 Wessling, Germany. Tel.: +49-8153-28-1297, E-mail:
flechtne@gfz-potsdam.de

14 days, the S1 atmospheric tide, and sub-diurnal periods; (ii) errors at small spatial
15 scales typically not covered by global models of atmosphere and ocean variability;
16 and (iii) errors due to physical processes not represented in currently available
17 de-aliasing products. The model is provided in two separate sets of Stokes coef-
18 ficients to allow for a flexible re-scaling of the overall error level to account for
19 potential future improvements in atmosphere and ocean mass variability models.
20 Error magnitudes for the different frequency bands are derived from a small en-
21 semble of four atmospheric and oceanic models. For the largest spatial scales up
22 to $d/o = 40$ and periods longer than 24 h, those error estimates are approximately
23 confirmed from a variance component estimation based on GRACE daily normal
24 equations. Future mission performance simulations based on the updated Earth
25 System Model and the realistically perturbed de-aliasing model indicate that for
26 GRACE-type missions only moderate reductions of de-aliasing errors can be ex-
27 pected from a second satellite pair in a shifted polar orbit. Substantially more
28 accurate global gravity fields are obtained when a second pair of satellites in an
29 moderately inclined orbit is added, which largely stabilizes the global gravity field
30 solutions due to its rotated sampling sensitivity.

31 **Keywords** Time-Variable Gravity Field · Future Missions · Atmosphere and
32 Ocean De-Aliasing Errors

33 1 Introduction

34 Observations from the twin-satellite mission GRACE (*Tapley et al.*, 2004) allow
35 one to solve for a unique global gravity field model from observations accumulated
36 over a time period of typically 30 days. After 13 years of operation, GRACE-based
37 time-series representing month-to-month changes in the external gravitational field
38 of the Earth offer the opportunity to infer variations of the water masses in the soil
39 and sub-surface aquifers (*Famiglietti and Rodell*, 2013), and the continental ice-
40 sheets (*Sasgen et al.*, 2013) down to spatial scales of a few 100 km. Such changes
41 in the amount of water stored on the continents are characterized by rather low
42 frequencies that are well resolved by a monthly sampling. In contrast to this,
43 mass redistributions in the atmosphere, as, e.g., associated with the advection of
44 cyclones, or in the oceans in response to time-variable surface winds, are dominated
45 by much shorter periods. Since those mass signals are of similar or even greater
46 magnitude than terrestrial water storage changes, their effects need to be properly
47 treated in the gravity field estimation process to avoid temporal aliasing of high-
48 frequency variability into the monthly mean gravity fields (*Thompson*, 2004; *Han*
49 *et al.*, 2004).

50 The Atmosphere and Ocean De-Aliasing Product (AOD1B; *Flechtner and Dob-*
51 *slaw*, 2013) is routinely calculated by the GRACE Science Data System to provide
52 a time-variable background model for the removal of non-tidal high-frequency mass
53 signals from the sensor data. The most recent release 05 of AOD1B is based on
54 6 hourly sampled global atmospheric analysis fields from the European Centre
55 for Medium Range Weather Forecasts (ECMWF) and an unconstrained simula-
56 tion with the global ocean circulation model OMCT (*Dobslaw et al.*, 2013). In
57 addition to this widely applied standard model, alternative atmosphere and ocean
58 de-aliasing (AOD) coefficients are calculated by, e. g., Groupe de Recherche de

59 Geodesie Spatiale (GRGS; *Boy, 2005*), TU Munich (*Zenner et al., 2010*), and Uni-
60 versity of Bonn (*Forootan et al., 2013*). Imperfections in any of those de-aliasing
61 models, however, will lead directly to residual atmospheric and oceanic signals in
62 the instrument data that inevitably alias into the monthly mean gravity field so-
63 lution. Such artifacts are frequently assumed to be one of the major reasons that
64 the GRACE baseline accuracy targeted prior to the launch of the mission (*Kim,*
65 2000) is still not reached by the most recent data releases (*Sakumura et al., 2014*).

66 While the GRACE Follow-On mission is progressing well towards its antici-
67 pated launch date in August 2017 (*Flechtner et al., 2014*), mission performance
68 simulation studies are carried out already now by several research groups in order
69 to eventually propose concepts for future gravity missions to the space agencies
70 (e.g., *Panet et al., 2012; Gruber et al., 2014*). Such simulations evaluate costs and
71 benefits of higher sensor accuracies, orbits at lower altitudes that even might re-
72 quire active drag compensation, or constellations of multiple satellites connected
73 by inter-satellite ranging links. In those studies, de-aliasing errors are typically ei-
74 ther approximated from the difference of two arbitrarily selected numerical models
75 (*Wiese et al., 2009, 2012; Loomis et al., 2011*) or by assuming that the errors are
76 linearly dependent on the signal magnitude (*Visser, 2010; Elsaka et al., 2014*).
77 Both approaches lead to widely different magnitudes of de-aliasing errors, and
78 consequently widely different conclusions about limiting factors on the overall ac-
79 curacy of a future mission candidate concept.

80 We therefore attempt in this study to consolidate the present-day level of non-
81 tidal atmosphere and ocean de-aliasing errors and prepare a realization of such
82 errors that is readily applicable in future mission simulation studies together with
83 the updated Earth System Model (ESM; *Dobslaw et al., 2015*) of the European
84 Space Agency (ESA). In the following, we discuss a strategy to approximate er-
85 rors at large spatial scales that are typically well represented by global numerical
86 atmosphere and ocean models (Sect. 2); and at small spatial scales that are only
87 resolved by high-resolution regional models (Sect. 3). We further assess effects of
88 physical processes that are not covered by a typical AOD model (Sect. 4), and
89 describe how all those errors are made publicly available as part of the updated
90 ESM (Sect. 5). Subsequently, a variance component estimation with present-day
91 daily GRACE normal equations is performed to demonstrate the consistency of
92 our error estimates with observational evidence (Sect. 6), and full-scale simulation
93 experiments for both a GRACE setting and two double-pair configurations are
94 carried out in order to assess the impact of those atmosphere and ocean non-tidal
95 background model errors on the performance of future gravity missions (Sect. 7),
96 before a brief summary is provided in Sect. 8.

97 2 AOD Errors at Large Spatial Scales

98 The mass variability in atmosphere and ocean represented by any AOD model cur-
99 rently available is inevitably contaminated by errors of the underlying geophysical
100 models. Those errors are typically correlated in time and space, and are distributed
101 all over the globe as a complex function of latitude, orography, and other factors.
102 Since those dependencies are difficult to describe analytically, hydrodynamic model
103 errors are typically approximated from an ensemble of model realizations that are
104 implicitly assumed to be equally probable. We follow this approach by analyzing

105 pressure anomaly data from a small ensemble of four up-to-date atmospheric and
106 oceanic model data-sets available to us for the year 2006.

107 It should be noted that errors at temporal scales larger than 30 days that
108 also include, e.g., apparent jumps in the atmospheric pressure due to changes in
109 the orography of the operational ECMWF model (e.g., *Duan et al.*, 2012) will
110 not be considered. Errors at those temporal scales map directly into the monthly
111 gravity fields as systematic biases, but do not contribute to the variability within
112 the month and therefore do not actually lead to de-aliasing errors. Such monthly
113 biases might be corrected a posteriori by means of subtracting an alternative
114 reference model, whereas de-aliasing errors cannot be removed from the monthly-
115 mean gravity fields with post-processing methods (*Forootan et al.*, 2014).

116 The updated ESM intended to serve as a source model in future mission sim-
117 ulations covers the period 1995 – 2006. Yet, it would be convenient to have a
118 perturbed de-aliasing model available for the same period. Since the original ESM
119 of *Gruber et al.* (2011) covers that time-frame as well, we use differences between
120 those two data-sets as a basis for a series of error realizations. This choice is more-
121 over justified by the fact that the original ESM is partly based on ECMWF’s
122 previous re-analysis ERA-40 (*Uppala et al.*, 2005) and an earlier OMCT experi-
123 ment (*Dobslaw and Thomas*, 2007), which have been both part of the now out-
124 dated GRACE AOD1B release 04. Thus, the difference between the two ESMs
125 with its very complex spatial and temporal correlations is representative for errors
126 contained in AOD1B only three years ago.

