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The forecasting skill of physics-based seismicity models
during the 2010-2012 Canterbury, New Zealand,
earthquake sequence

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Abstract

The static Coulomb stress hypothesis is a widely known physical mechanism for earthquake triggering, and thus a prime candidate for physics-based Operational Earthquake Forecasting (OEF). However, the forecast skill of Coulomb-based seismicity models remains controversial, especially in comparison to empirical statistical models. A previous evaluation by the Collaboratory for the Study of Earthquake Predictability (CSEP) concluded that a suite of Coulomb-based seismicity models were less informative than empirical models during the aftershock sequence of the 1992 M_w 7.3 Landers, California, earthquake. Recently, a new generation of Coulomb-based and Coulomb/statistical hybrid models were developed that account better for uncertainties and secondary stress sources. Here, we report on the performance of this new suite of models in comparison to empirical Epidemic Type Aftershock Sequences (ETAS) models during the 2010-2012 Canterbury, New Zealand, earthquake sequence. Comprising the 2010 M 7.1 Darfield earthquake and three subsequent $M \geq 5.9$ shocks (including the February 2011 Christchurch earthquake), this sequence provides a wealth of data (394 $M \geq 3.95$ shocks). We assessed models over multiple forecast horizons (1-day, 1-month and 1-year, updated after $M \geq 5.9$ shocks). The results demonstrate substantial improvements in the Coulomb-based models. Purely physics-based models have a performance comparable to the ETAS model, and the two Coulomb/statistical hybrids perform better or as well as the corresponding statistical model. On the other hand, an ETAS model with anisotropic (fault-based) aftershock zones is just as informative. These results provide encouraging evidence for the predictive power of Coulomb-based models. To assist with model development, we identify discrepancies between forecasts and observations.

Introduction

Recent earthquakes in Italy, New Zealand, Japan and Nepal have demonstrated that forecasts of the space time evolution of seismic sequences provide information that can expand seismic risk reduction strategies beyond building codes, and enhance preparedness and re-

29 silence. This is the main goal of Operational Earthquake Forecasting (OEF) introduced by
30 the International Commission on Earthquake Forecasting (*Jordan et al.*, 2011) appointed by
31 the Italian government after the 2009 L'Aquila, Italy, earthquake.

32 For these applications, forecasts should be consistent with future seismicity and they
33 should be the most skilful amongst alternatives (i.e. perform better than other forecasts,
34 according to well defined quantitative measures such as the information gain). The evaluation
35 of consistency and skill of forecast models is the main goal of the Collaboratory for the Studies
36 of Earthquake Predictability (CSEP, *Jordan* (2006); *Zecher et al.* (2010)).

37 To date, the first results of a prospective CSEP experiment (*Nanjo et al.*, 2012), retro-
38 spective CSEP experiments (*Woessner et al.*, 2011; *Rhoades et al.*, 2016), and applications
39 to ongoing earthquake sequences (*Marzocchi et al.*, 2017; *Kaiser et al.*, 2017) showed that
40 statistical models of clustered seismicity like the epidemic-type aftershock sequence mod-
41 els (ETAS, *Ogata*, 1998) and the short-term earthquake probability models (*Gerstenberger*
42 *et al.*, 2005, STEP) provide informative forecasts of future seismicity. In our view, these rep-
43 resent the first generation of earthquake forecasting models, and a benchmark for measuring
44 any improvements in forecasting capability.

45 Ongoing model development aims to improve the skill of the forecasts (e.g. *Field et al.*,
46 2015; *Segou et al.*, 2013). One of the most promising approaches is based on Coulomb stress
47 transfer, the most widely accepted mechanism for aftershock triggering (e.g. *Stein et al.*,
48 1992; *King et al.*, 1994; *Toda et al.*, 1998). The predictive power of this hypothesis, however,
49 remains a subject of debate (*Hardebeck et al.*, 1998; *Marsan*, 2003). To date, most evalu-
50 ations of the Coulomb hypothesis are retrospective, with stress changes often calculated at
51 the locations of subsequent events without considering locations which experienced positive
52 Coulomb stress changes without an increase in seismicity. There is a need to rigorously eval-
53 uate the Coulomb hypothesis (*Strader and Jackson*, 2014; *Toda and Enescu*, 2011). When
54 coupled with Dieterich's rate-state friction formulation (or another framework for convert-
55 ing stress to seismicity), Coulomb-based models can generate probability forecasts, enabling
56 evaluations of forecast reliability and skill against alternative models.

57 A previous CSEP evaluation of the predictive skills of forecast models during the 1992
58 M_w 7.3 Landers earthquake sequence found that the Coulomb-based models performed worse
59 than statistical models (*Woessner et al.*, 2011), even though they were comparable at short
60 times after the mainshock. Subsequent studies have confirmed that physical models have a
61 lower overall performance than statistical ones, but they can be comparable for at short times
62 after the mainshock, and beyond the near-source region (*Segou et al.*, 2013). To increase our
63 understanding of the physics of triggering, it is important to understand whether the poor
64 performance is due to a failure of the Coulomb stress hypothesis - i.e., static stress changes
65 are not an important mechanism for aftershock triggering - or whether the implementations
66 of the hypothesis involved inappropriate model choices. For instance, large uncertainties
67 exist in Coulomb stress calculations, due to errors in the slip models and receiver fault
68 orientations (e.g. *Steacy et al.*, 2005; *Hainzl et al.*, 2010, 2009). Here, we test recently
69 developed Coulomb models designed to address some of these issues (*Cattania et al.*, 2014).

70 We investigate the forecasting consistency and skill of this new generation of physics-
71 based forecasting models, as well as new non-parametric models and hybrid Coulomb/statistical
72 models, during the 2010-2012 Canterbury earthquake sequence. The September 3, 2010,
73 M 7.1 Darfield earthquake initiated a vigorous and damaging aftershock sequence, including
74 the damaging M 6.2 Christchurch earthquake in February 2011 (Fig. 1).

75 CSEP Experiment Design

76 Before submitting models, participants agreed on forecast formats, target data and perfor-
77 mance measures. Three forecast horizons were considered (1-day, 1-month, 1-year); models
78 update their forecasts at the end of each forecast horizon, and after each of the four $M \geq 5.9$
79 earthquakes of the sequence (Fig. 1). We test the effect of data quality with three data-
80 availability scenarios. In the first scenario, most interesting scientifically, models were pro-
81 vided best-available data (a reviewed earthquake catalog, focal mechanisms, and published
82 slip models) to generate forecasts. In the second scenario, the slip models were provided with
83 a 10-day delay to mimic delays in finalising a slip model; no slip models were provided in

84 the first 10 days. In the third scenario, only preliminary data were made available, namely
85 preliminary slip models and catalogs, to mimic the real-time situation of operational earth-
86 quake forecasting. All scenarios were evaluated against the best available earthquake catalog
87 data. For brevity, here we focus on the results from the two extreme setups (scenario 1 and
88 3), and present all results in the electronic supplement.

