



Originally published as:

Reinsch, T., Thurley, T., Jousset, P. (2017): On the mechanical coupling of a fiber optic cable used for distributed acoustic/vibration sensing applications — a theoretical consideration. - *Measurement Science and Technology*, 28, 12.

DOI: <http://doi.org/10.1088/1361-6501/aa8ba4>

1 **On the mechanical coupling of a fiber optic cable**
2 **used for distributed acoustic/vibration sensing**
3 **applications - a theoretical consideration**

4 **Thomas Reinsch, Tom Thurley, Philippe Jousset**

5 GFZ German Research Centre for Geosciences, Telegrafenberg, Potsdam, Germany

6 E-mail: Thomas.Reinsch@gfz-potsdam.de

7 **Abstract.** In recent years, fiber optic cables are increasingly used for the acquisition
8 of dynamic strain changes for seismic surveys. When considering seismic amplitudes,
9 one of the first questions arising is the mechanical coupling between optical fiber and
10 the surrounding medium. Here we analyse the interaction of ground movement with
11 a typical telecom-grade fiber optic cable from an existing telecommunication network
12 deployed in a sand filled trench at the surface. Within the cable, the optical fiber is
13 embedded in a gel filled plastic tube. We apply Hooke's law to calculate the stress
14 needed to strain the optical fiber throughout the cable structure. In case the stress
15 magnitude at the cable-sand interface as well as the gel-optical fiber interface is below
16 the yield strength of the respective material, sand and gel, it can be regarded as an
17 elastic medium. Hence, a multilayer radial symmetric model can be used to calculate
18 the coupling of the optical fiber with the surrounding medium. We show that the
19 transfer function has a -3 dB lower cut-off wavelength of about 22 m. The magnitude
20 response of this telecom-grade fiber optic cable is therefore almost perfect at typical
21 low frequency seismic waves. The approach presented here can be applied to various
22 cable designs to estimate the strain transfer between ground movement and an optical
23 fiber.

24 PACS numbers: 07.60.Vg,93.85.Rt,91.30.Hc,91.30.Dk

25 Submitted to: *Meas. Sci. Technol.*

26 *Keywords:* distributed acoustic/vibration sensing, coupling, seismic survey, seismology,
27 instrumental response

1. Introduction

Fiber optic cables are increasingly used for the acquisition of dynamic strain changes in geophysical applications (e.g. Mestayer et al., 2011; Daley et al., 2013; Götz et al., 2017). This technology is often referred to as distributed acoustic sensing (DAS) or distributed vibration sensing (DVS) (Hartog et al., 2013), where phase-OTDR (optical time domain reflectometry, φ -OTDR) data is used to measure local strain changes at several positions along an optical fiber (Masoudi and Newson, 2016). A coherent laser pulse is coupled into the optical fiber and the phase difference of elastically backscattered photons from neighbouring positions along the fiber can be analysed. When analysing local strain changes, the mechanical coupling of the sensing fiber with the surrounding environment has to be known.

Several studies analysing the strain transfer across to an optical fiber exist (e.g. Pak, 1992; Ansari and Libo, 1998; Schmidt-Hattenberger et al., 2003; Li et al., 2006, 2009). Here, we apply and extend the approach of Li et al. (2006) to analyse the mechanical coupling of an optical fiber inside a standard telecommunication cable deployed in the subsurface. Inside the cable structure, the optical fiber is embedded in a gel, and the gel is confined in a plastic tube. To increase the strength of the fiber optic cable, steel wires are used as strength members. In addition, a plastic outer layer protects the cable from e.g. alteration and corrosion (Refi, 1998). In our application, the fiber optic telecommunication cable is embedded in a trench and covered with sand (Figure 1).