127 To account for likely model improvements after the release 04 of AOD1B and/or
128 common errors in the original and the updated ESM, we will subsequently scale
129 those differences with constant coefficients for land and ocean regions in order
130 to match the spread of the ensemble of current state-of-the-art model data-sets
131 for the year 2006 at large spatial scales up to approximately spherical harmonic
132 degree and order (d/o) 60. Since we expect a frequency dependence of the scaling
133 coefficients, we will dissect all model differences from the ensemble and also the
134 two ESMs by a series of 3rd order Butterworth filters into five different bands
135 covering periods of 10 – 30 days, 3 – 10 days, 1 – 3 days, the S1 atmospheric tide,
136 and the sub-diurnal periods. In the following, we describe our procedure for the 10
137 – 30 days band explicitly. Please note that the ensemble spread related to signals
138 at intermediate spatial scales over the oceans will be covered in particular by the
139 third error component due to omitted physical processes, which will be introduced
140 later in Sect. 4.

141 One member of the small ensemble of atmospheric model data-sets is the cur-
142 rent global re-analysis of the ECMWF, ERA-Interim (*Dee et al.*, 2011) that is
143 also used for the updated ESM. In addition, we use the Climate Forecast System
144 Reanalysis (CFSR; *Saha et al.*, 2010) prepared by the National Center for Environ-
145 mental Prediction, and NASA’s Modern Era Retrospective Analysis for Research
146 and Applications (MERRA; *Rienecker et al.*, 2011). The fourth data source is the
147 operational ECMWF analysis data-set that is also used for AOD1B. In order to
148 assess the current level of uncertainties in global models of atmospheric mass vari-
149 ability over the continents, we calculate pairwise rms differences of atmospheric
150 surface pressure that are bandpass filtered to contain only variability with periods
151 of 10 – 30 days over the year 2006 (Fig. 1). For closely related models, as, e.g.,
152 operational and re-analysis data from ECMWF, we note globally fairly homoge-
153 neous residuals of about 0.2 hPa only, which is equivalent to a 2 mm change in

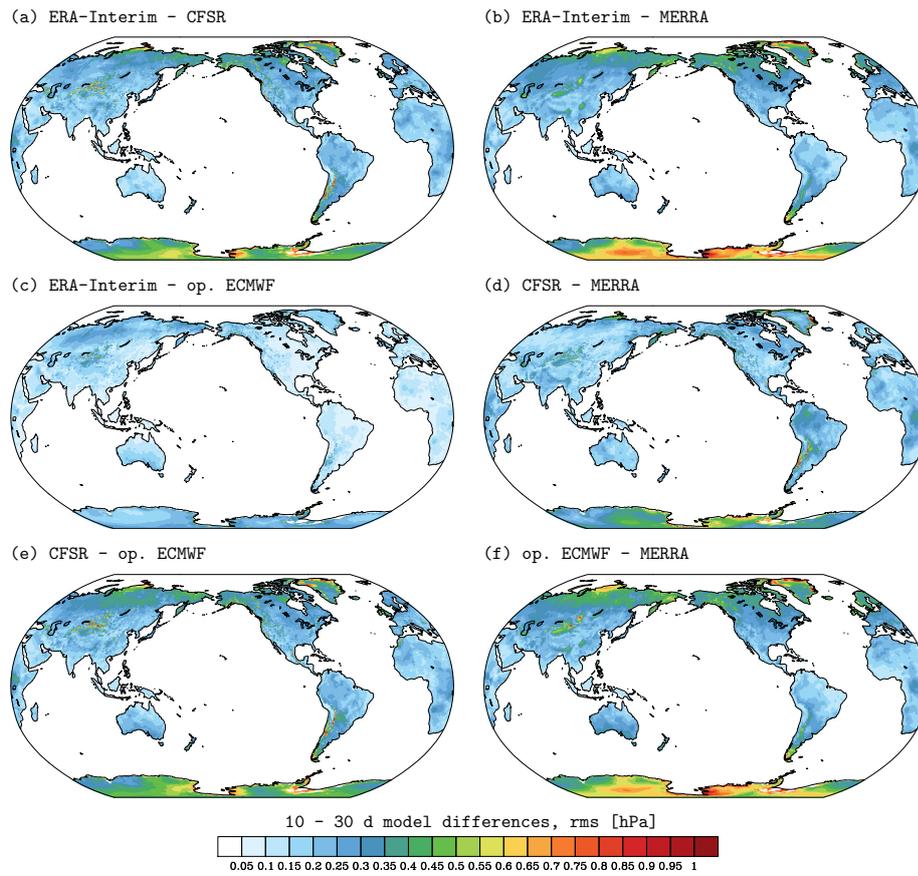


Fig. 1 Rms differences of band-pass filtered atmospheric mass variability from four different global atmospheric surface pressure data-sets at periods of 10 – 30 days for the year 2006: ERA-Interim re-analysis, MERRA re-analysis, CFSR re-analysis, and ECMWF operational analysis

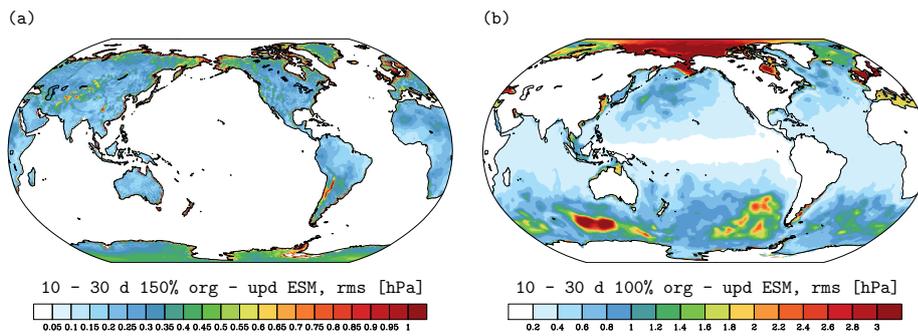


Fig. 2 Rms of re-scaled atmospheric surface pressure (left) and ocean bottom pressure (right) errors with large spatial scales at periods of 10 – 30 days as averaged over the period 1995 – 2006

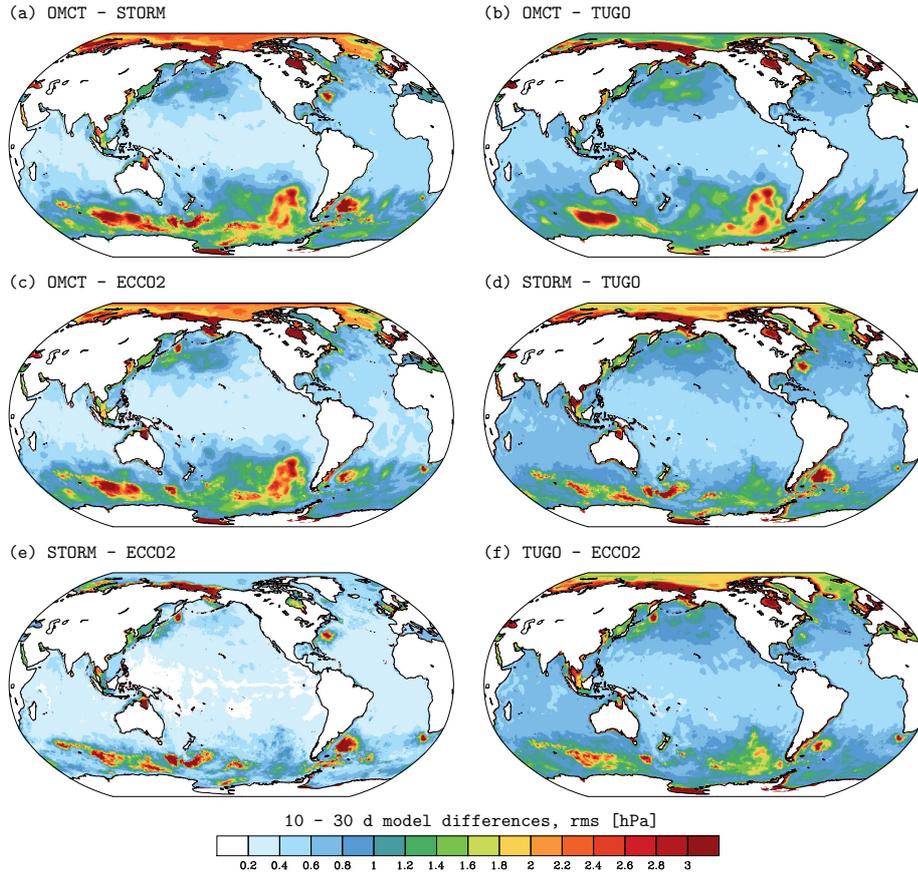


Fig. 3 Rms differences of band-pass filtered ocean bottom pressure variability from four different global ocean models at periods of 10 – 30 days for the year 2006: OMCT as used for AOD1B RL05, TUGO, ECCO2, and MPIOM STORM

154 water height. The spread is substantially larger when any of the ECMWF models
 155 is compared with CFSR or MERRA, culminating in rms differences of up to 0.7
 156 hPa in Antarctica. Based on the pairwise model comparison, we conclude that
 157 the spread between the two ESMS underestimates the uncertainty. We therefore
 158 up-scale the atmospheric component by 150% for periods of 10 – 30 days (Fig.
 159 2a).