89 Forecasts were specified as numbers of earthquakes in space and magnitude bins (*Schor-*
90 *lemmer et al.*, 2007). The spatial region extends between 170.5° and 174.0° longitude, and
91 -44.5° and -42.5° latitude, and a single layer extending to 40 km depth. Spatial cells are
92 0.05° by 0.05° wide. Magnitude bins are 0.1 units wide, starting from M 3.95; the last bin
93 has no upper bound.

94 **Data**

95 The data sets associated with the three data-input scenarios (best-available, delayed best-
96 available, and near real-time) are summarized in Table: 1 and shown in Fig. 2 and Fig.S1,
97 S2 of the electronic supplement. The target data set comprises 394 $M \geq 3.95$ earthquakes
98 between the September 3, 2010 (UTC), M 7.1 Darfield earthquake and the end of the experi-
99 ment at midnight on February 29, 2012 (the last date of reviewed data available at experiment
100 conception). The catalog was later reviewed by GeoNet; magnitudes were initially given as
101 local magnitudes, and later replaced by moment magnitudes when available. The input data
102 sets for the scenario with best available data comprise (i) the reviewed GeoNet catalog, (ii)
103 published slip models of the four large earthquakes (*Beavan et al.*, 2012), and (iii) a GeoNet
104 focal mechanism catalog. The same data sets are provided in the second scenario, except
105 that slip models are provided to models 10 days after each of the four largest quakes. The
106 data sets of the near-real-time data scenario include (i) a very preliminary GeoNet cata-
107 log that was downloaded intermittently by one of us during the sequence (*Christophersen,*
108 *private communication*) and (ii) preliminary slip models (*Holden et al.*, 2011, ; Beaven and
109 Holden, private communication). The preliminary model and best slip model were computed
110 from the same dataset of near-source strong motion data. The preliminary models are based

111 on a single fault inversion, while the best models invert for kinematic parameters for three
112 newly-defined fault planes. These models provide more details about the overall rupture
113 process and better overall waveform fits (see Fig. S1 in the electronic supplement).

114 **Evaluation Metrics**

115 We evaluated the model forecasts with several CSEP methods (*Rhoades et al.*, 2011; *Schor-*
116 *lemmer et al.*, 2007; *Zechar et al.*, 2010; *Werner et al.*, 2011), which test for consistency
117 of the observations with the probabilistic forecasts and compare the predictive skills of the
118 models. We focus here on the comparison of the forecast skills and on a qualitative con-
119 sistency check between the numbers of observed and forecast earthquakes. The electronic
120 supplement contains remaining results.

121 We measure the skill of forecasts with the information gain per earthquake, which com-
122 pares a model’s predictive skill against a benchmark (*Rhoades et al.*, 2011). The benchmark
123 is a time-independent and spatially-uniform Poisson (SUP) process with a Gutenberg-Richter
124 magnitude distribution ($b = 1$). The SUP model is updated at each time step, so that the
125 total forecast rate for the next time step matches the average rate over the past catalog
126 in the test region. The information gain per earthquake calculates the average difference,
127 per earthquake, of the log-likelihood scores of a model and the benchmark. We use 95%
128 confidence bounds estimated by *Rhoades et al.* (2011) to assess statistical significance.

129 **Models**

130 Modelers submitted a total of sixteen models as software to the CSEP testing center, which
131 generated and evaluated forecasts. Due to a bug in STEP-cff, the first 1-day forecast was
132 produced offline. Here, we focus on the results of eight representative models (table 2),
133 described next.

134 **Non-parametric kernel smoothing models: K2, K3**

135 Models K2 and K3 represent statistical end-members on the spectrum of competing models.
136 None of the usual assumptions about earthquake clustering are explicitly included, such as
137 the Omori law or the Utsu-Seki clustering law (large earthquakes generate exponentially
138 more aftershocks). Instead, the models employ Gaussian kernels to estimate seismicity as
139 a function of time, space and magnitude (*Helmstetter and Werner, 2014*). K2 does assume
140 a Gutenberg-Richter magnitude distribution with $b = 1$, while K3 uses kernels to estimate
141 the (space-time dependent) magnitude distribution. The widths of the kernels adapt to the
142 activity level: sparse seismicity (in space and time) widens kernels; concentrated seismicity
143 narrows kernels. The models thereby adjust to the current seismicity rate, which is ex-
144 trapolated over the forecast horizon. These non-parametric kernel models offer maximum
145 flexibility, at the cost of dispensing with commonly observed empirical laws. Further details
146 can be found in the Supplementary Material.

147 **ETAS implementations: ETAS, ETAS-fault, ETAS-cff**

148 The empirical ETAS model and its hybrid model versions (ETAS-fault, ETAS-cff) are imple-
149 mented in an identical framework for the setup and parameter estimation which is explained
150 in detail in the Supplementary Material. Any difference in the performance is therefore
151 directly related to the ignorance or use of additional source information and stress calcu-
152 lations. In particular, the only difference is the spatial triggering kernel which is in the
153 case of the ETAS model one or a sum of isotropic power-law kernels centered at the loca-
154 tion of the preceding events. In contrast, for events with available slip models, ETAS-fault
155 uses an anisotropic power-law kernel as a function of the nearest distance to the mainshock
156 fault plane and ETAS-cff uses a probability distribution based on calculated Coulomb stress
157 changes (*Bach and Hainzl, 2012*).

158 **STEP and STEP-cff**

159 The hybrid model proposed by *Steacy et al.* (2014) is based on the STEP (Short-Term
160 Earthquake Probability) model, a purely statistical approach proposed by *Gerstenberger*
161 *et al.* (2005). STEP is a weighted sum of models with increasing spatial complexity, includ-
162 ing background seismicity and Omori decay. In addition, STEP-cff redistributes seismicity
163 according to the sign of the Coulomb stress change: 93% and 7% of events in regions of
164 positive and negative stress changes respectively. More details can be found in the electronic
165 supplement.

166 **Coulomb-rate-state models**

167 We focus here on two of the submitted Coulomb/rate-state (CRS) models, with the following
168 features:

- 169 • CRS-oop: uses Coulomb stresses imparted by the mainshocks ($M \geq 5.95$) on planes
170 optimally oriented with respect to the total Coulomb stress (optimally oriented planes,
171 OOPs).
- 172 • CRS-unc: in addition to mainshocks, this model includes stress changes from smaller
173 earthquakes. Instead of using OOPs, CRS-unc accounts for the variability of receiver
174 fault orientations by resolving stress changes on a set of faults from the regional focal
175 mechanisms catalog (electronic supplement and *Cattania et al.* (2014))

176 For both model versions, seismicity rates are calculated by considering the response
177 of a population of faults with rate-and-state dependent friction (*Dieterich*, 1994), where
178 parameter setting and estimations are done in an identical manner. Both models use an
179 internal grid with a higher resolution than the output grid. CRS-oop is similar to earlier
180 implementations (*Woessner et al.*, 2011), except for the use of an internally refined grid,
181 while the additional features in CRS-unc are new.