The study is motivated by the continuous acquisition of DAS data along a 15.3 km fiber optic cable within the telecommunication network of Iceland (Reinsch et al., 2016; Jousset et al., 2016) where the mechanical coupling between the optical fiber and the subsurface is unknown, and the cable was not accessible for calibration measurements. Data was collected in the framework of the EC funded FP7 project IMAGE (Integrated Methods for Advanced Geophysical Monitoring) van Wees et al. (2015). To interpret the data, we estimated the strain transfer from the surrounding rock mass to the optical fiber. Information about the multi-layer cable design and its deployment in the 1990's is available, the mechanical properties of individual layers like the gel or the unconsolidated sand surrounding the optical cable, however, are unknown and assumptions have to be made. Evaluating the mechanical properties of sand and gel, we were able to populate the model of Li et al. (2006) and calculate the strain transfer to the optical fiber for seismic applications.

2. Theory

In order to estimate the mechanical coupling between fiber optic cable and the surrounding medium, the force applied to the cable needed to create the measured strain has to be calculated. Assuming elastic rheology, Hooke's Law (Hooke, 1678) is applied:

$$\sigma = \epsilon E \tag{1}$$

67 where σ is the stress, ϵ the strain and E the Young's modulus. The stress σ is
 68 defined as:

$$69 \quad \sigma = \frac{F}{A} = \epsilon E = \epsilon \frac{k}{A} \quad (2)$$

70 where F is the force applied to a cross-sectional area A and $k = AE$ is defined
 71 as the axial stiffness. For a multilayer cable, the axial stiffness per unit length can be
 72 written as:

$$73 \quad k = \sum_{i=1}^n k_i \quad (3)$$

74 where k_i is the is the axial stiffness for each component i .

75 *2.1. Mechanical coupling with surrounding medium*

76 Here, we consider a fiber optic communication cable, buried in a trench and embedded
 77 in sand at a depth of about 1 m. To evaluate whether the sand is plastically or elastically
 78 deformed from seismic waves passing the trench, the force necessary to overcome the
 79 yield strength of sand at the sand-cable interface is estimated. We assume that (a) the
 80 sand can be treated as a homogeneous medium, (b) the contact force between cable and
 81 sand is created due to the lithostatic stress, i.e. the weight of the overburden sand, (c)
 82 cohesion between individual sand grains is neglected. According to Mohr (1900), the
 83 shear strength τ is defined by the Mohr-Coulomb failure criterion:

$$84 \quad \tau = c + \sigma_N \tan(\varphi) \quad (4)$$

85 where σ_N is the lithostatic stress normal to the medium, c is a measure of the internal
 86 forces and φ is the friction angle. Typical friction angles for non-cohesive sand are
 87 in the order of 30-40° (Grabe, 2012). The lithostatic stress was calculated according to
 88 the density ρ of the overburden material and the depth of burial h :

$$89 \quad \sigma_N = \rho gh \quad (5)$$

90 For non-cohesive sand ($c = 0$), typical densities are in the range of 1600-2000 kg/m³
 91 (Grabe, 2012). Consequently, in 1 m depth, the shear strength of the sand can be
 92 calculated to be in the range of 9000-16500 Pa. For the interface friction angle of a flat
 93 sample of 8in HDPE pipe versus Fraser River sand, Huber (2014) measured about 26°
 94 in a direct shear test. The shear strength of the interface between cable and sand is
 95 therefore above 7500 Pa.

96 *2.2. Viscosity*

97 Optical fibers inside a fibre optic cable are often embedded in non-Newtonian, oil-based
 98 visco-elastic fluid inside a tube. To maintain optical transmission properties over a
 99 long time, such fluids are designed to isolate the optical fibers from external forces and
 100 to avoid any water migration along the cable. On the one hand, such a gel shall be
 101 low viscous to maintain fiber mobility within the cable as well as to allow for an easy