160 The ensemble of ocean model data-sets consists of a simulation from OMCT
 161 forced with ERA-Interim atmospheric data, the STORM experiment performed
 162 with the Max-Planck-Institute for Meteorology Ocean Model (MPIOM) that uses
 163 NCEP2 atmospheric data (*von Storch et al., 2012*), and a recent run with the french
 164 TUGO model (*le Bars et al., 2010*) under ERA-Interim forcing that is contained
 165 in the alternative AOD coefficients provided by GRGS. Finally, we utilize the
 166 ECCO2 ocean synthesis performed at JPL (*cube92; Menemenlis and Campin,*
 167 *2008*) that assimilates a vast range of different oceanographic observations. We
 168 note rather small differences between STORM and ECCO2 that both lack forcing

Table 1 Scaling Coefficients for continental and oceanic regions applied to the model differences of the original and the updated ESM for the realization of de-aliasing errors at large spatial scales for five different frequency bands

	Continents	Oceans
10 – 30 days	150%	100%
3 – 10 days	250%	100%
1 – 3 days	350%	150%
S ₁ (p)	250%	150%
sub-diurnal	180%	150%

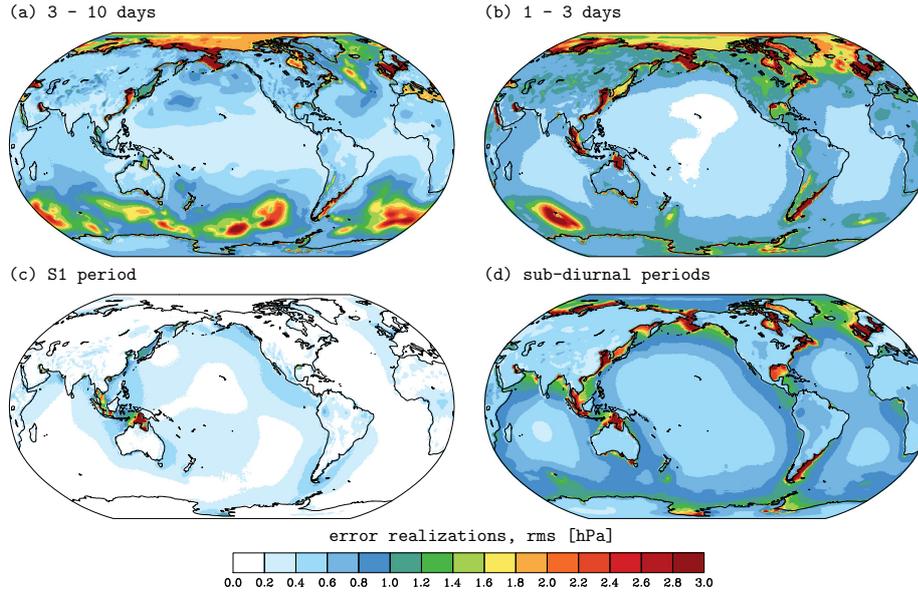


Fig. 4 Rms of the sum of the re-scaled atmospheric surface and ocean bottom pressure errors at large spatial scales with periods of (a) 3 – 10 days, (b) 1 – 3 days, (c) the S1 atmospheric tide, and (d) sub-diurnal periods over the period 1995 – 2006

169 from atmospheric pressure variability and thereby miss a small but non-negligible
 170 part of the energy input from the atmosphere into the oceans (Fig. 3). This is in
 171 particular relevant for the Arctic Ocean, where the inverse barometric adjustment
 172 to atmospheric pressure changes is hampered by the narrow straits connecting
 173 the basin to the Pacific and the North Atlantic. Residuals of all pairwise model
 174 differences show a clear zonal structure, where larger errors are found in regions
 175 with larger signal variability. Based on those comparisons, we conclude that the
 176 uncertainty associated with ocean bottom pressure variability for periods of 10 –
 177 30 days is well represented by the spread between the two ESMs (Fig. 2b).

178 Similar assessments are performed for the periods of 3 – 10 days, 1 – 3 days,
 179 the S1 atmospheric tide, as well as the sub-diurnal periods; details are given in
 180 *Bergmann-Wolf et al. (2015)*. The estimated scaling coefficients (Tab. 1) indicate
 181 that the original and the updated ESM tend to converge for higher frequencies,
 182 thereby requiring stronger up-scaling at shorter periods. This is consistent with

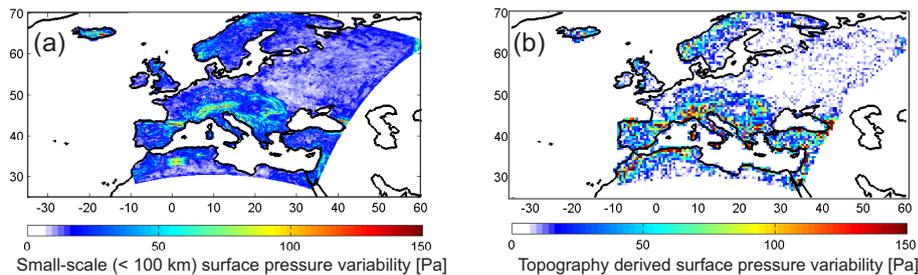


Fig. 5 Rms variability of pressure signals at small spatial scales from the high-resolution regional model COSMO-EU that are typically missed out by global atmospheric re-analysis data-sets (a), and a reconstruction of small-scale atmospheric error variability from a regression model based on the local orographic roughness (b)

183 our findings from the pairwise model comparison, since similar modelling assump-
 184 tions and common errors included in two closely related model realisations tend
 185 to have an increasing influence at higher frequencies. The spatial variability of
 186 errors in the different frequency bands are summarized in Fig. 4. For the conti-
 187 nents, error patterns of atmospheric surface pressure are fairly homogeneous over
 188 all latitudes at all periods considered. For the oceans instead, we find distinctly
 189 latitude-dependent errors at in particular lower frequencies with almost no errors
 190 in the tropics and highest values at extra-tropical latitudes ($40^\circ - 60^\circ$ N and S)
 191 characterized by strong westerly winds. In addition, we note that a number of sub-
 192 surface oceanic basins have distinctly different resonance frequencies that lead to
 193 peak error levels in particular sectors of the Southern Ocean for the various periods
 194 considered, as, e.g., the Bellingshausen Basin (60° S, 60° W) or the Australian
 195 Antarctic Basin (50° S, 100° E). For more rapid processes at diurnal periods and
 196 beyond, errors are rather dominated by basin-wide modes, resembling the typical
 197 tidal resonances at those frequencies.

198 3 AOD Errors at Small Spatial Scales

199 Global atmospheric re-analyses often used for the derivation of AOD coefficients
 200 are confined to the representation of atmospheric dynamics, land-atmosphere in-
 201 teractions, and orography at large spatial scales only, since the discretization of
 202 the model equations in time and space requires the parametrization of local pro-
 203 cesses from a certain threshold on. Both ESM data-sets are expanded in Stokes
 204 coefficients up to $d/o = 180$, which roughly corresponds to a spatial resolution of
 205 1° , but there is certainly variability at spatial scales below that threshold in both
 206 atmosphere and oceans.

