

Originally published as:

Diao, F., Wang, R., Wang, Y., Xiong, X., Walter, T. R. (2018): Fault behavior and lower crustal rheology inferred from the first seven years of postseismic GPS data after the 2008 Wenchuan earthquake. - *Earth and Planetary Science Letters*, 495, pp. 202—212.

DOI: http://doi.org/10.1016/j.epsl.2018.05.020

Manuscript to be submitted to EPSL

1	Fault behaviour and lower crustal rheology inferred from the first seven years of
2	postseismic GPS data after the 2008 Wenchuan earthquake
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12 Abstract

Long-term and wide-area geodetic observations may allow identifying distinct postseismic 13 deformation processes following large earthquakes, and thus can reveal fault behaviour 14 and permit quantifying complexities in lithospheric rheology. In this paper, the first 7 15 years of GPS (Global Positioning System) displacement data following the 2008 Mw7.9 16 Wenchuan earthquake are used to study the relevant mechanisms of postseismic 17 deformation. Two simple models that consider either afterslip or viscoelastic relaxation as 18 the unitary mechanism of the postseismic deformation are tested at first. After analysing 19 the limitations and complementarity of these two separated models, a combined model 20 incorporating the two main mechanisms is presented. In contrast to previous studies, 21 which mostly assume that afterslip and viscoelastic relaxation are independent, our 22 combined model includes the secondary viscoelastic relaxation effect induced by transient 23

24	afterslip. Modelling results suggest that the middle- to far-field postseismic deformation is
25	mainly induced by viscoelastic relaxation of the coseismic stress change in the lower
26	crust, whereas the near-field displacement is dominantly caused by stress-driven aseismic
27	afterslip. The seismic moment released by the transient afterslip corresponds to an Mw7.4
28	earthquake, or 25% of that released by the Mw7.9 main shock. With a characteristic decay
29	time of 1.2 years, most afterslip (~ 80%) is released in the first 2 years. Negligible
30	aseismic afterslip is observed in the seismic gap between the Wenchuan and Lushan
31	earthquakes, indicating the locked state of the fault within this segment. The effective
32	lower crustal viscosity of the eastern Tibetan Plateau is estimated to be 2.0×10^{18} Pa·s, at
33	least two orders of magnitude smaller than that of the adjacent Sichuan Basin. This
34	finding is consistent with previous observations of low seismic velocity and high electrical
35	conductivity in this region, all of which support the assumption that the crustal thickening
36	in the eastern Tibetan Plateau is dominantly caused by ductile lower crustal flow, with
37	important implications for understanding both long- and short-term crustal deformation
38	processes.

Keywords: Wenchuan earthquake, afterslip, viscoelastic relaxation, postseismic
 deformation, seismic hazard

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42 **1. Introduction**

Postseismic deformation following large earthquakes has been routinely observed and
explained as the response of the lithosphere to coseismic stress redistribution (Bürgmann and
Dresen, 2008). The main mechanisms proposed for postseismic deformation include
viscoelastic relaxation of coseismically induced stress changes in the ductile lower crust and
upper mantle (e.g., Bürgmann and Dresen, 2008; Wang et al. et al., 2012; Freed et al., 2017)

and aseismic fault afterslip within or near the coseismic rupture area (e.g., Hsu et al., 2006; 48 49 Freed, 2007). In addition, the poroelastic rebound effect can also cause observable postseismic deformation, but this mainly occurs in the near-field area with a relatively lower 50 magnitude (Jónsson et al., 2003), which is often ignored due to a lack of observations affected 51 by it. In most cases, aseismic afterslip and viscoelastic relaxation are widely used to interpret 52 postseismic observations. However, due to a similar deformation pattern that is induced by 53 54 these two mechanisms, it remains challenging to distinguish their relative contributions, especially if the observations have an insufficient temporal and spatial coverage (Bürgmann 55 and Dresen, 2008). 56

Aseismic afterslip releases stress in weak regions on faults adjacent to coseismic slip 57 58 asperities (Hsu et al., 2007), potentially reflecting the stress state of the faults. In many 59 previous studies, static or time-dependent afterslip is usually directly derived from surface postseismic displacements via a linear inversion. This approach is convenient and most 60 effective for explaining the observations, but it is not physically based and requires a large 61 number of free parameters. In addition, by ignoring the contribution of viscoelastic relaxation, 62 the obtained afterslip is often overestimated in terms of magnitude and depth (Diao et al., 63 2014). From recent decades of studies on postseismic deformation, afterslip has been regarded 64 as the dominant mechanism governing the near-field in the initial stage following the 65 66 earthquake. However, a study on the 2011 Mw9.0 Tohoku earthquake indicates that this paradigm is no longer valid because viscoelastic relaxation can contribute as much as afterslip 67 to postseismic displacements in the first 3.5 years (Freed et al., 2017), suggesting that relative 68 contributions of afterslip and viscoelastic relaxation on postseismic deformation depend on 69 the fault behaviour and regional rheology. 70