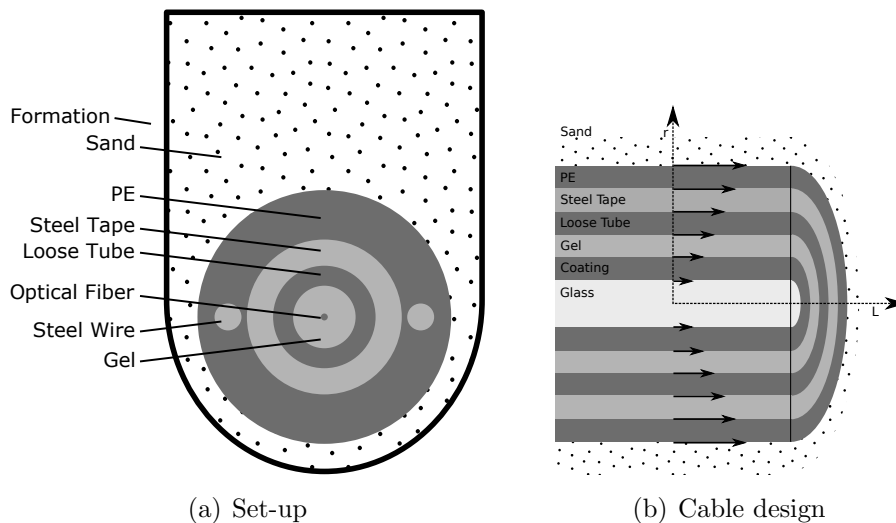


Figure 1. (a) Schematic view of a multilayer fiber optic cable embedded in a sand filled trench. (b) Schematic cross sectional view of the fiber optic cable. Arrows indicates the strain transfer across the cable to the optical fiber. L denotes the length of the cable under investigation and r is the radial distance from the center; dimensions not to scale.

102 intrusion of the fluid into the tube during cable manufacturing. On the other hand, a
 103 high viscosity is desired to avoid the fluid from dripping out of the tube once intruded
 104 into it. As a gel with low viscosity is mechanically unstable in a cable structure, a high
 105 viscosity gel reduces the ability of the fiber to quickly move relative to the plastic tube
 106 in response to external forces. A thixotropic fluid is generally used to protect the optical
 107 fibers during movement of the cable (e.g. Khan et al., 1991). For seismic applications,
 108 the time of excitation is very short, the thixotropic behaviour will therefore be neglected
 109 in the following.

110 Neglecting creep (e.g. plastic deformation of the gel below the yield strength),
 111 different rheological models can be applied to describe such non-Newtonian fluids. One
 112 of the most simple models to describe a non-Newtonian fluid is the Bingham fluid
 113 model (Bingham, 1916), where a yield strength τ_0 of the gel has to be overcome before
 114 it behaves like a viscous fluid:

$$115 \quad \tau = \tau_0 + \dot{\gamma}\nu_p \quad (6)$$

116 where τ is the resulting stress to generate a shear rate $\dot{\gamma}$ for a gel with a plastic viscosity
 117 ν_p . Another model, taking into account the non linear plastic viscosity of the gel, is the
 118 Hershel-Bulkley model (Herschel and Bulkley, 1926):

$$119 \quad \tau = \tau_0 + K\dot{\gamma}^n \quad (7)$$

120 where K the consistency factor and n the flow behaviour index. Again, the yield strength
 121 τ_0 has to be overcome before the gel behaves like a fluid. Below the yield strength, the
 122 fluid behaves like an elastic medium. Typically, gels are engineered to have a yield point
 123 of 35-140 Pa (Evonik Industries, 2015). A linear relation between yield point and the

124 elastic modulus of typical gels can be observed. For yield points greater than 30 Pa, an
 125 elastic modulus of more than 100 Pa can be assumed (Evonik Industries, 2015).