182 Results

183 Temporal performance

184 Most models successfully forecast the total number of events and the main features of day-
185 by-day evolution (Fig. 3a). All ETAS models, the STEP models and CRS-unc forecast the
186 total event number to within (Poissonian) uncertainty. Models K2 and K3 underpredict by a
187 factor of about two, while CRS-oop strongly overpredicts. K2 and K3 heavily underestimate
188 the number of triggered earthquakes during the first day of each sequence, but they otherwise
189 forecast the rates well. Since these models do not include mainshock magnitude, but estimate
190 aftershock number from the observed seismicity, starting the forecast exactly at the time of
191 each mainshock (before aftershocks have occurred) hinders their performance on the first
192 day. This is a weakness in the present experimental design.

193 All models underestimate the number of events triggered by the three large quakes that
194 followed Darfield. The STEP models and, more so, the CRS models, forecast a slower decay
195 after the large shocks than is observed. We note that both CRS models tend to select high
196 values of the aftershock duration t_a (close to 27 yrs, the upper end of the parameter search
197 range): lower t_a values, which would give a faster decay, give a worse fit during the inversion
198 period and are not selected. CRS-oop severely overestimates the number of shocks after the
199 Darfield earthquake, and it does not predict any aftershocks of the last mainshock: this is
200 because the model only considers stress sources from events above a user-defined minimum
201 magnitude, which was set to a value 5.95 (greater than the last large shock's magnitude
202 $M5.9$). The use of a predefined "mainshock" magnitude has been eliminated in later versions
203 of the code (*Cattania and Khalid, 2016*).

204 The ETAS models match the temporal evolution most closely. They forecast identical
205 numbers because they differ only in their spatial densities. Their success is a result of an
206 Omori p-value $p > 1$ (one of the models' free parameters) and a rather high α value, which
207 reduces the relative importance of secondary triggering by smaller events.

208 Spatio-temporal performance

209 The largest differences between model forecasts occur during the first day after each main-
210 shock (e.g. Fig. 4). The expression of the Coulomb component appears remarkably different
211 between CRS-unc and the hybrid models ETAS-cff and STEP-cff, illustrating the sensitivity
212 of these forecasts to the specific implementation of the hypothesis. ETAS-cff and STEP-cff
213 display the more commonly expected Coulomb lobes of a predominantly strike-slip earth-
214 quake, while CRS-unc displays much smoother lobes. For ETAS-cff, seismicity rates are
215 linearly related to stress changes; STEP-cff considers only the sign of the stress change, and
216 hence it presents sharp transitions along the nodes, not seen in ETAS-cff (Fig. 4); CRS
217 models are strongly nonlinear in stress, due to the rate-state equations. Moreover, the dif-
218 ferent treatment of uncertainties (such as receiver faults and subgrid variability) introduces
219 additional differences. The overall pattern for model CRS-oop (not shown) is similar to
220 CRS-unc. The four Coulomb-based models forecast the first day of seismicity much more
221 successfully along the Darfield rupture than the three statistical models; ETAS-fault model
222 forecasts the seismicity about as well, and indeed better after the first few days (Fig. 3).

223 As already discussed, K2 and K3 do not use the mainshock magnitude to forecast after-
224 shocks and therefore forecast very low seismicity on the first day (Fig. 4). On the second
225 day, however, they forecast a spatial pattern similar to the ETAS model and consistent with
226 observations. This highlights the ability of these models to adapt quickly once enough quakes
227 have occurred (about 10 events).

228 ETAS-based models are the most successful at reproducing the spatial distribution of
229 seismicity with distance from the fault integrated over the entire time period (Fig. 5). Mod-
230 els K2 and K3 underestimate seismicity, but they have an overall trend similar to the catalog,
231 with most seismicity within cells centered at 0.5-10km from the mainshock fault. (Fig. 3).
232 Both CRS-models underestimate seismicity rates within the first few kilometers from the
233 fault, and overestimate rates beyond 5 km from the fault. ETAS-cff also tends to overesti-
234 mate rates beyond 10 km from the mainshock faults, while ETAS and ETAS-fault predict a
235 faster spatial decay. In contrast, the difference between STEP and STEP-cff is minimal.

237 **Model ranking**

238 The STEP models and the hybrid models ETAS-cff and ETAS-fault generated the most
 239 informative forecasts across all three (1-day, 1-month, 1-year) forecast horizons (see Fig. 6,
 240 best-available data scenario). CRS-unc and CRS-oop performed slightly less well, but they
 241 were quite close to the hybrid models, and better than the simple ETAS model over longer
 242 forecast horizons. Nonetheless, the Coulomb component as implemented in STEP-cff affected
 243 its performance very slightly (lowering the information gain), and the ETAS-cff model did
 244 not provide additional skill over the ETAS-fault model. CRS-oop consistently performed
 245 slightly worse than CRS-unc, to some extent because CRS-oop did not use the last large
 246 shock as a stress source (see Fig. 3b).

247 K2 and K3 presented the lowest information gains, because they performed poorly dur-
 248 ing the first time window after each mainshock (Fig. 4). Because the log-likelihood score
 249 is dominated by earthquake occurrences rather than empty bins, the slower-than-observed
 250 decay predicted by most models did not affect their ranking significantly.

251

252 Most models (except the STEP models, and the 1-month forecasts of K2 and K3) per-
 253 formed better when they were provided the best-available input data, due to either a more
 254 complete and accurate catalog (K2, K3 and ETAS) and also to better slip models (CRS and
 255 hybrid models). We found that even the CRS models were more sensitive to the quality of
 256 the catalog than to the slip models (fig. S10). Models ETAS-cff and ETAS-fault performed
 257 identically to the simple ETAS model in the near-real-time data scenario: in the absence of
 258 preliminary slip models (not provided until day 10), these models reverted to simple ETAS
 259 models, and the first 10 days heavily dominated the information gain. For a few models, the
 260 difference in information gain with best and preliminary data is smaller than 95% confidence
 261 intervals (Fig. 6); and even with preliminary data, all models do significantly better than
 262 the SUP model, as previously observed for Japanese sequences (*Omi et al.*, 2016).

263 We can gain some insight into the model performance from the spatial distribution of
264 information gains (Fig. 7). Near the Darfield fault, ETAS-cff was the best performing model,
265 followed by the CRS models. ETAS-cff also better forecasted the few aftershocks to the
266 north-west of the Darfield earthquake, but it overpredicted in the remainder of this enhanced
267 Coulomb lobe (region (1)). ETAS performed worse than its hybrid counterparts, except for
268 the aforementioned lobe of ETAS-cff and a small region near the epicenter of the Darfield
269 earthquake (since its isotropic kernel leads concentrates the forecast for the first day in this
270 area; region (2) in Fig. 7). STEP and STEP-cff present small differences in information
271 gains, indicating that the Coulomb mask has only a subtle effect; this occurs because most
272 of the events forecasts by the STEP model already occur in regions where stress changes
273 resolved on OOPs are positive, so that the redistribution of events does not change the rates
274 significantly. We verified that few points of negative information gains for STEP-cff fall into
275 cells where STEP-cff calculated negative stress changes (region (4) in Fig. 7), near a node of
276 the stress field; in contrast, ETAS-cff does not present a stress shadow and performed better
277 than ETAS in the same cells. This can be due to two reasons: ETAS-cff considers stresses
278 at multiple depth layers, and resolves it on a set of receiver faults; and since STEP-cff only
279 considers the sign of the stress change, it overestimates its effects near nodes of the stress
280 field, where the absolute value is low.