Rheologically weak regions within the lower crust and upper mantle cannot sustain the
 coseismic stress change and release it in the form of viscoelastic relaxation, which can result

in significant crustal deformation (Bürgmann and Dresen, 2008). With measurements of 73 74 postseismic deformation, the viscosity in regions beneath the coseismic slip area can be explored, which can provide an important *a-priori* information for dynamic studies. Given the 75 nonlinear relation to postseismic displacements, the optimal viscosity value that best fits the 76 data is generally solved using a grid search approach among various forward simulations 77 (Diao et al., 2014; Freed et al., 2017; Huang et al., 2014). Viscoelastic relaxation can 78 generally predict the long-term deformation in the middle- to far-field, but this underestimates 79 observations in the near-field, which implies that other mechanisms might be involved. 80 Considering the limitations of separated afterslip and viscoelastic relaxation models, there has 81 82 been a consensus that a combined model incorporating these two main mechanisms is more realistic. 83

84 On 12 May 2008, the devastating Mw7.9 Wenchuan earthquake occurred on the eastern margin of the Tibetan Plateau adjacent to the Sichuan Basin (Fig. 1), which is a region 85 86 of high topographic gradient with relatively low strain accumulation rates (Zhang, 2010). Geologic field data, geodetic and seismic monitoring observations have offered valuable 87 constraints on the coseismic slip distribution of the earthquake (e.g., Xu et al., 2009; Shen et 88 al., 2009; Wang et al., 2011). As shown in most slip models, this earthquake first ruptured the 89 southwestern segment with an oblique thrusting slip. Then, the rupture propagated towards 90 the northeast with an increased right-lateral slip component, resulting in a total rupture length 91 of ~ 300 km (Fig. 1). The stress state of the medium surrounding the coseismic rupture was 92 severely changed by the 2008 event and significant postseismic deformation occurred after 93 the earthquake (Shao et al., 2011; Shen et al., 2011; Huang et al., 2014; Xu et al., 2014; Jiang 94 et al., 2017). Observations related to this postseismic deformation are not only unique for 95 understanding the postseismic mechanisms of this earthquake, but also valuable for probing 96 the regional rheology and fault behaviour in the eastern region of the Tibetan Plateau. 97

The postseismic deformation of the Wenchuan earthquake has been analysed in 98 99 several previous studies (Huang et al., 2014; Jiang et al., 2017; Shao et al., 2011; Shen et al., 2011; Xu et al., 2014). As summarized in Table 1, all studies except that by Xu et al. (2014) 100 used combined models that incorporate afterslip and viscoelastic relaxation to explain the 101 postseismic displacements. However, the parameters of these combined models, i.e., the lower 102 crustal viscosity and the afterslip distribution, were not determined simultaneously but 103 successively. In these previous studies, a viscoelastic relaxation model is optimized first to fit 104 the accumulated postseismic displacement data without considering the contribution from the 105 afterslip. In the second step, the residual data obtained by subtracting the viscoelastic 106 107 relaxation effect are inverted for a static afterslip distribution. The time dependence of the 108 afterslip and the secondary effect of viscoelastic relaxation induced by the afterslip are generally neglected. Moreover, previous studies may suffer from the limited spatial and 109 110 temporal coverage of the observations and the simplified earth models (Table 1).

111 Here, we present a set of well-distributed postseismic GPS data, which covers the first 7 years following the Wenchuan earthquake. We initially carry out a separate set of 3D 112 viscoelastic relaxation simulations and kinematic afterslip inversions to identify how and to 113 what extent the observations can be explained by each of the two main mechanisms. Then, we 114 demonstrate a combined model incorporating viscoelastic relaxation (due to both coseismic 115 116 slip and afterslip) and stress-driven afterslip, and we estimate their relative contributions to the observed postseismic deformation. Finally, we discuss the lower crustal viscosity of the 117 Tibetan Plateau and seismic hazards in the southwestern end of the Wenchuan rupture based 118 119 on the results of this study.

120 **2. GPS Data Processing**

Campaign GPS data, which are used to evaluate the postseismic deformation of the
Wenchuan earthquake, were collected from 1999 to 2015 by the Crustal Movement

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Observation Network of China (CMONOC). Pre-earthquake observations are required to obtain the background tectonic deformation rate, which should be removed to isolate the postseismic transient signals. Therefore, only GPS stations with at least four pre-earthquake surveys that cover a period of ~ 10 years (Fig. S1) were selected for use in this study.

We processed the GPS data following the approach presented in Gan et al. (2007). The daily raw GPS data were first processed to obtain loosely constrained daily coordinates and satellite orbits using the GAMIT software (Herring et al., 2010a). Then, the obtained, loosely constrained daily solutions were combined with loosely constrained global solutions of 80 IGS tracking stations (released by the Scripps Orbital and Position Analysis Center, http://sopac.ucsd.edu/) using the GLOBK software (Herring et al., 2010b), which fix the solutions to the ITRF08 reference frame (Gan et al., 2007).