207 To assess signals at small spatial scales not covered explicitly by the updated
 208 ESM, we utilize data from a recent experiment with the non-hydrostatic high-
 209 resolution numerical weather prediction model COSMO-EU (*Baldauf et al., 2011*).
 210 The model is discretized over Europe with a horizontal resolution of approximately
 211 7 km. Available to us is an operational analysis run for the year 2006 with daily
 212 sampling. To isolate small-scale variability at high frequencies, a high-pass fil-
 213 ter with 30 days cutoff period is applied to the surface pressure anomalies. Sub-

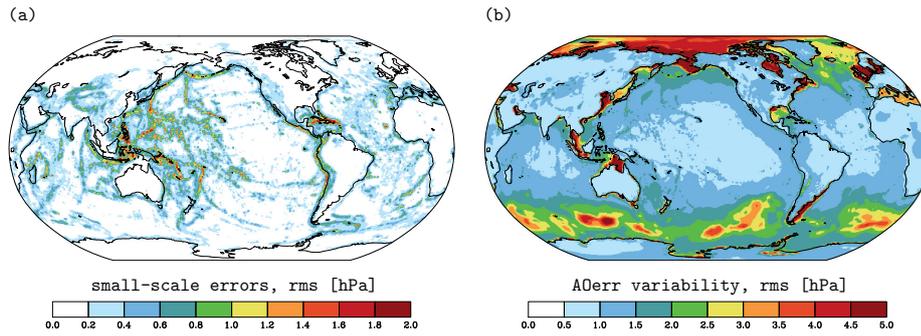


Fig. 6 Rms errors at small spatial scales obtained from a regression model based on COSMO-EU (a), and rms variability of the component AOerr of the realistically perturbed de-aliasing model that summarizes errors at both large and small spatial scales (b)

214 sequentially, 1° spatial averages are removed from the high-pass filtered pressure
 215 anomalies to retain only signals at spatial scales smaller than approximately 100
 216 km. Rms values of those residuals are as large as 1.5 hPa in mountainous regions
 217 as, e.g., the Alps and the Pyrenees (Fig. 5a). Those signals are highly localized
 218 and in particular confined to areas with very rough orography. Most other regions
 219 instead, show residual pressure signals that are well below 0.5 hPa.

220 Since high-resolution model data is only available to us from an atmospheric
 221 model over a limited domain, we extrapolate small-scale errors to the whole globe
 222 by hypothesizing a linear relationship between the roughness of the orography
 223 and residual pressure signals at small spatial scales (see also, e.g., *van Dam*
 224 *et al.*, 2010). To obtain a roughness scale, we smooth the global orography data-set
 225 ETOPO2 that also includes sub-surface ocean bathymetry at its original spatial
 226 resolution of $2'$ by a two-dimensional moving average filter with a window length
 227 of 30 grid cells. The difference between the original topography and the smoothed
 228 version is taken as an orographic roughness indicator for the regression model.
 229 Small-scale pressure variability as obtained from orographic roughness and the
 230 regression model at a 0.5° spatial resolution approximates the signal quite well
 231 over land (Fig. 5b). We now randomly sample normal distributions with the er-
 232 ror variances at small spatial scales derived from the regression model to arrive
 233 at a series of 250 independent daily realizations that are subsequently cycled for
 234 about 18 times to cover the full 12 year-long period of the updated ESM (Fig.
 235 6a). Regions covered by high mountains as the Himalaya or the Andes dominate
 236 the small-scale errors on the continents, whereas deep-sea trenches and sub-ocean
 237 ridges are clearly discernable in the oceanic regions. Errors are spatially highly
 238 variable and approach 2 hPa in isolated places only.

239 In order to validate our extrapolation approach, we show rms variability of
 240 pressure signals at small spatial scales from two moderate resolution global mod-
 241 els – the operational analyses of ECMWF and the MPIOM STORM experiment
 242 (*von Storch et al.*, 2012) for January 2011 and 2006, respectively – that have been
 243 obtained by applying a Gaussian high-pass filter with 100 km half-wavelength to
 244 the daily pressure anomaly grids (Fig. 7). Both data-sets have an approximate spa-
 245 tial resolution of 0.1° which is more than twice the sampling of COSMO-EU. We
 246 note very good correspondence with the orography-based error predictions over the

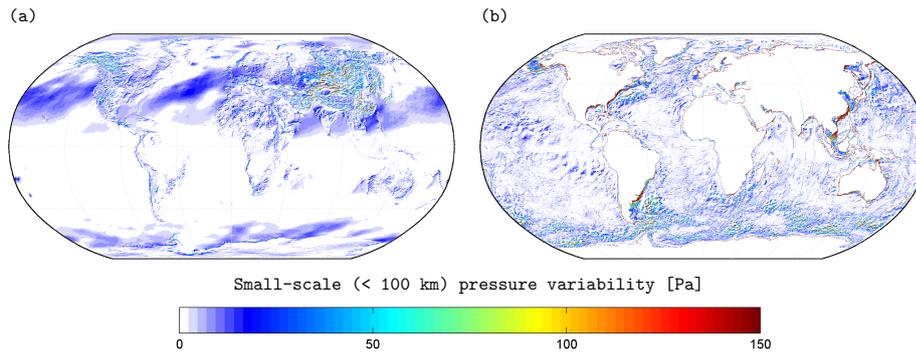


Fig. 7 Rms variability of pressure signals at small spatial scales from the global ECMWF numerical weather prediction model (a), and the global ocean model experiment STORM performed with the MPIOM ocean model (b), both obtained by high-pass filtering the original daily pressure anomalies at 0.1° with a Gaussian smoother of 100 km half-wavelength from 30 days of data during January 2011 and 2006, respectively

247 continents that fully justifies our approach in those regions. For the oceans, how-
 248 ever, pattern deviate substantially: small-scale signals in the STORM experiment
 249 are rather concentrated along the coasts and do show only moderately enhanced
 250 amplitudes at bathymetric structures as deep-sea trenches. Since nevertheless the
 251 signal amplitudes approaching 150 Pa in isolated places only are consistent with
 252 our extrapolation approach, we conclude that the small-scale errors predicted over
 253 the oceans are indeed suitable to initially test the relevance of such small-scale
 254 features for satellite gravimetry. For a more rigorous assessment of the impact of
 255 small-scale errors on future satellite gravity mission concepts, as, e.g., a satellite-
 256 mounted quantum gradiometer, a revision of this error model component based
 257 on high-resolution global atmosphere and ocean model experiments should be at-
 258 tempted.

259 4 Oceanic Signals Omitted in AOD Models

260 The oceanic components of all currently available AOD models rely on either
 261 barotropic (*Carrère and Lyard, 2003*) or baroclinic (*Thomas et al., 2001*) global
 262 ocean circulation models forced with atmospheric data from numerical weather
 263 models that do not assimilate any type of high-frequency observational data. Such
 264 simulations typically exclude meso-scale variability and small-scale eddies, which
 265 are primarily near-surface features but partially have also bottom pressure signa-
 266 tures in particularly energy-rich areas of the world’s ocean (*Kuhlmann et al., 2013*).
 267 In addition, those models typically do not include changes in global barostatic sea-
 268 level (*Chambers, 2004*), which causes a seasonal variation of the globally averaged
 269 ocean bottom pressure of about 1 hPa.