281 While CRS-unc outperformed STEP along much of the Darfield fault, STEP better
282 captured the Christchurch and Pegasus Bay sequences. As noted above, CRS-oop did not
283 consider the Pegasus Bay $M5.9$ earthquake as a stress source, and it therefore predicted no
284 aftershocks (Fig. 3a). Both CRS models did poorly at intermediate distances ($\gtrsim 10km$ from
285 the mainshock faults), where higher seismicity rates were forecasted than observed. The
286 good performance of the CRS models along the Darfield fault may seem surprising since
287 the CRS models predicted lower near-fault rates than others (Fig. 5): this occurred because
288 the log-likelihood is space and time dependent, and CRS models predicted higher seismicity
289 rates than others on the first day of the forecast, when about a third of the aftershocks took
290 place.

Discussion and Conclusions

The ranking of the models indicates that including physical information, such as fault geometry or Coulomb stress changes, can lead to better overall model performance. This is particularly clear from the comparison ETAS to ETAS-cff and ETAS-fault, in agreement with a retrospective case study for California mainshocks (*Bach and Hainzl, 2012*).

Coulomb rate-state models, and in particular CRS-unc, have a performance comparable to the hybrid models, in stark contrast to a previous retrospective evaluation (*Woessner et al., 2011*), and in agreement with a comparative study of seismicity in Northern California (*Segou et al., 2013*). This result indicates that when resolving Coulomb stresses in more detail, by using an internally refined grid and including uncertainties and secondary stress sources, the overall performance of physics-based models greatly improves. On the other hand, their spatial and temporal fit indicate that some aspects of the triggering mechanism are not yet captured, as discussed below.

Most of the ETAS-based models (ETAS, ETAS-fault and STEP) prescribe a functional form for the spatial decay of seismicity from the mainshock sources (a power law), and they reproduce the observed decay reasonably well. The inclusion of Coulomb stress changes in ETAS-cff leads to overestimation of off-fault seismicity, analogous to the CRS models. STEP-cff, on the other hand, only considers the sign of the stress change and therefore preserves the power-law decay prescribed by STEP, so that the two models exhibit a similar decay (Fig. 5). Models K2 and K3, which estimate the spatial distribution directly from the catalog itself, also provide a good fit when accounting for the fact that rates are underestimated everywhere due to the lack of information on the first day.

A major simplification in the CRS models was to assume spatially uniform background rate. With this assumption, the model does not distinguish between areas with fault structures capable of hosting seismicity, and areas without pre-existing faults, leading to overestimation in the far-field. Moreover, in the rate-state formulation, weakly stressed regions contribute to seismicity later in the sequence (*Dieterich, 1994; Helmstetter and Shaw, 2006*),

319 so that assuming a uniform background rate leads to a slower decay (*Cattania et al.*, 2015),
320 consistent with Fig. 3. One of the models submitted for testing (electronic supplement) is
321 a variation of CRS-unc, including heterogeneous background rate derived from smoothed
322 seismicity (*Helmstetter et al.*, 2007). However, this model has a poorer performance than
323 CRS-unc, because before the Darfield earthquake the seismic activity was dominated by the
324 Alpine fault system, with relative little seismicity in the area of the Darfield-Canterbury
325 sequence (see model K2 in Fig. 4, first day). Estimating the spatially variable background
326 seismicity rate, especially when the mainshock hits relatively quiescent regions, remains one
327 of the challenges of Coulomb rate-state models (e.g. *Bhloscaidh et al.*, 2014; *Cocco et al.*,
328 2010). Another challenging aspect in modeling Coulomb stress triggering is the heterogene-
329 ity in stress, especially in the near field. In addition to the variable orientation of receiver
330 faults, a source of stress heterogeneity is the small scale variability of seismic slip, gener-
331 ating locally high stresses on the fault plane and seismicity within the rupture area, where
332 the average stress is negative (*Helmstetter et al.*, 2007). We note that considering multiple
333 fault orientations has a similar effect in terms of stress shadow reduction (*Cattania et al.*,
334 2014), leading to reasonable information gains even near the mainshock faults (Fig. 6); how-
335 ever, underestimation of near-field stresses may contribute to the overall underestimation
336 of seismicity in these regions (Fig. 5).

337 The better performance of CRS-unc over CRS-oop is consistent with previous stud-
338 ies (*Cattania et al.*, 2014, 2015), and was due to the inclusion of secondary triggering and
339 uncertainties due to receiver fault orientation (for a comparison with models including only
340 one of these aspects, see electronic supplement). The use of OOPs instead of a fixed re-
341 ceiver fault typically leads to a better performance in the near field and short time after the
342 mainshock (e.g. *Hainzl et al.*, 2009; *Woessner et al.*, 2011; *Segou et al.*, 2013), since they can
343 reproduce high rates near the mainshock. Here we find that model CRS-unc has a better
344 performance than CRS-oop across all temporal and spatial scales (Fig. 3, 5). This result
345 suggests that using known information on the local fault geometry (from focal planes, as
346 done here; or from mapped faults, when available) may be the optimal forecasting strategy,

347 as long as the variability of fault orientations is also modeled.

348 We note that CRS-oop has better scores, relative to statistical models, than a similar
349 model tested in the Landers retrospective experiment (*Woessner et al.*, 2011). This is most
350 likely due to the use of a refined grid for Coulomb stress calculations. Considering multiple
351 depth layers, for example, accounts for the fact that stress changes (resolved on the mainshock
352 fault plane) are negative within the rupture area and positive above and below it; therefore,
353 calculating stresses at a single intermediate depth will likely result in underestimation of
354 on-fault rates. Since we did not test a CRS model without grid refinement, we can not
355 directly measure the improvement due to this aspect. An indirect test, however, comes from
356 the study by *Steacy et al.* (2014), who compared STEP, STEP-cff and a classic CRS model
357 (using OOPs, and no grid refinement) for the Canterbury sequence. While the different time
358 window and target magnitude prevent us from comparing information gains exactly, we note
359 that the overall performance of the CRS model was significantly lower than the STEP model,
360 with a difference in information gains per event of about 2 – 3. The relative performance
361 between STEP and STEP-cff, on the other hand, is close to what we find here: a small
362 difference in information gain per event (< 0.1), with the STEP model performing slightly
363 better.

364 ETAS-cff shares certain aspects of model implementation with CRS-unc: vertical grid
365 refinement and consideration of receiver fault variability (even though the set of receiver
366 faults was different; see Electronic Supplement). The improvement of ETAS over ETAS-cff,
367 in contrast with models STEP-cff and STEP, confirms that these aspects have a first-order
368 effect on information gains.