After the position time series have been achieved, a station-dependent exponential function is employed to describe the time-dependent behaviour of the obtained postseismic displacement time series (Freed et al., 2017),

137
$$U(t) = U_0 + V_0 \cdot \left(t - t_{eq}\right) + \left\{ D_e + D_v \cdot \left[1 - e^{-\frac{(t - t_{eq})}{\tau}}\right] \right\} H\left(t - t_{eq}\right)$$
(1)

where t_{eq} (hereafter $t_{eq} = 0$) is the earthquake occurrence time, U_0 is the pre-event station 138 position, V_0 is the regional secular velocity of GPS stations, D_e is the coseismic elastic 139 displacement, D_v is the total postseismic viscoelastic displacement, τ is the characteristic 140 relaxation time, and H(t) is the Heaviside function. The parameters U_0 , V_0 , D_e , D_v and τ are 141 142 station dependent and obtained with a non-linear regression of the observed GPS time series. For all selected stations, the secular velocity can be stably estimated from the pre-earthquake 143 observations that were generally surveyed 3-4 times since 1999. Note that a few stations are 144 affected by the coseismic displacements caused by the 2013 Lushan (Mw6.6) and 2014 145 Kangding (Mw6.2) earthquakes, which have been corrected theoretically based on the slip 146

models inferred from geodetic data (Jiang el al., 2014; Jiang el al., 2015). It should be
mentioned that only annual campaign-mode GPS observations are available for this study, and
it is known that the effect caused by seasonal hydrological variation is hard to estimate
directly from such campaign data. Therefore, we model this effect based on the approach of
Chanard et al. (2014) and remove it from the GPS time series, although the resulted correction
to the horizontal components (< 2 mm) is rather small.

Because the first postseismic GPS observations in the CMONOC dataset were 153 obtained ~ 1 year after the event (~ 2009.5), the unknown parameters in Eq. (1) are only 154 loosely constrained, especially the characteristic decay time (τ). As shown in Fig. S2 in the 155 supplementary material, a range of curves with different relaxation times can fit the 156 157 postseismic displacements equally well. The quasi coseismic offsets, which consist of the coseismic offset and the postseismic deformation in the initial few weeks, are important for 158 estimating the unknown parameters in Eq. (1) (Fig. S2). For this reason, the quasi coseismic 159 displacement offsets shown in Wang et al. (2011) are included in our displacement time 160 series, too. All GPS time series in comparison with their regression curves (after removing the 161 pre-event position and the secular trend) are shown in the supplementary material (Fig. S1). 162

The accumulative postseismic displacements for the first 7 years after the Wenchuan 163 earthquake are shown in Fig. 1b and Table S1, which are derived from the observed time 164 series using the exponential fitting approach, and equal to $D_v \cdot \left[1 - e^{-\frac{7}{\tau}}\right]$ (Fig. S1). The 165 signals are more significant on the hanging wall than on the footwall and show the south-166 eastward thrusting pattern of the Tibetan Plateau relative to the stable Sichuan Basin. 167 Generally, the postseismic displacement field shows a gentle gradient across the 168 Longmenshan fault (LMSF, Fig. 1b), indicating the likelihood of a fully locked shallow fault 169 after the coseismic rupture. 170

In comparison, the vertical component of the postseismic data shows a significantly larger uncertainty and a lower signal-to-noise ratio than the two horizontal components (Fig. S1). In addition, the vertical component can be severely influenced by local nontectonic effects (e.g. surface water loading variation and atmospheric pressure), which are hard to remove from the observations in campaign-mode (Fu and Freymueller, 2012). For these reasons, we do not include the vertical displacements to constrain our postseismic model in this study.

178 **3. Modelling Results**

Viscoelastic relaxation of the coseismically induced stress change in the lower crust and upper mantle and the aseismic fault slip on the coseismic rupture plane (afterslip) are the two major mechanisms which are widely proposed to explain the postseismic deformation observed after large earthquakes (e.g., Freed et al., 2006; Freed, 2007; Hsu et al., 2006; Ryder et al., 2014). In this section, we investigate how and to what extent the observed postseismic deformation can be explained by these two mechanisms, both separately and jointly.

185 3.1. Pure Viscoelastic Relaxation Model

The finite element modelling (FEM) software Pylith (Aagaard et al., 2013) is used to simulate 186 the viscoelastic relaxation effect, which is based on a lateral heterogeneous earth structure. 187 The FEM extends for approximately 3000 km and 600 km in the horizontal and depth 188 directions, respectively, to sufficiently minimize the artefacts induced by using the fixed 189 condition at the bottom and side boundaries (Fig. 2a). Based on seismic tomography results 190 191 (Li et al., 2011), layered structures are defined for either side of the fault (Fig. 2b), which 192 consists of an elastic upper crust and two underlying viscoelastic layers representing the lower crust and upper mantle, and for both, a Maxwell rheology is used (e.g., Freed et al., 2017). 193

The thickness of the elastic upper crust is set to 35 km for the eastern Tibetan Plateau and 25 km for the Sichuan Basin based on the inferred seismic structure (Li et al., 2011).

A coseismic slip model is imposed as the driving source of the pure viscoelastic relaxation model (PVR). Considering the self-consistency in the methodology and parameter setting used for modelling the co and postseismic deformation, which can avoid potential error propagation (Freed et al., 2017), we re-estimate the coseismic slip distribution by inverting the coseismic GPS data (Wang et al., 2011).