126 2.3. Strain Transfer

127 The strain transfer between the surface of the cable and the optical fiber can be evaluated
 128 according to the analytical model of Li et al. (2006). Li et al. (2006) take into account
 129 an optical fiber embedded in a perfectly bonded multilayer, radial symmetric design of
 130 elastic materials, where stress is uniformly applied to the outside of the system. In this
 131 model the boundary condition for the strain measured along the optical fiber is zero at
 132 both ends of the investigated cable section. The average strain transfer rate $\bar{\alpha}$ can be
 133 estimated according to:

$$134 \quad \bar{\alpha} = 1 - \frac{\sinh(\beta \frac{L}{2})}{\beta \frac{L}{2} \cosh(\beta \frac{L}{2})} \quad (8)$$

135 where L is the length of the multilayer system under stress and β is the shear lag
 136 parameter:

$$137 \quad \beta^2 = \frac{2}{r_g^2 E_g \left\{ \sum_{i=1}^n \frac{1}{G_i} \ln\left(\frac{r_i}{r_{i-1}}\right) \right\}} \quad (9)$$

138 where r_g and E_g are the outer radius and Young's modulus of the cladding, respectively.
 139 For simplicity, it is assumed that core and clad have the same mechanical properties.
 140 r_i and G_i is the outer radius and shear modulus of the i -th layer and $i=1$ denotes the
 141 cladding material.

142 2.4. Experimental set-up in Iceland

143 In our study the cable has a centrally situated loose tube with a reasonable excess length
 144 that minimizes the influence of lateral crush. A double wire acts as strength member
 145 to provide sufficient strain resistance. In Iceland, an about 1 m deep trench was dug in
 146 solidified basaltic lava. The fiber optic cable was embedded in sand and the trench was
 147 re-filled with sand and the excavated lava. Figure 1 shows the model set-up used in this
 148 study.

149 To calculate the mechanical properties of the fiber optic cable, cable dimensions
 150 are estimated based on commercially available GYXTW cable types (Table 1). For
 151 simplicity, the properties of the steel tape are assumed to be similar to the properties
 152 of steel.

153 For seismic applications, relevant frequencies are typically ranging from 0.1-100 Hz
 154 (Aki and Richards, 2002). When a seismic wave passes by, the cable is compressed and
 155 stretched with an amount depending on the angle of incidence for the seismic wave as
 156 well as the type wave. During our measurement campaign in Iceland, strain magnitudes
 157 typically much lower than $1 \mu\text{m}/\text{m}$ were observed.

158 For every site, the mechanical coupling between the subsurface and the optical fiber
 159 has to be evaluated. The workflow to evaluate the strain transfer between surrounding

Table 1. Design of a GYXTW cable together typical young’s moduli (E) are given. Cited references are: ¹Ansari and Libo (1998),²Evonik Industries (2015),³Kern GmbH (2016a),⁴Kern GmbH (2016a), ⁵Kuchling (2011),⁶Kern GmbH (2016b). The shear modulus of PE, PBT and the gel was calculated according to $E = 2G(1 + \nu)$ with a Poisson’s ratio of $\nu = 0.5$.

Component	Material	Outer Radius (mm)	E (kN/mm ²)	G (kN/mm ²)
Fiber Core	Glas	0.004	72 ¹	
Fiber Clad	Glas	0.0625	72 ¹	
Fiber Coating	Silicon	0.250	2.55e ⁻³ ¹	8.50e ⁻³ ¹
Tube Filling Compound	Gel	0.9	1e ⁻⁷ ²	3.3e ⁻⁸
Loose Tube	PBT	1.3	2.5 ³	0.83
Steel Tape	Steel	1.43	211 ⁵	80 ⁵
Outer Sheath	PE	4.25	1.35 ⁶	0.45
2 Wires	Steel	0.8	211 ⁵	80 ⁵