369

370 The relatively good performance of CRS models is encouraging in terms of our physical
371 understanding of earthquake triggering. Like earlier versions (*Woessner et al.*, 2011), these
372 models are based on two widely accepted concepts: that aftershocks are mainly caused by
373 static stress changes, and that time-dependence of their nucleation is controlled by rate-state
374 friction (*Dieterich*, 1994). The drawback of physical models is that several of the quantities

375 involved in Coulomb stress calculations and rate-state seismicity evolution are not known
376 precisely. The improvement in performance compared to earlier studies suggests that the
377 main issues with physics-based models was not in the fundamental process but rather in
378 specific details of model implementation.

379

380 There are still multiple ways in how physical models can be refined: for example, we
381 identified the spatial dependence of background seismicity as a particularly challenging as-
382 pect. Other improvements such as the inclusion of aseismic stresses or consideration of the
383 spatial variability in receiver fault orientations can in the future be tested in the context
384 of CSEP. Another important question to address is the practical use of these models in the
385 context of operational earthquake forecasting: we note that currently, hybrid models with a
386 similar performance require less computation time, making them more suitable for real-time
387 applications.

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396 **Data and Resources**

397 The GeoNet catalog and the focal mechanism catalog can be accessed at: <https://www.geonet.org.nz/>

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Table 1: Overview of data sets.

	data type	use	features
Geonet-Cat	catalog	target data/model input (mode 1/2)	$M \geq 3.95$
Preliminary-Cat	catalog	model input (mode 3)	Captured by NZ CSEP Testing Center Beaven et al. 2012
Best-Slip	slip models	model input (mode 1/2)	
Preliminary-Slip	slip models	model input (mode 3)	Holden et al. 2011 & personal comm. GeoNet
Focal Mechanisms	focal mechanisms	model input (mode 1-3)	

Table 2: Overview of models participating in the test with reference whether or not they use the Gutenberg-Richter distribution (GR), the Omori-Utsu decay function (OU), an exponential productivity function ($N(M)$), fault information (fault), Coulomb Failure Stress (CFS), rate- and state-dependent frictional response (RS), or focal mechanisms (FM).

index	model name	empirical relations			fault information & physics				reference
		GR	OU	$N(M)$	Fault	CFS	RS	FM	
0	SUP	x							Rhoades et al, 2018, this issue
1	K2	x							<i>Helmstetter and Werner (2014)</i>
2	K3								<i>Helmstetter and Werner (2014)</i>
3	ETAS	x	x	x					<i>Ogata (1988)</i>
4	ETAS-fault	x	x	x	x				<i>Bach and Hainzl (2012)</i>
5	ETAS-cff	x	x	x		x			<i>Bach and Hainzl (2012)</i>
6	STEP	x	x	x	x				<i>Gerstenberger et al. (2005)</i>
7	STEP-cff	x	x	x	x	x			<i>Steacy et al. (2014)</i>
8	CRS-oop	x				x	x		<i>Cattania et al. (2014)</i>
9	CRS-unc	x				x	x	x	<i>Cattania et al. (2014)</i>

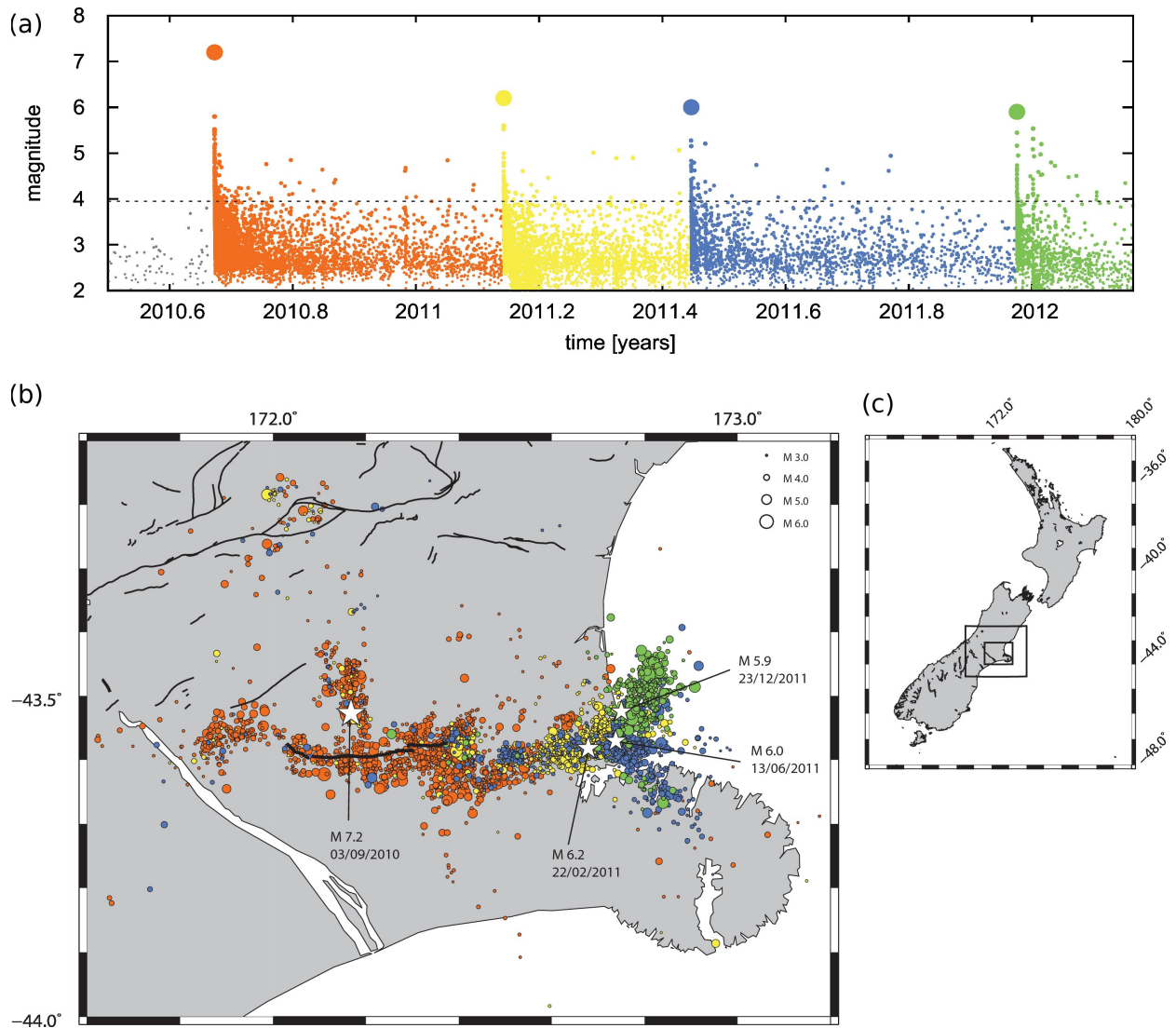


Figure 1: (a) Magnitude vs. time from the reviewed GeoNet catalog. The different colors indicate the sequences of events with $M \geq 5.9$. (b) Map view of the seismic sequence, colorcoded in agreement with the top panel. Fault lines are from the New Zealand Active Fault Database, and the thicker line is the Greendale fault. (c) map of new Zealand, with the boxes marking the forecast area (larger box) and the map on the left.