270 To account for those principal deficits of presently available AOD products in
 271 future mission studies, we derive a set of unperturbed AOD coefficients for the
 272 updated ESM. In continental regions it is identical to the atmospheric part of the
 273 updated ESM, but the oceanic component is only based on the OMCT model
 274 run and excludes the small-scale variability from MPIOM STORM (*von Storch*

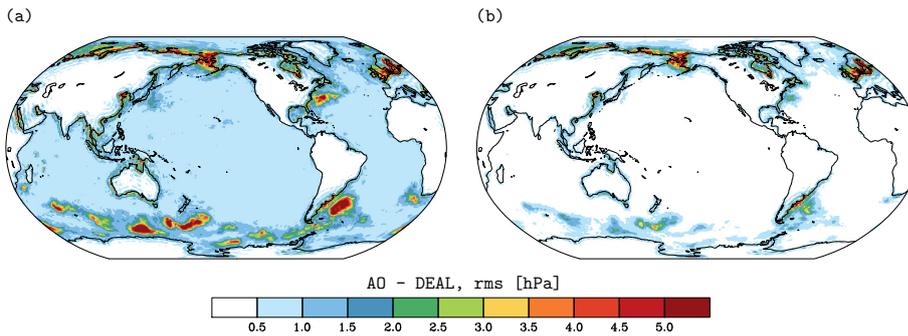


Fig. 8 Rms differences (a) and rms of 30 days high-pass filtered differences (b) between the atmospheric and oceanic (AO) component of the updated ESM and the DEAL component of the realistically perturbed de-aliasing model calculated over the period 1995 - 2006, which reflect the impact of physical processes in the ocean that are currently not represented in release 05 of AOD1B

275 *et al.*, 2012) and the barostatic sea-level variability that have been added into the
 276 updated ESM. Rms differences therefore deviate from zero only over the oceans
 277 (Fig. 8a), where a globally homogeneous signal of about 0.7 hPa rms related to
 278 seasonal changes in total ocean mass dominates. In addition, we find areas with
 279 distinctly large residuals of more than 5 hPa in oceanic regions characterized by
 280 high eddy kinetic energy, and small-scale mass variability all along the coasts.
 281 Signal magnitudes drop substantially, however, when only high-pass filtered vari-
 282 ability at periods shorter than 30 days is considered (Fig. 8b), which is the part
 283 of the signal that actually contributes to the de-aliasing errors. High-pass filtered
 284 signals in the open ocean are well below 1 hPa, thereby indicating that global
 285 ocean simulations that are not eddy-resolving are sufficiently well suited to serve
 286 as de-aliasing model for high-frequency bottom pressure variability.

287 **5 Two New Components of the Updated Earth System Model: AOerr** 288 **and DEAL**

289 To allow for a straightforward implementation of those AOD-related errors in
 290 future satellite mission simulation studies, two series of Stokes coefficients up to
 291 $d/o = 180$ with a time sampling of 6 hours over the full 12-year time-period of the
 292 updated ESM are made publicly available via doi:10.5880/GFZ.1.3.2014.001. First,
 293 the "DEAL" coefficients contain the unperturbed de-aliasing model introduced
 294 in the previous section. It might be used as a background model for satellite
 295 simulation studies in case that – apart from the principal deficits indicated above
 296 – perfect model-based predictions of atmospheric and oceanic mass variability
 297 should be assumed. If even the removal of those principal model deficits appears
 298 plausible until the anticipated launch date of a future mission, a replacement of
 299 DEAL with the sum of the atmospheric and oceanic components of the updated
 300 ESM might be considered.

301 Secondly, the "AOerr" coefficients consist of the sum of the error series at
 302 small spatial scales as well as at large spatial scales at all five frequency bands

discussed above. AOerr needs to be added to DEAL to arrive at a reasonably perturbed AOD model. Time series of AOerr at individual geographic positions have zero mean over the 12-year period of data provided and an approximately stationary variance when individual years are considered (*Bergmann-Wolf et al., 2015*). The overall variability of AOerr (Fig. 6b) is dominated by large-scale errors over the oceans with particularly high error magnitudes in the Arctic Ocean and in the vicinity of the Antarctic Circumpolar Current. Continental regions show a comparatively homogeneous error level, although slightly higher rms values are found in glaciated regions at polar latitudes. Errors at small spatial scales instead tend to have a minor influence only. For satellite simulation experiments aiming at satellite scenarios in a rather distant future, a reduction of those AOerr coefficients by means of a global factor might be performed in order to account for anticipated future improvements in the accuracy of those models until launch.

6 Variance Component Estimation from GRACE

In the following, we are going to contrast the newly derived series of error realizations against GRACE-based estimates of high-frequency mass variability. The recently published ITSG-Grace2014 release (*Mayer-Gürr et al., 2014*) contains in addition to long-term and monthly-mean global gravity fields also a series of daily solutions based on a Kalman Smoother approach (*Kurtenbach et al., 2009, 2012*). The large-scale signals at periods down to a few days contained in the predecessor of this daily series have been found to be geophysically plausible both on the continents (*Zenner et al., 2014*) and over the oceans (*Bergmann and Dobslaw, 2012*) from comparisons with independent in situ observations and numerical model estimates. From the daily normal equations of ITSG-Grace2014 expanded into spherical harmonics up to $d/o = 40$ we remove all signals at periods longer than 30 days by estimating and subtracting the long-term mean and drift as well as subsequently also the monthly-mean components. Since all standard time-variable background models for solid Earth tides, ocean tides, and also non-tidal variability from AOD1B release 05 have been applied, we expect that the remaining residuals only contain GRACE system noise, sub-monthly variations in terrestrially stored water, episodic seismic events, and errors in background models applied.

To separate different contributors to the GRACE residuals we describe their stochastic behaviour by a number of variance-covariance matrices, implicitly assuming that each of the components is stationary in time. Those covariance matrices for the remaining atmospheric and oceanic signals are calculated from the 12 year-long series of non-tidal atmosphere and ocean error realizations. We dissect the AOerr series into nine different regions of the world: the Arctic Ocean, the oceans at the Northern Hemisphere between $23^{\circ}\text{N} - 67^{\circ}\text{N}$, the tropical oceans, and the southern oceans at latitudes below 23°S . Continental areas are divided into Antarctica, Greenland, the northern hemispheric continents at latitudes above 23°N except for Greenland, the tropical continents, and the southern hemispheric continents at $23^{\circ}\text{S} - 60^{\circ}\text{S}$. To additionally account for sub-monthly water storage changes on the continents that might be eventually correlated to atmospheric errors, we also derive empirical covariances out of the continental components of the updated ESM containing ice mass and terrestrial water storage changes that are

Table 2 Area-weighted average rms at large spatial scales up to $d/o = 40$ and for periods of 1 – 30 days in nine different regions of the globe as contained in the updated ESM (a priori) and derived from two variance component estimation experiments with consistently pre-processed daily normal equations from ITSG-GRACE2014 (a posteriori)

		ESM a priori [hPa]	ITSG-GRACE2014 a posteriori [hPa]	
			individual	summarized
Antarctica	surface pressure	0.81	1.16	-
	ice mass change	0.13	0.52	-
	sum	0.82	-	1.26
Greenland	surface pressure	0.66	1.27	-
	ice mass change	0.21	0.50	-
	sum	0.69	-	1.36
northern continents (23°N – 80°N)	surface pressure	0.61	1.15	-
	terrestrial water	0.35	0.64	-
	sum	0.70	-	1.26
southern continents (23°S – 60°S)	surface pressure	0.58	1.06	-
	terrestrial water	0.48	0.68	-
	sum	0.75	-	1.23
tropical continents (23°S – 23°N)	surface pressure	0.50	0.63	-
	terrestrial water	0.43	1.04	-
	sum	0.66	-	1.27
tropical oceans	bottom pressure	0.65	0.76	0.76
Arctic Ocean	bottom pressure	2.83	3.01	3.03
northern oceans	bottom pressure	1.41	1.21	1.21
southern oceans	bottom pressure	1.68	1.27	1.28

348 filtered for periods of 1 – 30 days as well, leading to in total 14 covariance matrices
349 solely for the non-tidal geophysical signals.

350 The sum of all covariance matrices is expected to describe the stochastic char-
351 acteristics of the daily GRACE solutions. The method of variance component
352 estimation (VCE; Koch and Kusche, 2002) allows for estimating pre-factors for
353 each covariance matrix in an iterative procedure. Those variance factors obtained
354 from the VCE are subsequently multiplied by the prior variances and transformed
355 into area-weighted rms (wrms) values. By comparing prior wrms estimates from
356 the updated ESM with a posteriori results from the VCE (Tab. 2), we note that the
357 level of errors derived in this study is generally confirmed by ITSG-Grace2014. For
358 the continents, GRACE-based wrms are typically 0.4 hPa larger than predicted
359 by AOerr, which is very well acceptable since additional non-atmospheric error
360 contributors to the GRACE normal equations at this level cannot be entirely ex-
361 cluded. A second VCE experiment utilizing only 9 instead of 14 empirical error
362 covariance matrices by summing atmospheric errors and short-term water storage
363 changes beforehand (Tab. 2; third column) indicates weak correlations between
364 atmospheric mass variability errors and sub-monthly water storage changes, but
365 generally confirms that ITSG-Grace2014 contains slightly higher variability over
366 the continents than predicted by AOerr and the sub-monthly water storage vari-
367 ability contained in the updated ESM. Over the oceans, the substantially higher
368 error levels contained in AOerr are also confirmed by GRACE. The exceptionally
369 strong wrms of about 3 hPa in the Arctic Ocean in particular is seen by GRACE as
370 well, whereas the errors in the extra-tropical oceans at both hemispheres are found

371 to be modelled slightly too large. We nevertheless note that differences between
 372 the prior estimates and the GRACE-based results never exceed 0.5 hPa in wrms
 373 both on continents and in the oceans, letting us conclude that the error realiza-
 374 tion series provided in AOerr realistically represents the current level of AOD1B
 375 accuracy for spatial scales up to $d/o = 40$ and periods of 1 – 30 days.