A constrained least squares method (Wang et al., 2013) was employed to infer the coseismic slip model, in which an *a-priori* smoothing constraint was employed to make the result stable. We applied the steepest descent method to search for the optimal solution and defined the cost function as

205

$$F(s) = \|Gs - y\|^2 + \beta^2 \|Hs\|^2$$
(2)

where s is the slip vector of each sub-fault, G is the Green's function matrix calculated based 206 on an averaged layered earth structure (Li et al., 2011), y is the coseismic GPS displacements; 207 H denotes the finite difference approximation of the Laplacian operator, and β is called the 208 smoothing factor, which controls the trade-off between model roughness and data misfit (Fig. 209 210 S3). The fault geometry used for the coseismic slip model is slightly modified from that of Wang et al. (2011), which consists of a shallow ramp fault (with a large dip of $\sim 55^{\circ}$) 211 212 connected by a deep décollement with a small dip angle (7°). To avoid the unrealistic sudden change of fault dip angle, we defined a listric ramp at depth, which makes the geometry 213 smoother compared to that of Wang et al. (2011) and closer to the geologically inferred listric 214 fault geometry (Zhang et al., 2010). 215

As will be shown in Section 3.3, the obtained coseismic slip is in general agreement with previous slip models (Shen et al., 2009; Wang et al., 2011). After imposing the coseismic slip model in the FEM, viscosities of the lower crust and upper mantle become the key

parameters that influence the model results. Further investigations in this section show that the postseismic surface deformation in the first 7 years is less sensitive to the assumed upper mantle viscosity (Fig. S4). Therefore, we fix the upper mantle viscosities used in the FEM at 1.0×10^{19} Pa·s for the Tibetan Plateau and 1.0×10^{20} Pa·s for the Sichuan Basin following Huang et al. (2014), but the lower crustal viscosities (LCVs) will be optimized to best fit the data.

A grid search approach is used to optimize the lower crustal viscosities of the eastern Tibetan Plateau (η_T) and Sichuan Basin (η_S) in the PVR model. By fitting the observed displacement time series, the misfit function to be minimized is defined by

228
$$F(\eta_T, \eta_S) = \sum_{i=1}^{M} \sum_{j=1}^{N_i} \sigma_{ij}^{-2} [U_{ij} - D_{ij}(\eta_T, \eta_S)]^2, \qquad (3)$$

where *i* and *j* represent the indices of the GPS stations and the station-dependent time samples of the observed time series, respectively, σ is the standard observation error, *U* and *D* are the observed and predicted data, respectively. The best-fit viscosities that minimize $F(\eta_T, \eta_s)$ are found to be $\eta_T = 1.0 \times 10^{18}$ Pa·s and $\eta_S \ge 10^{20}$ Pa·s (Fig. 3a). The optimal PVR model can generally explain the middle- to far-field observations (Fig. 3b), but it significantly underestimates the near-field observations.

235 3.2. Pure Afterslip Model

For testing the pure afterslip (PAS) mechanism, we investigate two different extensions of the fault plane in depth (Fig. 4). In the first model (PAS-1), the fault geometry is constructed by extending the coseismic rupture described in section 3.1. A sub-horizontal décollement fault (with a dip angle of 7° and a width of 90 km) is added at the lower edge of the shallow listric fault to include possible deep slip (Wang et al., 2011). An initial test shows that far-field displacements will be significantly underestimated using this 90-km wide décollement (Fig. S5). We therefore extend the width of the décollement to 300 km for modelling the far-field observation. In the second model (PAS-2), the shallow listric fault is extended to the bottom
of the crust (~ 65 km) with a variable dip angle, which decreases with depth from ~ 55° to 5°.
From a mechanical point of view, afterslip should appear around or between the coseismic
slip asperities rather than in the ductile lower crust. With these two typical but speculative
extensions of the fault, we may be able to see whether significant, deep afterslip, though
physically unreasonable, are required here to explain the observations.

Compared with model PAS-1, model PAS-2 permits afterslip at larger depths within a 249 250 narrower zone. For both PAS models, the finite fault is divided into sub-faults with a uniform size of $10 \text{ km} \times 10 \text{ km}$. The afterslip distributions are estimated by inverting the 7-year 251 cumulative postseismic GPS displacement data with the same inversion method (Wang et al., 252 2013) as used for the coseismic slip inversion in Section 3.1. Green's functions needed for the 253 inversion are calculated based on an average layered earth structure (Li et al., 2011). The 254 elastic inhomogeneity across the fault is neglected here because its influence on the elastic 255 deformation is insignificant (Diao et al., 2014). 256

The obtained afterslip models and their data fitting are shown in Fig. 4 and Fig. S6. A common and notable feature of both models is the widely distributed afterslip on deep patches of the fault extensions. The postseismic displacements are better fitted by PAS-1 than by PAS-2 (Fig. 4). However, if the width of the horizontal décollement in model PAS-1 is not large enough, the far-field displacements on the hanging wall will be significantly underestimated (Fig. S5). In addition, it should be noted that the apparent afterslip distributions in both models are beyond the depth of the aftershock seismicity (Fig. 4d).

264 3.3. Combined Model Incorporating Viscoelastic Relaxation and Stress-Driven Afterslip

As mentioned in the Introduction, it has been recognized that viscoelastic relaxation and

afterslip could jointly govern the postseismic process (Diao et al., 2014; Huang et al., 2014;

²⁶⁷ Freed et al., 2017). Our results shown in Sections 3.1 and 3.2 also suggest that separated

mechanisms either fail to explain the observations or lead to physically unreasonable results.
Therefore, a combined model that incorporates afterslip and viscoelastic relaxations is
necessary in the present case. Contrary to previous studies that generally assume that the
afterslip and viscoelastic relaxation are independent, we consider the viscoelastic relaxation
induced by time-dependent afterslip and build the combined model in a more reasonable and
realistic way:

(1) As mentioned in Section 3.2, like the aftershocks, afterslip should appear in areas 274 surrounding the large coseismic slip where the Coulomb failure stress (CFS) is dramatically 275 raised by coseismic stress adjustment (e.g., Hsu et al., 2006). Therefore, it can be assumed 276 that afterslip is potentially induced by the coseismic stress change. Following this assumption, 277 we fix the spatial pattern of the afterslip to be the same as the positive coseismic CFS change. 278 279 The latter is calculated based on the coseismic slip model obtained in Section 3.1 (Figs. 5a and S7). A uniform frictional coefficient of 0.5 is assumed, but a range of values for this 280 parameter (0.2~0.8) were tested to confirm that the pattern of CFS change remains stable. 281 Areas of high coseismic slip (Fig. 5a) show a clear drop in CFS, whereas the adjacent patches 282 around them are stress-enhanced (Fig. 5b). 283

(2) Following the exponential model used in Eq. (1), we assume that the rate of afterslip
 decays exponentially and uniformly with time,

286
$$\dot{s}(\mathbf{r},t) = \dot{s}_0(\mathbf{r}) \cdot e^{-\frac{t}{\tau}}, \quad (t > 0),$$
 (4)

where \boldsymbol{r} is the spatial location of the afterslip, $\vec{s}_0(\boldsymbol{r})$ is the initial afterslip rate that is assumed to be proportional to the positive coseismic CFS change, τ is the uniform characteristic decay time. After fixing the spatial and temporal pattern of the afterslip, the magnitude of afterslip is related to the positive coseismic CFS change (ΔS) with a scaling factor (α),

291
$$\int_0^T \dot{s}(r,t')dt' = \alpha \cdot max[\Delta S(r), 0], \qquad (5)$$

where *T* is the observation period (here 7 years). Thus, the space and time dependent afterslip is controlled by only two free parameters, i.e., the scaling factor α and the decay time τ .

(3) The transient postseismic displacements predicted by the combined model thatincorporates viscoelastic relaxation and afterslip can be expressed in the form

296
$$D(\mathbf{r},t) = D_{cv}(\mathbf{r},t) + \alpha \left[D_{ae}(\mathbf{r})\tau \left(1 - e^{-\frac{t}{\tau}} \right) + \int_{0}^{t} D_{av}(\mathbf{r},t')e^{-\frac{(t-t')}{\tau}}dt' \right],$$
(6)

where $D_{cv}(\mathbf{r}, t)$ represents the postseismic viscoelastic displacement due to the viscoelastic relaxation induced by the coseismic slip, $D_{ae}(\mathbf{r})$ and $D_{av}(\mathbf{r}, t)$ are the afterslip-induced elastic and viscoelastic displacements, respectively. All the three variables D_{cv} , D_{ae} and D_{av} are calculated using the same FEM and fault geometry described in Section 3.1.

(4) As known from Section 3.1, the lower crustal viscosity of Sichuan Basin can only be constrained with a lower bound at about 1.0×10^{20} Pa·s (Fig. 3a). Therefore, we fix this parameter to further reduce the degrees of freedom of the combined model. Finally, the combined model is controlled by three parameters, i.e., the lower crustal viscosity of the Tibetan Plateau (η_T), the scaling factor (α) and the decay time (τ) for afterslip, which need to be optimized by minimizing the misfit function of the whole displacement time series, as defined by Eq. (3),

308
$$F(\eta_T, \alpha, \tau) = \sum_{i=1}^{M} \sum_{j=1}^{N_i} \sigma_{ij}^{-2} [U_{ij} - D_{ij}(\eta_T, \alpha, \tau)]^2.$$
(7)

It should be noted that only the viscosity η_T and the decay time τ are nonlinearly related to the data. For any given combination of η_T and τ , the scaling factor α can be determined uniquely by the linear approach. Therefore, the present 3D nonlinear fitting problem is actually solved with a grid search in a 2D parameter space.

The best-fit lower crustal viscosity for the Tibetan Plateau is found to be $\eta_T = 2.0 \times$ 10^{18} Pa·s, which is twice as large as that of the PVR model (Figs. 3a and 6a). The best-fit 314 decay time of the afterslip is given by $\tau \sim 1.2$ years, which implies that more than 80% of the 315 afterslip was released already within the first 2 years after the event (Fig. S8). The consequent 316 317 best-fit scaling factor between the magnitude of afterslip and the positive coseismic CFS change is $\alpha = 0.35$ m/MPa, which results in the afterslip distribution shown in Fig. 5c and 318 Fig. S7. The viscosity η_T and the scaling factor α are better resolved than the decay time τ 319 (Figs. 6a-c). The reasons are the insufficient time sampling rate and limited data length. 320

The seismic moment released by the afterslip is 1.6×10^{20} Nm, ~ 25% of that of the 321 main shock and corresponding to an Mw7.4 earthquake. We found that most of the afterslip is 322 distributed at the down-dip extension of the coseismic rupture, whereas the local afterslip is 323 observed on shallow patches among the coseismic slip concentrations. The aftershocks with 324 $M \ge 3$ for the same 7-year period (data obtained from China Earthquake Networks Center, 325 http://www.csndmc.ac.cn/newweb/data.htm#, last accessed on 01/08/2017), released only ~ 326 1.1% of the seismic moment of the main shock, which indicates that the moment release 327 during the postseismic process is essentially aseismic. 328