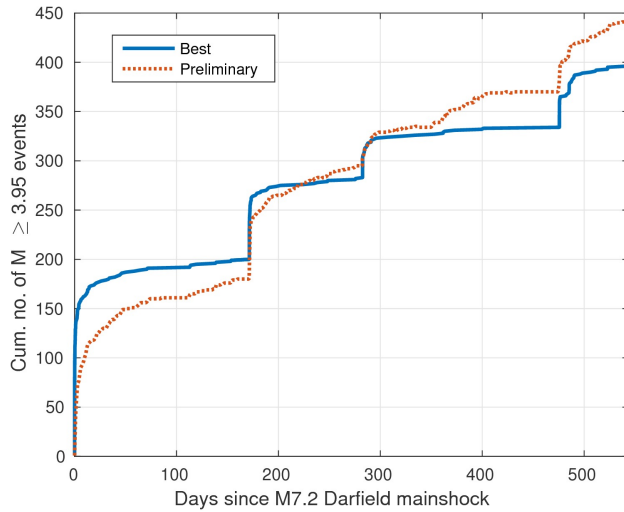


Figure 2: Cumulative number of $M \geq 3.95$ events in the reviewed reviewed catalog and real time data. For the real time data, we report the total number of events in the catalog used on each day: as the catalog is revised, the number of events may vary because the catalog becomes more complete, or because magnitudes are revised. Magnitude were initially reported as M_L , and later replaced by M_w ; since M_L were systematically overestimated until the end of 2011, the number of events in the real time catalog can exceed the reviewed catalog.

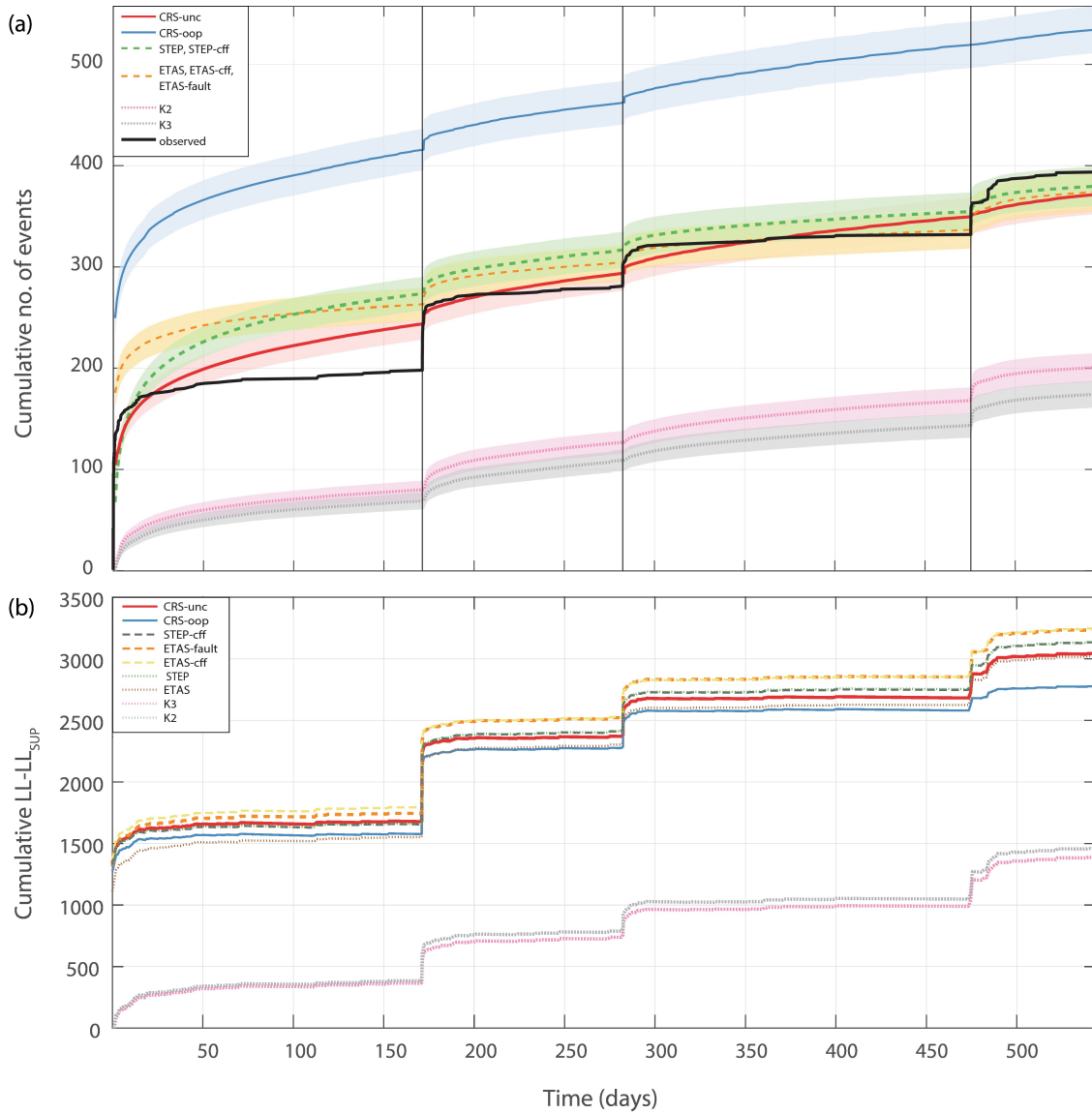


Figure 3: (a) Forecasted and observed temporal evolution. Shaded areas indicate Poissonian errors; vertical lines are events with $M \geq 5.9$. (b) Cumulative difference in log-likelihood with respect to the SUP model, obtained from the sum of the log-likelihood calculated for (space, time, magnitude) bins.