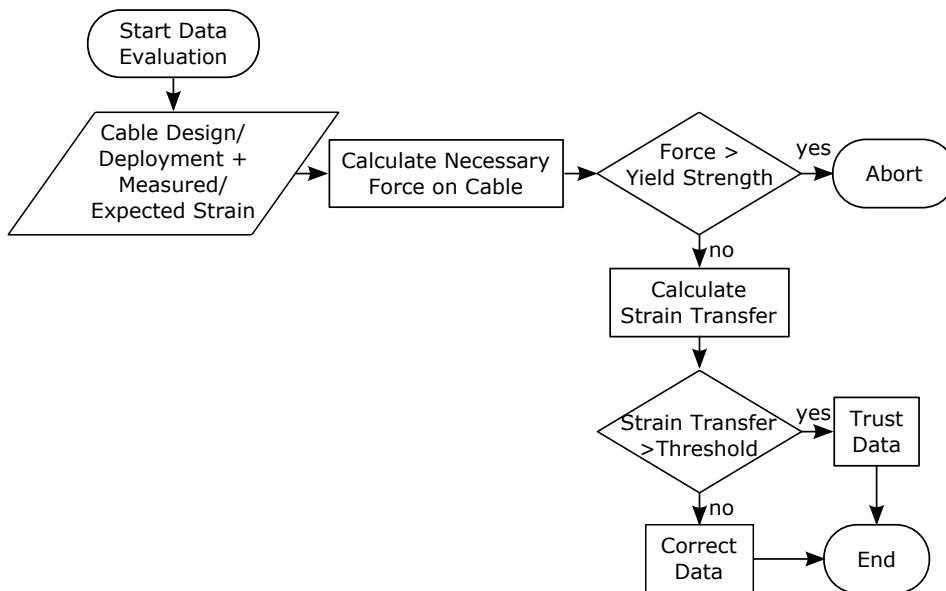


Figure 2. Flow chart showing the procedure to evaluate the mechanical coupling between the subsurface and the optical fiber.

160 medium and optical fiber is detailed in Figure 2. The single steps are based on equations
 161 1 - 9. The threshold in Figure 2 refers to ground motions, where the mechanical coupling
 162 influences measured amplitudes significantly. In our study, we set the threshold to -
 163 3 dB. Below the threshold, data must be corrected to compensate the influence of poor
 164 coupling.

165 3. Results

166 According to the specifications of the fiber optic cable listed in Table 1, the force needed
167 to stretch the cable to $1 \mu\text{m}/\text{m}$ is about 1.2 N (Equations 2 and 3). This force is applied
168 at the surface of the fiber optic cable. The corresponding cable length is more than 10 m
169 for seismic frequencies we observed in Iceland (0.1-50 Hz) and subsurface velocities larger
170 than 700 m/s Raab et al. (2017).

171 3.1. Mechanical coupling of cable to ground

172 For a length of more than 10 m, the surface area of the cable is at least 0.28 m^2 (see
173 Table 1). The shear stress at the interface is therefore greater than 4.5 Pa. This value
174 is three orders of magnitude lower than the shear stress needed to shear non-cohesive
175 sand (see Section 2.1).

176 3.2. Mechanical coupling of optical fiber to cable

177 Equation 3 can be used to calculate the force acting at the surface of the optical fiber,
178 i.e., at the interface between optical fiber and tube filling compound, here a gel, as well
179 as at the interface between gel and loose tube. With regards to the surface area at each
180 interface, the acting shear stresses are calculated according to Equation 2 to be 0.12 Pa
181 and 0.02 Pa, respectively. The yield point of the gel is more than 2 orders of magnitude
182 larger (see Section 2.2). Hence, the gel behaves elastically.

183 3.3. Strain Transfer

184 The average strain transfer from the outer surface of the cable to the optical fiber
185 was calculated according to Equation 8. As the steel wires are embedded in the outer
186 sheath and are not in direct contact with the next inner layer, they are neglected for
187 the calculation of the strain transfer. In case of a cable design as listed in Table 1 and a
188 cable length of 10 m, about 45 % of the strain applied to the outer surface of the cable is
189 transmitted to the optical fiber. For a 100 m long cable, an average of about 95 % of the
190 strain is transmitted to the optical fiber (Figure 3 (a)). For passive seismic recordings,
191 frequencies of less than 10 Hz are of major interest. Assuming a propagation velocity
192 of 3000 m/s, the wavelength is larger than 300 m. For a longitudinal wave half of this
193 length is in compression, half is in extension. Hence, for a cable section of 150 m, on
194 average, about 96 % of the strain applied to the surface of the cable are transmitted to
195 the optical fiber. Figure 3 (b) shows the strain transfer for different wavelength. The
196 -3 dB cut-of wavelength for the transfer function is at a wavelength of about $\approx 22 \text{ m}$.