376 7 Simulated Aliasing Errors for Future Missions

377 We finally attempt to estimate the spatial characteristics of temporal aliasing
 378 errors for a set of different GRACE-type satellite constellations to provide evidence
 379 for the relative importance of such type of errors for a small range of plausible
 380 future mission scenarios. We use the simulation setting described by *Flechtner*
 381 *et al.* (2015), who recently performed an in-depth assessment of the upcoming
 382 GRACE-FO mission and its expected contributions to Earth science applications.
 383 We perform four twin experiments ranging each over the three months April – June
 384 2005 that all assume satellites in pearl-string formation connected by microwave
 385 intersatellite links at altitudes of 480 km. The scenarios include (i) a single polar
 386 mission, (ii) a single polar mission with an orbital plane shifted by 90° in right
 387 ascension, (iii) a combination of two polar pairs separated by 90° in right ascension,
 388 as well as (iv) one polar pair and a second pair in an 70° inclined orbit, which is
 389 typically referred to as a Bender-type constellation.

390 For any twin experiment, two monthly global gravity fields are retrieved up
 391 to $d/o = 100$ that differ only in the applied a priori background model: by ap-
 392 plying the atmospheric and oceanic (AO) component of the updated ESM, we
 393 perfectly remove all atmosphere and ocean mass variability signals, whereas the
 394 use of the realistically perturbed de-aliasing model AOerr+DEAL is expected to
 395 cause residual variability. Differences between the two retrievals are interpreted as
 396 de-aliasing errors. Note that further processes leading to temporal aliasing – as,
 397 e.g., ocean tides and sub-monthly water storage variations – included in *Flechtner*
 398 *et al.* (2015) are treated identically in both retrievals and are therefore not visible
 399 in the results discussed below.

400 Differences between finally retrieved gravity fields for May 2005 reveal the spa-
 401 tial characteristics of the de-aliasing errors introduced by imperfect time-variable
 402 background models (Fig. 9). For the unfiltered geoid heights of both single-pair
 403 polar mission experiments, we note substantial de-aliasing errors at moderate lat-
 404 itudes – where much of the signal variability in atmosphere and oceans resides
 405 and thus the errors are larger – and also at tropical latitudes, where the spa-
 406 tial sampling of a polar-orbiting mission is poor. For a combination of two polar
 407 pairs, we see a moderate reduction of de-aliasing errors on the finally retrieved
 408 gravity field by 14% that are distributed quite homogeneously over all latitudes.
 409 For the Bender-type constellation, we find a substantial reduction by more than
 410 69%, indicating that the rotated error anisotropy due to the modified azimuth of
 411 the inter-satellite link of an inclined pair stabilizes the monthly gravity field solu-
 412 tion substantially. Nevertheless, for the Bender-type constellation, the reduction
 413 of aliasing artifacts is not homogeneous over the globe: de-aliasing errors in regions
 414 not sampled by the second pair are as high as for a single pair mission in a polar
 415 orbit.

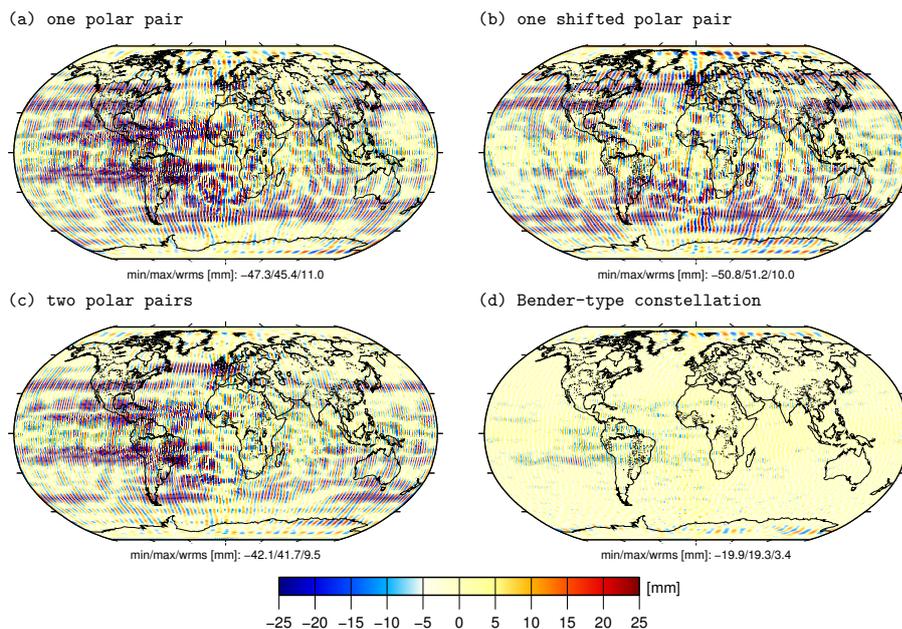


Fig. 9 Geoid height errors introduced into simulated gravity field retrievals from imperfect background models of non-tidal atmosphere and ocean variability for May 2005: (a) polar pair, (b) polar pair shifted by 90° in right ascension, (c) two polar pairs, and (d) a polar pair and a 70° inclined pair

416 Similar levels of reduction are obtained for April and June 2005 with 38%
 417 and 34%, respectively, for two polar pairs, and 62% and 77% for the Bender-type
 418 constellation (not shown). In terms of degree amplitudes (Fig. 10), we note that
 419 AOD1B errors start dominating the solutions at around $d/o = 30$ on a global
 420 average. For two polar pairs, improvements are visible starting at $d/o = 15$, and
 421 AOD1B errors only dominate from $d/o = 40$ onwards. Best results are expected
 422 from a Bender-type constellation with a second pair in a moderately inclined orbit,
 423 that extends this threshold even further to $d/o = 45$ and reduces the impact of
 424 AOD1B errors on the monthly solutions even on higher degrees and orders by a
 425 substantial amount.

426 8 Summary

427 A realistically perturbed synthetic de-aliasing model consistent with the updated
 428 Earth System Model is now available over the period 1995 – 2006. The model
 429 contains (i) errors at large spatial scales that were assessed individually for five
 430 different frequency bands, (ii) errors at small spatial scales typically not treated
 431 by global atmospheric and oceanic models, and (iii) errors due to mass variability
 432 in the oceans that is currently not included in any of the available AOD products.
 433 The realistically perturbed de-aliasing model is publicly available in two sets of
 434 spherical harmonic coefficients: the AOerr coefficients represent a summarized re-

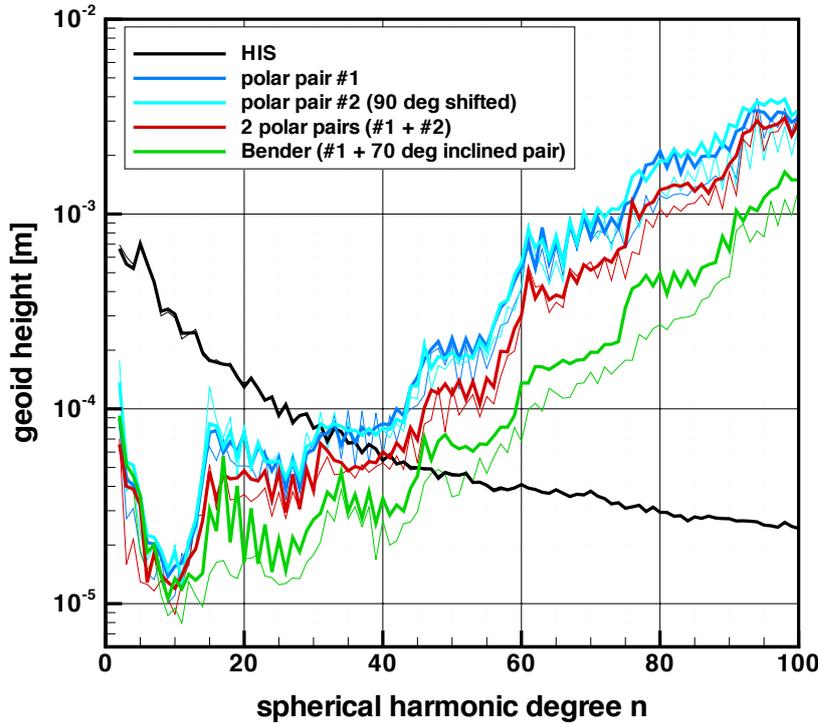


Fig. 10 Geoid height error degree amplitudes introduced into the simulated gravity field retrievals from imperfect background models of non-tidal atmosphere and ocean variability for May 2005 (thin lines) and averaged over three simulated months covering April – June 2005 (thick lines)