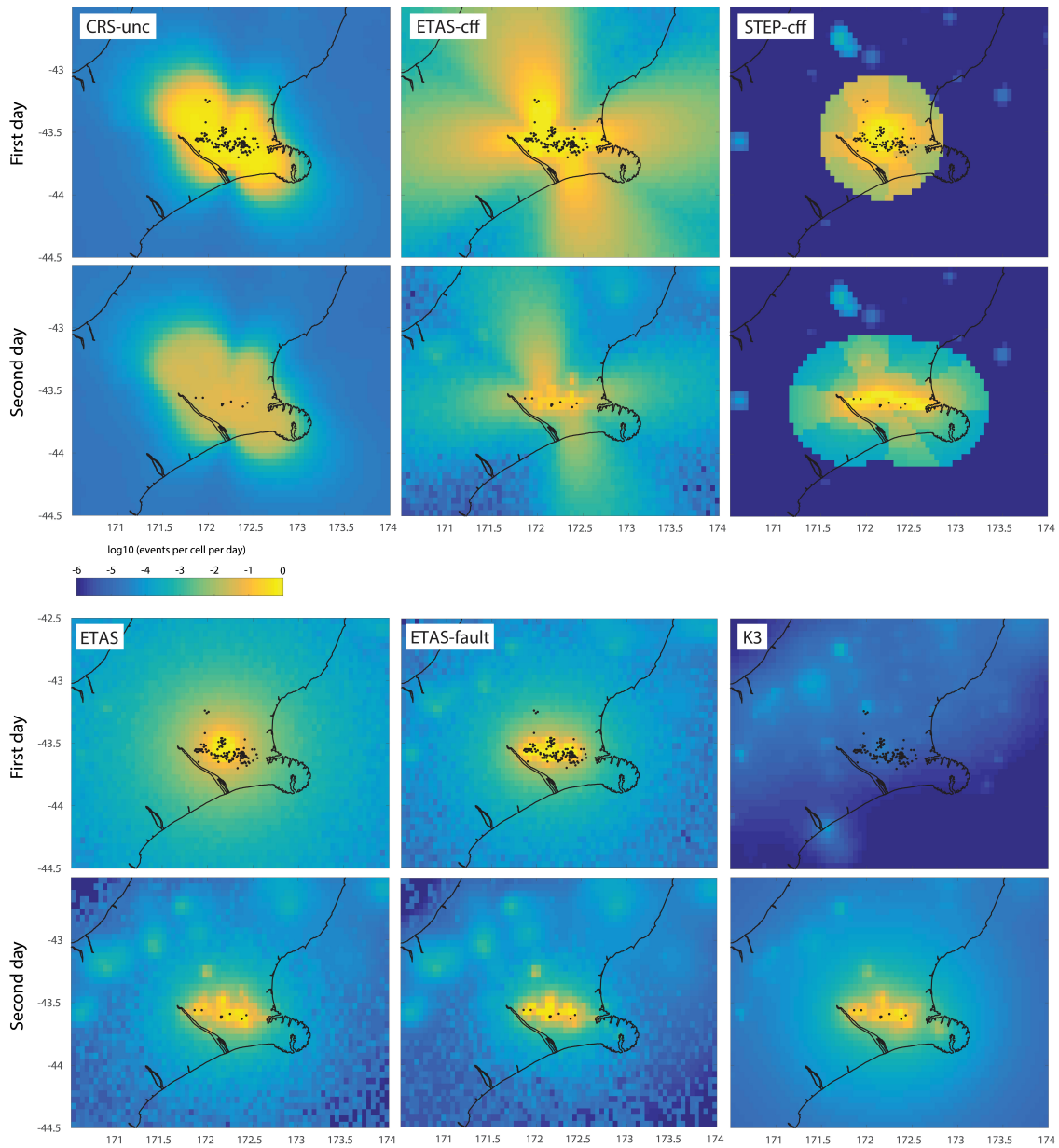


Figure 4: Examples of 1-day forecasts for selected models. The top row is the forecast starting at the time of the Darfield mainshock, and the second line forecasts are the second day: the difference highlights how each model incorporates information from the early part of the sequence. The dots are the observed events in each time period.

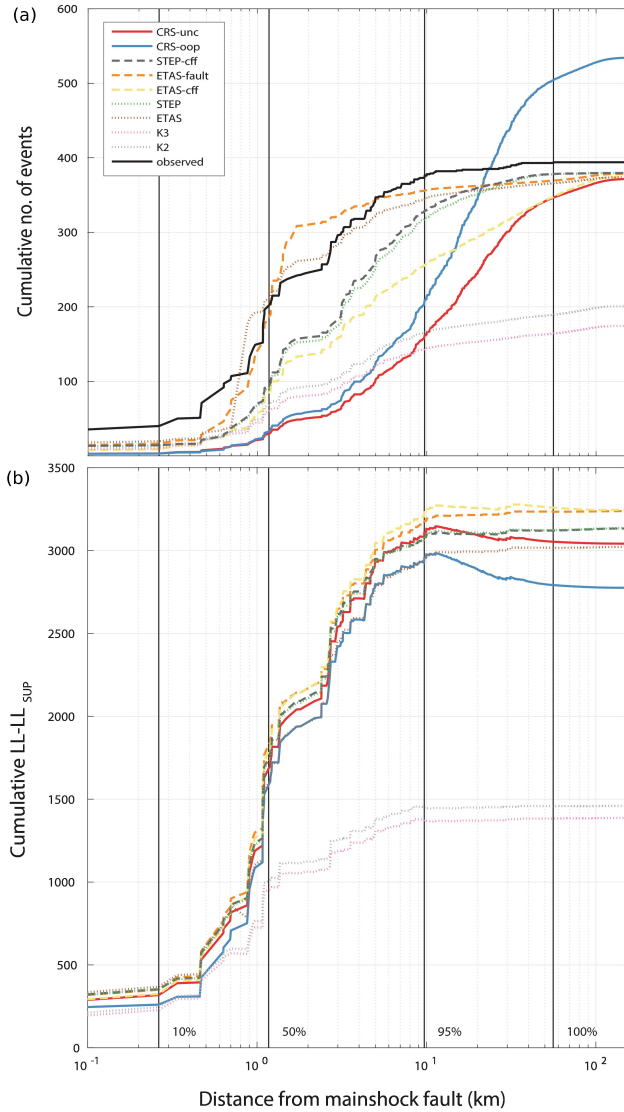


Figure 5: (a) Cumulative number of earthquakes as a function of distance from the nearest mainshocks fault trace, based on the location of the cell center. For consistency, the catalog was also binned into the forecast cells, so that distances does not reflect exact earthquake locations but rather the cell to which they are assigned. (b) cumulative difference in log-likelihood from the SUP forecast, obtained from the sum of the log-likelihood calculated for (space, time, magnitude) bins. The vertical lines indicate the percentage of earthquakes within cells at a given distance from the faults.