197 4. Discussion

198 Within this study, we used rather conservative estimates on the mechanical properties
199 of individual components to analyse the lower boundary of the strain transfer. For the

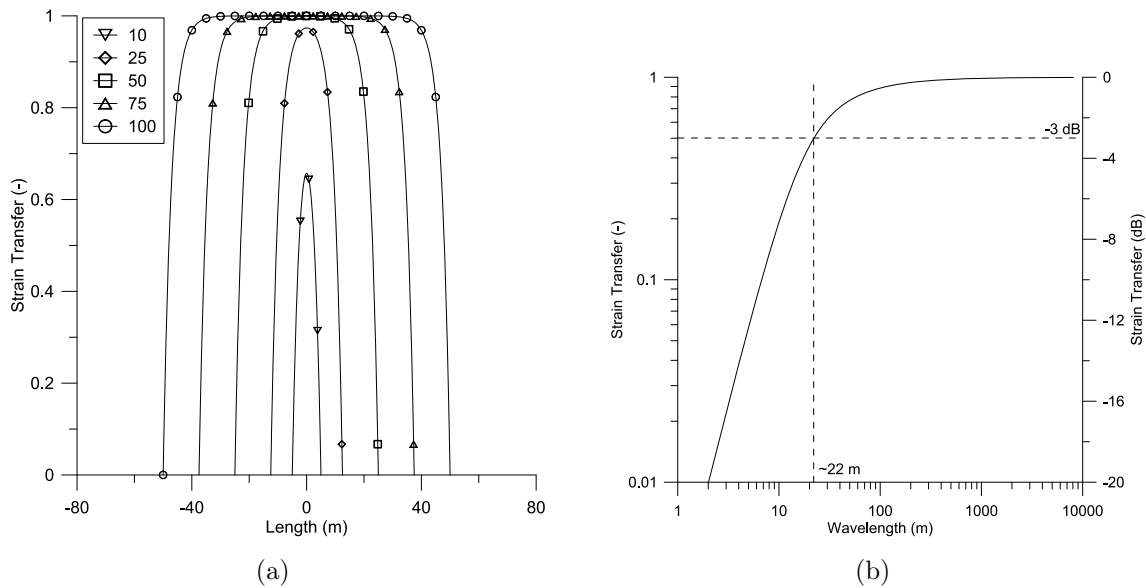


Figure 3. Strain transfer to the optical fiber for the cable design as described in the text. a) Distributed strain transfer along the cable for different cable lengths (10, 25, 50, 75 and 100 m). b) Average strain transfer at different wavelength.

200 evaluation, a perfect mechanical coupling at all interfaces was assumed. The estimated
 201 yield strength of the gel is at the lower boundary of what is reported as an industry
 202 standard (Evonik Industries, 2015, and references therein). Generally, optical fibers
 203 are overstuffed in a fiber optic cable. In this study, the overstuffing was neglected. In
 204 principle, an excess fiber length leads to an increased surface area between optical fiber
 205 and gel as well as an increase in the number of contact points between optical fiber
 206 and tube wall. Both effects would lead to an increased mechanical coupling between
 207 cable structure and fiber as the gel is the weakest component in the set-up and a elastic
 208 behaviour has to be guaranteed. However, even with the properties as assumed in this
 209 study, the yield strength of the gel is much higher than the expected shear stress.