435 alisation of errors at large and small spatial scales, whereas the DEAL coefficients
 436 differ from the AO component of the updated ESM only by absent bottom pressure
 437 anomalies related to barostatic sea-level changes and meso-scale dynamics.
 438 By providing those two separate components, we offer users some flexibility to
 439 account for potential improvements in AOD model quality until the anticipated
 440 launch date of mission – be it by a reduction of the overall noise level or by the
 441 additional inclusion of further ocean processes into AOD products.

442 Magnitudes of the errors contained in the perturbed de-aliasing model were de-
 443 rived from analyzing a small ensemble of four atmospheric and oceanic data-sets.
 444 For large spatial scales up to $d/o = 40$ and periods longer than 24 h, those error
 445 levels were generally confirmed from a variance component estimation based on
 446 ITSG-Grace2014 daily normal equations. The perturbed error model is thus be-
 447 lieved to be representative for the present-day errors contained in AOD1B RL05.
 448 Initial full-scale future satellite mission performance simulations with this per-
 449 turbed de-aliasing model indicated that for GRACE-like satellites only moderate
 450 reductions of de-aliasing errors can be expected from two satellite pairs in polar

451 orbits. Substantially larger improvements were found for a so-called Bender con-
452 stellation, where a polar pair is augmented by a second pair of satellites in an
453 moderately inclined orbit, which greatly stabilizes the gravity field solution due to
454 its rotated sampling sensitivity. Since both the source model and the perturbed
455 de-aliasing model are publicly available, similar experiments might be now per-
456 formed with alternative gravity mission simulation tools in order to increase the
457 robustness of conclusions derived from such assessments.

458 **Acknowledgements** This study was performed under contract No. 4000109421 with the
459 European Space Agency (ESA). We thank Deutscher Wetterdienst, Offenbach, Germany,
460 and the European Centre for Medium-Range Weather Forecasts for providing data from
461 ECMWF's latest reanalysis ERA-Interim. Numerical simulations were performed at Deutsches
462 Klimarechenzentrum, Hamburg, Germany. The updated Earth System Model and the corre-
463 sponding realistically perturbed AOD model as described in this paper are publicly available
464 at doi:10.5880/GFZ.1.3.2014.001.

465 References

- 466 Baldauf, M., A. Seifert, J. Förstner, D. Majewski, M. Raschendorfer, and T. Rein-
467 hardt (2011), Operational Convective-Scale Numerical Weather Prediction with
468 the COSMO Model: Description and Sensitivities, *Mon. Weather Rev.*, *139*,
469 3887–3905, doi:10.1175/MWR-D-10-05013.1.
- 470 Bergmann, I., and H. Dobslaw (2012), Short-term transport variability of the
471 Antarctic Circumpolar Current from satellite gravity observations, *J. Geophys.*
472 *Res.*, *117*(C5), 1–12, doi:10.1029/2012JC007872.
- 473 Bergmann-Wolf, I., E. Forootan, V. Klemann, J. Kusche, and H. Dobslaw (2015),
474 Updating ESA's Earth System Model for Gravity Mission Simulation Studies: 3.
475 A Realistically Perturbed Non-Tidal Atmosphere and Ocean De-Aliasing Model,
476 *Tech. rep.*, Scientific Technical Report 14/09, GFZ, doi:10.2312/GFZ.b103-
477 14091.
- 478 Boy, J.-P. (2005), Precise evaluation of atmospheric loading effects on
479 Earth's time-variable gravity field, *J. Geophys. Res.*, *110*(B8), 1–10, doi:
480 10.1029/2002JB002333.
- 481 Carrère, L., and F. Lyard (2003), Modeling the barotropic response of the global
482 ocean to atmospheric wind and pressure forcing - comparisons with observations,
483 *Geophys. Res. Lett.*, *30*(6), 1997–2000, doi:10.1029/2002GL016473.
- 484 Chambers, D. P. (2004), Preliminary observations of global ocean mass variations
485 with GRACE, *Geophys. Res. Lett.*, *31*(13), 1–4, doi:10.1029/2004GL020461.
- 486 Dee, D. P., et al. (2011), The ERA-Interim reanalysis: configuration and perfor-
487 mance of the data assimilation system, *Q. J. Roy. Met. Soc.*, *137*(656), 553–597,
488 doi:10.1002/qj.828.
- 489 Dobslaw, H., and M. Thomas (2007), Simulation and observation of
490 global ocean mass anomalies, *J. Geophys. Res.*, *112*(C5), C05,040, doi:
491 10.1029/2006JC004035.
- 492 Dobslaw, H., F. Flechtner, I. Bergmann-Wolf, C. Dahle, R. Dill, S. Esselborn,
493 I. Sasgen, and M. Thomas (2013), Simulating high-frequency atmosphere-ocean
494 mass variability for dealiasing of satellite gravity observations: AOD1B RL05,
495 *J. Geophys. Res.*, *118*(7), 3704–3711, doi:10.1002/jgrc.20271.

- 496 Dobslaw, H., I. Bergmann-Wolf, R. Dill, E. Forootan, V. Klemann, J. Kusche,
497 and I. Sasgen (2015), The updated ESA Earth System Model for future gravity
498 mission simulation studies, *J. Geod.*, doi:10.1007/s00190-014-0787-8.
- 499 Duan, J., C. K. Shum, J. Guo, and Z. Huang (2012), Uncovered spurious
500 jumps in the GRACE atmospheric de-aliasing data: potential contamination
501 of GRACE observed mass change, *Geophys. J. Int.*, *2*, 1–5, doi:10.1111/j.1365-
502 246X.2012.05640.x.
- 503 Elsaka, B., J.-C. Raimondo, P. Brieden, T. Reubelt, J. Kusche, F. Flechtner, S. Iran
504 Pour, N. Sneeuw, and J. Müller (2014), Comparing seven candidate mission
505 configurations for temporal gravity field retrieval through full-scale numerical
506 simulation, *J. Geod.*, *88*(1), 31–43, doi:10.1007/s00190-013-0665-9.
- 507 Famiglietti, J. S., and M. Rodell (2013), Water in the balance, *Science (80-.)*,
508 *340*(61), 1300–1301, doi:10.1126/science.1236460.
- 509 Flechtner, F., and H. Dobslaw (2013), GRACE AOD1B Product Description Doc-
510 ument for Product Release 05, *Tech. rep.*, Rev. 4.0, GRACE Document 327-750,
511 GeoForschungsZentrum Potsdam.
- 512 Flechtner, F., P. Morton, M. Watkins, and F. Webb (2014), Status of the GRACE
513 Follow-On Mission, in *IAG Symp. Gravity, Geoid, Height Syst. Int. Assoc. Geod.*
514 *Symp. 141*, edited by U. Marti, Springer, doi:10.0007/978-3-319-10837-7-15.
- 515 Flechtner, F., K.-H. Neumayer, C. Dahle, H. Dobslaw, A. Güntner, J.-C. Rai-
516 mondo, and E. Fagiolini (2015), What can be expected from the GRACE-FO
517 Laser Ranging Interferometer for Earth Science Applications?, *Surv. Geophys.*,
518 doi:10.1007/s10712-015-9338-y.
- 519 Forootan, E., O. Didova, J. Kusche, and A. Löcher (2013), Comparisons of at-
520 mospheric data and reduction methods for the analysis of satellite gravimetry
521 observations, *J. Geophys. Res.*, *118*(1), 2382–2396, doi:10.1002/jgrb.50160.
- 522 Forootan, E., O. Didova, M. Schumacher, J. Kusche, and B. Elsaka (2014), Com-
523 parisons of atmospheric mass variations derived from ECMWF reanalysis and
524 operational fields, over 2003-2011, *J. Geod.*, *88*, 503–514, doi:10.1007/s00190-
525 014-0696-x.
- 526 Gruber, T., et al. (2011), Simulation of the time-variable gravity field by
527 means of coupled geophysical models, *Earth Syst. Sci. Data*, *3*(1), 19–35, doi:
528 10.5194/essd-3-19-2011.
- 529 Gruber, T., et al. (2014), *Earth System Mass Transport Mission (Square) ? Con-*
530 *cept for a Next Generation Gravity Field Mission*, 318, 210 pp., Deutsche
531 Geodätische Kommission, Series B, Bayerische Akademie der Wissenschaften, Mu-
532 nich.
- 533 Han, S.-C., C. Jekeli, and C. K. Shum (2004), Time-variable aliasing effects of
534 ocean tides, atmosphere, and continental water mass on monthly mean GRACE
535 gravity field, *J. Geophys. Res.*, *109*(4), B04,403, doi:10.1029/2003JB002501.
- 536 Kim, J. (2000), Simulation study of a low-low satellite-to-satellite tracking mission,
537 Ph.d. thesis, University of Texas, Austin.
- 538 Koch, K. R., and J. Kusche (2002), Regularization of geopotential determina-
539 tion from satellite data by variance components, *J. Geod.*, *76*(5), 259–268, doi:
540 10.1007/s00190-002-0245-x.
- 541 Kuhlmann, J., H. Dobslaw, C. Petrick, and M. Thomas (2013), Ocean bottom pres-
542 sure signals around Southern Africa from in situ measurements, satellite data,
543 and modeling, *J. Geophys. Res.*, *118*(10), 4889–4898, doi:10.1002/jgrc.20372.