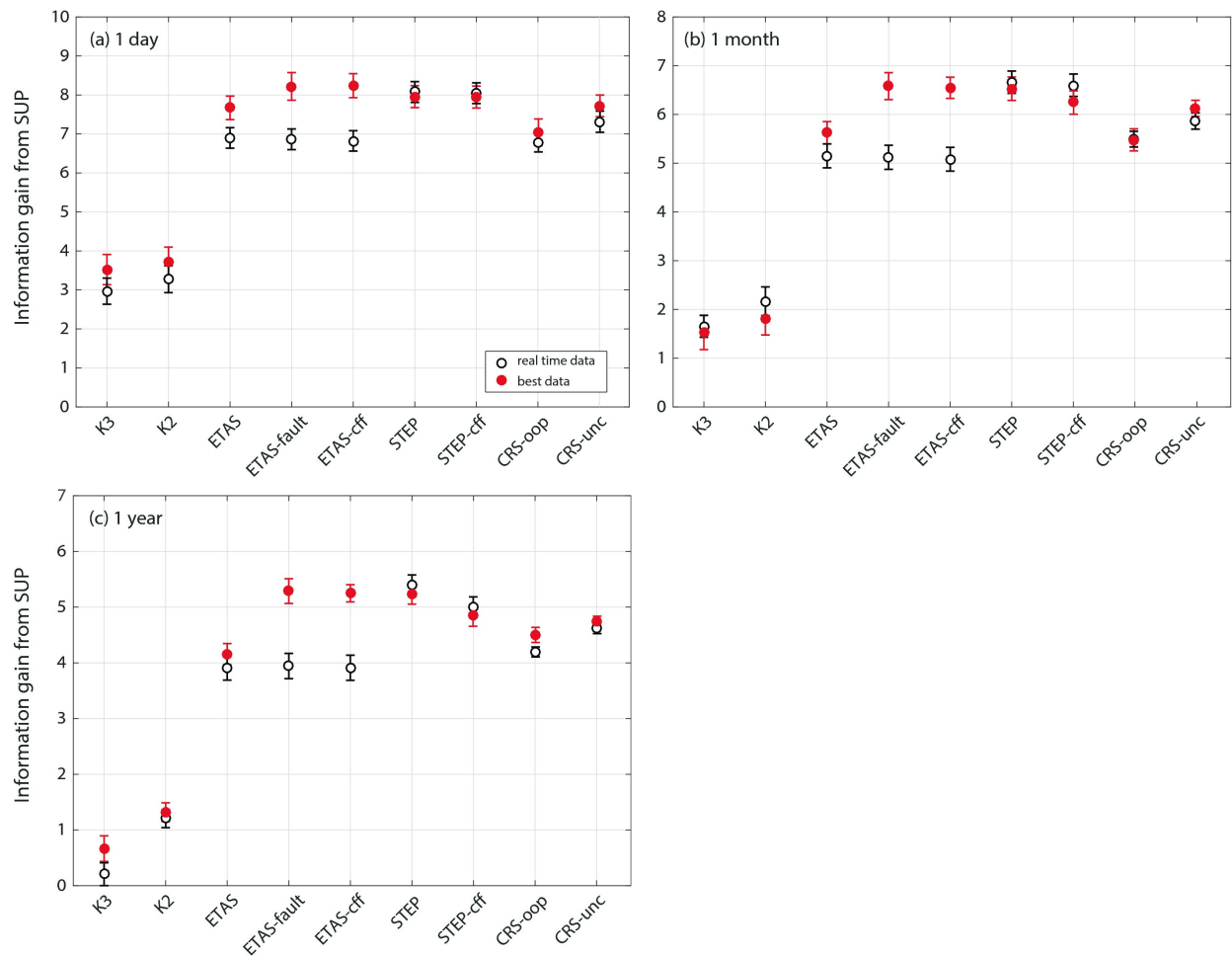


Figure 6: Information gain per earthquake relative to the SUP model, for real time and best available data. Each panel shows a forecast horizon. Error bars represent 95% confidence levels from a paired student-T test (*Rhoades et al., 2011*).

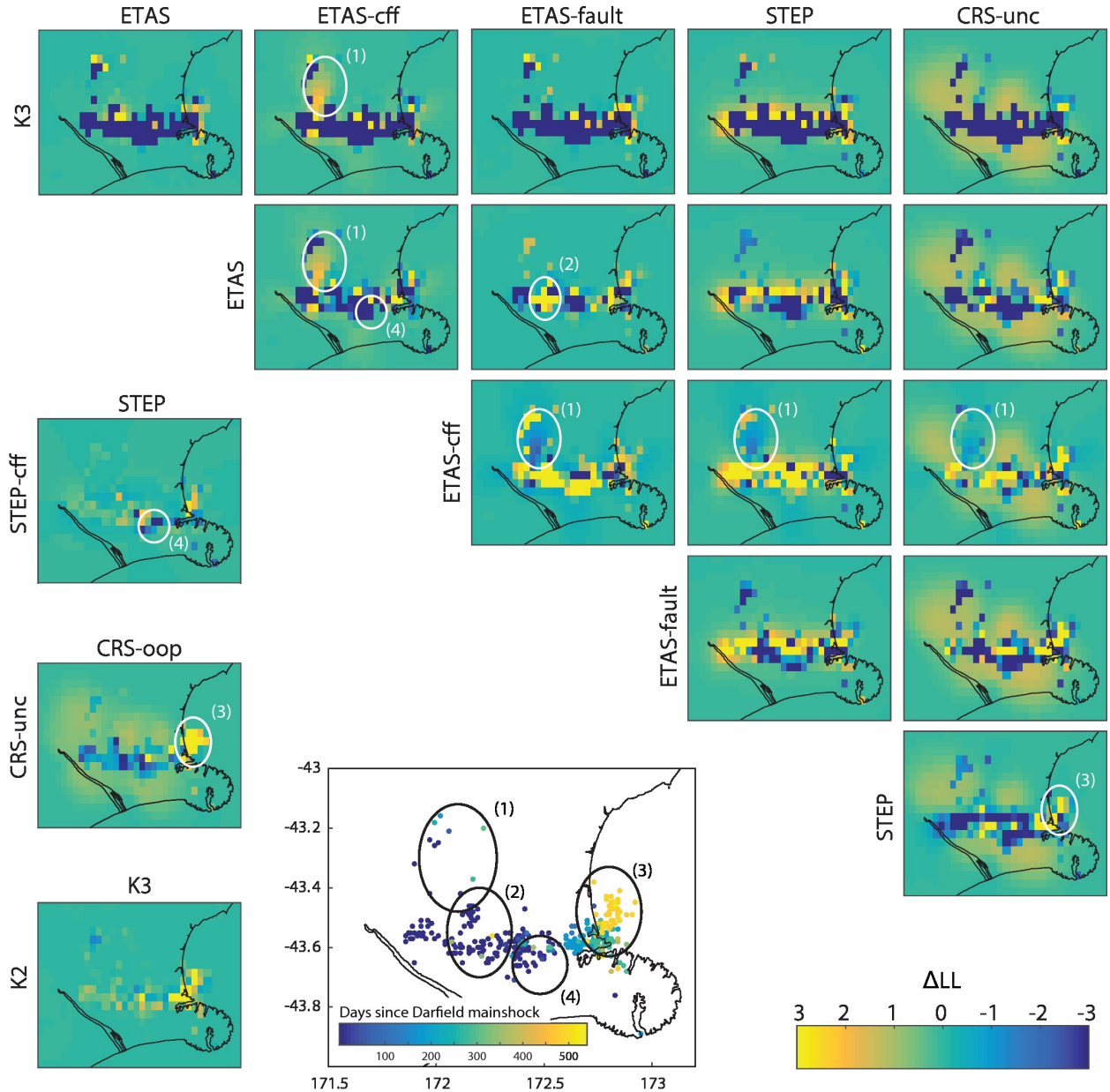


Figure 7: Map of log-likelihood differences between pairs of models, for a subset of the model domain near the mainshock faults. The color indicates $\Sigma_n LL_{n,i} - \Sigma_n LL_{n,j}$, where i and j are the row and column index, and the summation is performed over all time steps and magnitude bins. Note that the values are capped at ± 3 for clarity. Positive values along a row indicate good model performance, and along a column they indicate poor performance. The ellipses mark the following features, discussed in the main text: (1) a Coulomb stress lobe; (2) the area near the $M7.2$ Darfield epicenter; (3) the aftershocks of December 23th $M5.9$ Pegasus Bay earthquake; (4) few cells near a node of the Coulomb stress field, where STEP-cff predicts a stress shadow.