210 For the study presented here, the thixotropic behaviour of the gel was neglected.
 211 The presented strain transfer, therefore, is valid for a continuous seismic excitation
 212 with a maximum strain amplitude of $1 \mu\text{m}/\text{m}$. Taking into account the thixotropic
 213 behaviour of the gel, hence an increased viscosity at the onset of an excitation, a higher
 214 strain transfer magnitude can be assumed for event-type excitations, i.e. earthquakes.

215 A detailed documentation of the deployed cable in Iceland was not available. Hence,
 216 the design of the cable was compared to commercially available cable types. For the
 217 calculation, therefore, a GYXTW type of cable was used. Diameters and properties of
 218 the cable were approximated using a standard cable geometry and standard material
 219 properties for each layer. The type of deployment, i.e. embedding in sand, was assumed
 220 as a detailed documentation was not accessible.

221 For seismological applications, the wavelengths of interest are generally rather long.
 222 Using the experimental set-up as described in Section 2.4, the transfer function has a

223 -3 dB cut-of wavelength of about 22 m.

224 5. Summary and Conclusions

225 The mechanical coupling between an optical fiber, incorporated in a buried fiber optic
226 cable, and the subsurface was analysed. A simple analytical model was used to
227 calculate the forces needed to stretch the cable. For typical magnitudes of low frequency
228 ground motions, an almost perfect mechanical coupling can be assumed. The presented
229 approach can be used to estimate the transfer function of different fiber optic cable
230 types for seismic applications and beyond.

231 Acknowledgments

232 This project has received funding from the European Union's Seventh Framework
233 Programme for research, technological development and demonstration under grant
234 agreement no 608553 (Project IMAGE). The authors would like to thank three
235 anonymous reviewers for their valuable contribution to improve the manuscript.

236 References

- 237 Aki, K. and P. G. Richards (2002). *Quantitative Seismology* (2 ed.). University Science
238 Books.
- 239 Ansari, F. and Y. Libo (1998). Mechanics of bond and interface shear transfer in optical
240 fiber sensors. *Journal of Engineering Mechanics* 124(4), 385–394.
- 241 Bingham, E. C. (1916). An investigation of the laws of plastic flow. *Bulletin of the*
242 *Bureau of Standards* 13, 309–353.
- 243 Daley, T. M., B. M. Freifeld, J. Ajo-Franklin, S. Dou, R. Pevzner, V. Shulakova,
244 S. Kashikar, D. E. Miller, J. Goetz, J. Henninges, and S. Lueth (2013). Field testing
245 of fiber-optic distributed acoustic sensing (DAS) for subsurface seismic monitoring.
246 *The Leading Edge* 32(6), 699–706.
- 247 Evonik Industries (2015). AEROSIL fumed silica for cable gels. Technical Information
248 1163, Evonik Industries.
- 249 Grabe, J. (2012). *Empfehlungen des Arbeitsausschusses "Ufereinfassungen" Häfen und*
250 *Wasserstraßen EAU 2012*. Ernst W. + Sohn Verlag.
- 251 Götz, J., S. Lueth, J. Henninges, and T. Reinsch (2017). Vsp imaging using fibre
252 optic cable as receiver array in four wells simultaneously at the ketzin co2 storage
253 pilot site. In *Proceedings, EAGE/DGG Workshop 2017, Potsdam, Germany*. EAGE
254 Publications BV.
- 255 Hartog, A. H., O. Kotov, and L. Liokumovich (2013). The optics of distributed vibration
256 sensing. In *Second EAGE Workshop on Permanent Reservoir Monitoring 2013 –*
257 *Current and Future Trends*. EAGE.