The predicted postseismic displacements of the combined model match the data well 329 330 in terms of total magnitude (Fig. 6d) and temporal evolution (Fig. 7 and Fig. S1). The complementary contributions from the two main mechanisms can be revealed as expected, 331 i.e., the near-field postseismic deformation is mainly caused by the stress-driven afterslip, 332 whereas in the middle- to far-field it appears to be more controlled by the viscoelastic 333 relaxation (Fig. 7). In the time domain, the contribution from the afterslip is mostly 334 accomplished in the initial 2 years, whereas the contribution from the viscoelastic relaxation 335 is sustained and plays a dominant role on the postseismic displacements in the following 336 periods (Fig. 7). 337

338 **4. Discussion**

4.1. Joint Process of Viscoelastic Relaxation and Stress-Driven Afterslip

To explain the far-field postseismic deformation, widespread, deep afterslip beyond the depth 340 of most aftershocks (Fig. 4) appears to be required by the two pure afterslip models. However, 341 the low-velocity anomalies at these depths (e.g., Pei et al., 2014; Wang et al., 2015) imply 342 ductile rock materials, which might not be able to generate brittle dislocations (Pei et al., 343 2014). Thus, despite both PAS models providing a satisfactory fit to the data, they are 344 345 physically unreasonable. The apparent widespread deep afterslip only indicates the spatial extension of the driving source. In contrast, the pure viscoelastic model can reasonably 346 explain the data in the middle to far-field, but it fails in the near-field, which suggests that the 347 afterslip cannot be completely ruled out. 348

The combined model that incorporates viscoelastic relaxation in the lower crust and 349 stress-driven afterslip on the coseismic rupture plane can overcome the limitations of the 350 unitary mechanisms. In the case of the Wenchuan earthquake, stress-driven afterslip is mostly 351 concentrated at depths less than 30 km, which is close to the depth of the aftershock cloud 352 (Fig. 4d). Moreover, we observe that the stress-driven afterslip exhibits a closer spatial 353 correlation with the aftershock seismicity than the free afterslip derived from the residual data 354 of the pure viscoelastic relaxation model (Figs. 5e and 5f), indicating that the combined model 355 works better, even with only three free parameters (η_T, α, τ). 356

The temporal evolution of the afterslip should be considered to incorporate the secondary viscoelastic relaxation induced by it. In this study, we assume that the rate of afterslip decays exponentially with time. It should be noted that the mathematical form of the afterslip rate is not unique and other functions have been used to describe the temporal history of afterslip. For example, Helmstetter and Shaw (2009) adopted the same Omori law of the aftershock seismicity for the afterslip rate, $\dot{s}(t) \propto 1/(1 + t/\tau)^p$, $(p \sim 1)$. To test the stability

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of our results, we also tested the performance of the Omori law for the afterslip rate in the present case. As a result, we find that the corresponding optimal lower crustal viscosity of the eastern Tibetan Plateau (η_T) and scaling factor (α) remains the same as that inferred from the exponential form. The optimal Omori decay time is found to be ~0.18 year that is in fact comparable with the exponential decay time of 1.2 years (Fig. S8).

In most previous studies, afterslip is normally assumed to be independent from 368 viscoelasticity and therefore only its elastic deformation is considered (Freed et al., 2007). 369 370 However, afterslip itself should in principle induce viscoelastic relaxation t in addition to that driven by the coseismic slip. To compare the three different effects on the postseismic 371 deformation, i.e., the viscoelastic effect induced by the coseismic slip, the elastic and 372 viscoelastic effects induced by the afterslip, we calculate surface displacements on a profile 373 perpendicular to the earthquake fault using the optimal combined model. As shown in Fig. 7a, 374 the displacement induced by the viscoelastic effect of the afterslip can reach ~ 2 cm on the 375 hanging wall, $\sim 15\%$ of the total postseismic displacements in the first 7 years after the 376 Wenchuan earthquake. Generally, this effect contributes ~ 8-15% of the total deformation in 377 the period, which is comparable with the results from the postseismic study of the 2011 378 Tohoku earthquake by Freed et al. (2017). Furthermore, it is notable that the afterslip and 379 viscoelastic relaxation should co-evolve, i.e., not only will there be viscoelastic relaxation 380 caused by the afterslip, but also the viscoelastic relaxation would change the stress conditions 381 under which the afterslip occurs. To evaluate the influence of the latter, we calculate the CFS 382 change caused by coseismic viscoelastic relaxation. As shown in the Fig. S9, the stress 383 induced by coseismic viscoelastic relaxation is rather small compared with that induced by 384 sudden coseismic slip (< 7% in the first two years). Therefore, the coseismic CFS change 385 plays a dominant role in governing the spatial pattern of the afterslip, whereas that induced by 386 viscoelastic relaxation within the observation period considered in this study is negligible. 387

4.2. Lateral Inhomogeneity of the Lower Crustal Viscosity across the LMS

The lower crust beneath the eastern Tibetan Plateau has been characterized by low seismic velocity and high electrical conductivity that may be attributed to high temperatures and elevated fluid content (e.g., Bai et al., 2010; Pei et al., 2014). The large contrast in the lower crustal viscosity across the Longmenshan fault, which is inferred from the postseismic GPS data in this study, suggests a weak, lower crust beneath the eastern Tibetan Plateau in contrast to the stable Sichuan Basin, and it is therefore consistent with previous seismic and magnetotelluric results.