- 544 Kurtenbach, E., T. Mayer-Gürr, and A. Eicker (2009), Deriving daily snapshots
545 of the Earth's gravity field from GRACE L1B data using Kalman filtering,
546 *Geophys. Res. Lett.*, *36*, L17,102, doi:10.1029/2009GL039564.
- 547 Kurtenbach, E., A. Eicker, T. Mayer-Gürr, M. Holschneider, M. Hayn,
548 M. Fuhrmann, and J. Kusche (2012), Improved daily GRACE gravity
549 field solutions using a Kalman smoother, *J. Geod.*, *59-60*, 39–48, doi:
550 10.1016/j.jog.2012.02.006.
- 551 le Bars, Y., F. Lyard, C. Jeandel, and L. Dardengo (2010), The AMANDES tidal
552 model for the Amazon estuary and shelf, *Ocean Model.*, *31*(3-4), 132–149, doi:
553 10.1016/j.ocemod.2009.11.001.
- 554 Loomis, B. D., R. S. Nerem, and S. B. Luthcke (2011), Simulation study of a follow-
555 on gravity mission to GRACE, *J. Geod.*, *86*(5), 319–335, doi:10.1007/s00190-
556 011-0521-8.
- 557 Mayer-Gürr, T., N. Zehentner, B. Klinger, and A. Kvas (2014), ITSG-Grace2014:
558 A new GRACE gravity field release computed in Graz, in *Proc. GRACE Sci.*
559 *Team Meet.*, Potsdam.
- 560 Menemenlis, D., and J.-M. Campin (2008), ECCO2: High resolution global ocean
561 and sea ice data synthesis, *Mercat. Ocean Newsl.*, *1*(10), 13–21.
- 562 Panet, I., et al. (2012), Earth System Mass Transport Mission (e.motion): A Con-
563 cept for Future Earth Gravity Field Measurements from Space, *Surv. Geophys.*,
564 *34*(2), 141–163, doi:10.1007/s10712-012-9209-8.
- 565 Rienecker, M. M., et al. (2011), MERRA: NASA's modern-era retrospective anal-
566 ysis for research and applications, *J. Clim.*, *24*, 3624–3648, doi:10.1175/JCLI-
567 D-11-00015.1.
- 568 Saha, S., et al. (2010), The NCEP climate forecast system reanalysis, *Bull. Am.*
569 *Meteorol. Soc.*, *91*(August), 1015–1057, doi:10.1175/2010BAMS3001.1.
- 570 Sakumura, C., S. Bettadpur, and S. Bruinsma (2014), Ensemble prediction and
571 intercomparison analysis of GRACE time-variable gravity field models, *Geophys.*
572 *Res. Lett.*, pp. 1389–1397, doi:10.1002/2013GL058632.1.
- 573 Sasgen, I., H. Konrad, E. R. Ivins, M. R. Van Den Broeke, J. L. Bamber, Z. Mar-
574 tinec, and V. Klemann (2013), Antarctic ice-mass balance 2003 to 2012: Re-
575 gional reanalysis of GRACE satellite gravimetry measurements with improved
576 estimate of glacial-isostatic adjustment based on GPS uplift rates, *Cryosphere*,
577 *7*, 1499–1512, doi:10.5194/tc-7-1499-2013.
- 578 Tapley, B. D., S. Bettadpur, M. Watkins, and C. Reigber (2004), The gravity
579 recovery and climate experiment: Mission overview and early results, *Geophys.*
580 *Res. Lett.*, *31*(22), 09,607, doi:10.1029/2004GL019920.
- 581 Thomas, M., J. Sündermann, and E. Maier-Reimer (2001), Consideration of ocean
582 tides in an OGCM and impacts on subseasonal to decadal polar motion, *Geo-*
583 *phys. Res. Lett.*, *28*(12), 2457–2460.
- 584 Thompson, P. F. (2004), Impact of short period, non-tidal, temporal mass
585 variability on GRACE gravity estimates, *Geophys. Res. Lett.*, *31*(6), doi:
586 10.1029/2003GL019285.
- 587 Uppala, S. M., et al. (2005), The ERA-40 re-analysis, *Q. J. Roy. Met. Soc.*,
588 *131*(612), 2961–3012, doi:10.1256/qj.04.176.
- 589 van Dam, T., Z. Altamimi, X. Collilieux, and J. Ray (2010), Topographically
590 induced height errors in predicted atmospheric loading effects, *J. Geophys. Res.*,
591 *115*(B7), B07,415, doi:10.1029/2009JB006810.

- 592 Visser, P. N. A. M. (2010), Designing Earth Gravity Field Missions for the Future:
593 A Case Study, in *IAG Comm. 2 GRAVITY, GEOID EARTH Obs. Chania*, pp.
594 131–138.
- 595 von Storch, J.-S., C. Eden, I. Fast, H. Haak, D. Hernández-Deckers, E. Maier-
596 Reimer, J. Marotzke, and D. Stammer (2012), An Estimate of the Lorenz Energy
597 Cycle for the World Ocean Based on the STORM/NCEP Simulation, *J. Phys.*
598 *Ocean.*, *42*(12), 2185–2205, doi:10.1175/JPO-D-12-079.1.
- 599 Wiese, D. N., W. M. Folkner, and R. S. Nerem (2009), Alternative mission archi-
600 tectures for a gravity recovery satellite mission, *J. Geod.*, *83*, 569–581, doi:
601 10.1007/s00190-008-0274-1.
- 602 Wiese, D. N., R. S. Nerem, and F. G. Lemoine (2012), Design considerations for a
603 dedicated gravity recovery satellite mission consisting of two pairs of satellites,
604 *J. Geod.*, *86*, 81–98, doi:10.1007/s00190-011-0493-8.
- 605 Zenner, L., T. Gruber, A. Jäggi, and G. Beutler (2010), Propagation of atmospheric
606 model errors to gravity potential harmonics-impact on GRACE de-aliasing, *Geo-*
607 *phys. J. Int.*, *182*(2), 797–807, doi:10.1111/j.1365-246X.2010.04669.x.
- 608 Zenner, L., I. Bergmann-Wolf, H. Dobslaw, T. Gruber, A. Güntner, M. Watten-
609 bach, S. Esselborn, and R. Dill (2014), Comparison of Daily GRACE Gravity
610 Field and Numerical Water Storage Models for De-aliasing of Satellite Gravime-
611 try Observations, *Surv. Geophys.*, doi:10.1007/s10712-014-9295-x.