- 258 Herschel, W. H. and R. Bulkley (1926). Konsistenzmessungen von Gummi-
259 Benzollösungen. *Kolloid-Zeitschrift* 39(4), 291–300.
- 260 Hooke, R. (1678). *Lectures de Potentia Restitutiva, Or of Spring Explaining the Power*
261 *of Springing Bodies*. St. Pauls Church-Yard, London: John Martyn.
- 262 Huber, M. (2014). An experimental investigation of thermal effects on the axial
263 resistance to relative ground movement of buried district heating pipes. Master’s
264 thesis, University of British Columbia.
- 265 Jousset, P., T. Reinsch, J. Henninges, H. Blanck, and T. Ryberg (2016). Strain and
266 ground-motion monitoring at magmatic areas: ultra-long and ultra-dense networks
267 using fibre optic sensing systems. 18(EGU2016-15707).
- 268 Kern GmbH (2016a). Datenblatt Polybutylenterephthalat (PBT).
- 269 Kern GmbH (2016b). Datenblatt Polyethylen hoher Dichte (PE-HD).
- 270 Khan, S. A., M. A. Maruca, and I. M. Plitz (1991). Rheology of fumed silica dispersions
271 for fiber-optic cables. *Polymer Engineering & Science* 31(24), 1701–1707.
- 272 Kuchling, H. (2011). *Taschenbuch der Physik: mit zahlreichen Tabellen*. Fachbuchverl.
273 Leipzig im Carl-Hanser-Verlag.
- 274 Li, D., H. Li, L. Ren, and G. Song (2006). Strain transferring analysis of fiber bragg
275 grating sensors. *Optical Engineering* 45(2), 024402–024402–8.
- 276 Li, W., C. Cheng, and Y. Lo (2009). Investigation of strain transmission of surface-
277 bonded fbgs used as strain sensors. *Sensors and Actuators A: Physical* 149(2), 201 –
278 207.
- 279 Masoudi, A. and T. P. Newson (2016). Contributed review: Distributed optical fibre
280 dynamic strain sensing. *Review of Scientific Instruments* 87(011501), 9.
- 281 Mestayer, J., B. Cox, P. Wills, D. Kiyashchenko, J. Lopez, M. Costello, S. Bourne,
282 G. Ugueto, R. Lupton, G. Solano, D. Hill, and A. Lewis (2011). Field trials of
283 distributed acoustic sensing for geophysical monitoring. *SEG Technical Program*
284 *Expanded Abstracts* 30(1), 4253–4257.
- 285 Mohr, O. (1900). Welche Umstände bedingen die Elastizitätsgrenze und den Bruch eines
286 Materials? *Zeitschrift des Vereins Deutscher Ingenieure* 44, 1524–1530.
- 287 Pak, Y. E. (1992). Longitudinal shear transfer in fiber optic sensors. *Smart Materials*
288 *and Structures* 1(1), 57.
- 289 Raab, T., T. Reinsch, P. Jousset, and C. Krawczyk (2017). Multi-station analysis of
290 surface wave dispersion using distributed acoustic fiber optic sensing. In *Proceedings,*
291 *EAGE/DGG Workshop 2017, Potsdam, Germany*. EAGE Publications BV.
- 292 Refi, J. J. (1998). *Fiber Optic Cable*. Geneva, USA: abc Tele Training, Inc.
- 293 Reinsch, T., P. Jousset, J. Henninges, and H. Blanck (2016). Distributed acoustic sensing
294 technology in magmatic geothermal areas - first results from a survey in iceland. In
295 *Proceedings, European Geothermal Congress 2017, Strasbourg*.

- 296 Schmidt-Hattenberger, C., M. Naumann, and G. Borm (2003). Fiber bragg grating
297 strain measurements in comparison with additional techniques for rock mechanical
298 testing. *IEEE Sensors Journal* 3(1), 50–55.
- 299 van Wees, J.-D., J. Hopman, C. Dezayes, R. Vernier, A. Manzella, D. Bruhn, M. Scheck-
300 Wenderoth, O. Flovenz, G. P. Hersir, S. Halldosdottir, and D. Liotta (2015, 4).
301 Image: the eu funded research project integrated methods for advanced geothermal
302 exploration. In *Proceedings World Geothermal Congress 2015*, Melbourne, Australia.