396 Postseismic geodetic data have also been used for investigating the transient lower crustal viscosity beneath the Tibetan Plateau. The obtained effective lower crustal viscosity 397 398 generally increases with time during the postseismic deformation process (Fig. 8), which is 399 consistent with the stress-dependent behaviour inferred from laboratory experiments (e.g., Freed et al., 2006). This time and stress-dependent behaviour may provide a clue to 400 reconciling the difference between the effective viscosities inferred from postseismic data 401 402 over short time scales and that from geodynamic/geological models over longer time scales (Clark and Royden, 2000). In addition, it should be noted that the lower crustal viscosity of 403 the Tibetan Plateau inferred from the combined model is higher than that of the pure 404 viscoelastic relaxation model (Figs. 3a and 6a), which indicates that ignoring the afterslip 405 effect may lead to an underestimation of the lower crustal viscosity. 406

Lateral variations in the crustal viscosity may exist, as other studies in the northeast Tibetan Plateau suggest (Ryder et al., 2011; Wen et al., 2012). Speculating that such heterogeneities might also occur elsewhere, possibly on different scales, complexity in postseismic deformation can be expected for adjacent active faults. Moreover, lateral variations of viscosity will change the strength of crust, and affect the pattern of fault

deformations at different stages of a seismic cycle, which might warrant further studies in thefuture.

4.3. Seismic Hazard Assessment at the Extensions of the Wenchuan Rupture 414 The combined model indicates shallow afterslip on the northeastern extension of the 415 coseismic rupture (Figs. 5d and 5e), which agrees with the near-field postseismic InSAR 416 (Interferometric Synthetic Aperture Radar) observations that reveal short wavelength 417 deformation across this fault segment (Huang et al., 2014). The shallow aftershocks also 418 419 suggest that the shallow part of this segment was active under the loading of deep coseismic slip (Luo et al., 2010). Combined with this evidence, we infer that the stress on the Qingchuan 420 segment might have been released (Fig. 5e), whereas the stress on the more northern segment 421 $(> 32.5^{\circ} \text{ N})$ could be enhanced. 422

Compared to the north-eastern segments, the southwestern extension of the Wenchuan 423 rupture has attracted more attention because the 2013 Mw6.6 Lushan earthquake left a 50-km-424 long seismic gap (Fig. 1). The fault behaviour and related seismic hazards of this segment 425 have been subject to debate. Low seismic velocity anomalies in this area that inferred from 426 geophysical observations suggest that this segment might not be strong enough to generate 427 brittle deformation (e.g., Pei et al., 2014; Wang et al., 2015). However, geological studies 428 indicated that historical earthquakes have in fact ruptured this segment (Dong et al., 2017; 429 Wang et al., 2014). 430

In this study, negligible afterslip is observed on this seismic gap (Fig. 5c and 5e), which indicates that this fault segment may be locked. This result is obviously confirmed by the near-field displacement time series of station QLAI (Fig. S10), which indicates that the eastward motion on the footwall has increased instead of decreased after the Wenchuan earthquake. Moreover, the absence of aftershocks in the region also contradicts the afterslip hypothesis. The Lushan earthquake is generally believed to have been triggered by the

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Wenchuan earthquake (Jia et al., 2014) or even regarded as a strong aftershock of the latter
(Zhu, 2016). Our observations reveal that the locked fault segment between the two
successive major earthquakes either acts as a barrier that resists rupture propagation, or it has
not yet accumulated a sufficient slip deficit. The latter seems more reasonable as historical
earthquakes have ruptured this segment (Dong et al., 2017; Wang et al., 2014).

Similar seismic gaps between two adjacent successive ruptures have been observed
both on continental faults and in subduction zones (e.g., Delouis et al., 2010). These gaps
perhaps were caused by separated asperities, which remain locked to accumulate strain for the
next rupture even under the boost of adjacent events. Such rupture segmentation along a fault
could be due to fault geometry complexities (Hubbard et al., 2016), frictional behaviour
(Manighetti et al., 2007) or different fault locking degrees (Moreno et al., 2011).

448 4.4 Limitations

In this study, we investigate separated and combined postseismic mechanisms based on the 449 GPS displacements in first 7 years following the Wenchuan earthquake. Due to the spatial and 450 temporal complexity of the co-evolved processes, we fix the spatial pattern of afterslip to 451 decrease the degree of freedom in the combined model by assuming that afterslip is 452 potentially controlled by coseismic CFS change. This assumption is made following a general 453 consensus that afterslip is mostly distributed in areas surrounding the coseismic slip asperity, 454 where the coseismic CFS are clearly enhanced. Besides the loading of the coseismic rupture, 455 the CFS is also affected by the pre-seismic stress state of the fault. However, due to the lack 456 of efficient observations, the pre-seismic stress state on the LMSF is not available and 457 remains hard to estimate. Therefore, we assume that the afterslip is governed by the coseismic 458 CFS change and the effect of pre-seismic stress state is negligible. As shown by the results, 459 the combined model can explain the observation well in terms of the total amplitude and the 460 temporal evolution (Fig. 6d and Fig. S1). Moreover, the stress-driven afterslip exhibits a 461

closer spatial correlation with the aftershock seismicity than the free afterslip derived from the
residual data of the pure viscoelastic relaxation model (Figs. 5e and 5f), suggesting that this
assumption works as a first-order approximation.

In addition, the inferred coseismic CFS change depends on the inverted coseismic slip 465 model, which may vary by using different inversion approaches and model configurations 466 (Shen et al., 2009; Wang et al., 2011 and many others). For example, the choose of the 467 smoothing factor [see Eq. (2)], which controls the spatial roughness of the coseismic slip 468 model, could change the coseismic slip distribution, especially for some local slip asperities 469 (Wang et al., 2011). To investigate how the CFS change will be affected by the coseismic slip 470 471 model, we choose different smoothing factors to infer coseismic slip model, and calculate the corresponding CFS change. As shown in Fig. S3, the first-order pattern of the CFS change 472 remains stable though the coseismic slip model changes with the smoothing factors varying 473 474 near the corner of the trade-off curve.

Finally, it should be mentioned that the poroelastic rebound effect of the permeable 475 uppermost crust is ignored in this study. In fact, the magnitude of poroelastic contribution to 476 the postseismic deformation can be estimated based on the simple approach as presented in, 477 e.g., Huang et al. (2014). In this approach, the complete poroelastic rebound effect can be 478 479 evaluated using the difference of elastic coseismic displacements associated with variable Poisson ratios under undrained and drained conditions. Our test model for the poroelastic 480 effect has the same layering structure as the elastic one used for the coseismic slip inversion, 481 but the uppermost part (up to 4 km) is assumed to be poroelastic. The Poisson ratio of this 482 layer can vary by up to 30% depending on whether it is under drained or undrained conditions. 483 Using the Poisson ratio of 0.25 under undrained conditions (the original seismic reference) 484 and 0.21 under drained conditions, the model results show that the postseismic displacement 485 caused by poroelastic rebound effect is dominated by the vertical component and the 486 maximum of the long-term accumulative value is estimated to be about 1 cm, which is at least 487

by one order of magnitude smaller than the observed accumulative displacements in the first 7
years after the Wenchuan earthquake. We also test different thickness (2-4 km) of the
poroelastic layer, but found no substantial change on the results. Therefore, we ignored the
poroelastic rebound effect in our models.

492 **5. Conclusions**

Using the GPS observations acquired during the first 7 years after the 2008 Wenchuan 493 earthquake, the related mechanisms are investigated, which include lower crustal viscoelastic 494 495 relaxation, afterslip and a combination of both. In contrast to previous postseismic studies, we built a combined model that incorporates the viscoelastic relaxation induced by both the 496 coseismic slip and stress-driven afterslip. The secondary viscoelastic relaxation induced by 497 afterslip is found to have a significant effect on both near- and far-field postseismic 498 displacements and therefore should be considered for the postseismic analysis of large 499 earthquakes elsewhere. Lower crustal viscoelastic relaxation is found to be responsible for 500 most of the middle- to far-field postseismic deformation, while the near-field displacements 501 are governed mainly by the stress-driven afterslip. The estimated seismic moment released by 502 stress-driven afterslip is up to 25% of that released by the coseismic rupture, which is 503 equivalent to an Mw7.4 earthquake. Negligible afterslip is observed in the seismic gap 504 between the Wenchuan and Lushan ruptures. Combined with the observation of very low 505 aftershock seismicity during the same period, we deduce that this segment is in a locked state. 506 The effective lower crustal viscosity of the eastern Tibetan Plateau is estimated to be $2.0 \times$ 507 10^{18} Pa·s, whereas that of the Sichuan Basin should be larger than 1.0×10^{20} Pa·s, 508 suggesting a significant lateral viscosity contrast across the LMS fault. 509

510 Acknowledgments

511	This work was supported by the National Natural Science Foundation of China (Grant Nos.
512	41674023 and 41731072) and the China Earthquake Science Experiment Program
513	(2017CESE0103). The data were obtained from the Crustal Movement Observation Network
514	of China (CMONOC). Jeffrey Freymueller and an anonymous reviewer provided valuable
515	comments and suggestions to improve the quality and readability of the revised manuscript.
516	We thank Kevin Fleming, Jing Zeng-Liu, and Yanqiang Wu for their valuable comments on
517	the original manuscript. The figures are drawn using the Generic Mapping Tools (GMT)
518	software (Wessel and Smith, 1998).
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1	Supplementary Material
2	for
3	Fault behaviour and lower crustal rheology inferred from the first seven years of
4	postseismic GPS data after the 2008 Wenchuan earthquake
5	
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13	Figures S1-10
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15	Introduction

- 16 Figures S1-10 show the 7-years GPS time series and additional model results, and Table S1
- 17 lists the cumulative postseismic displacements measured at all selected stations.



18

Figure S1. Displacement time series used in this study. The blue circles show the postseismic displacements with error bars of 95% confidence. Black and red curves represent the best-fit curve defined by Eq. (1) in the main text and the theoretical time series of the combined model. The green lines are the reference level obtained by removing the preseismic background trend.









Figure S1 (continued)







Figure S2. Example showing the curve fitting of the postseismic displacement time series. Curves with different colors represent the fitting with variable rise times. Blue circles are the

data provided by Crustal Movement Observation Network of China (CMONOC), and the red

circle is the quasi coseismic offset adopted from Wang et al. (2011).



Figure S3. Results of coseismic slip inversion. Left: coseismic slip model and CFS change with variable smoothing factors; Right: trade-off curve with data misfit plotted as function of

with variable smoothing factors; Right: trade-off curve with data misfit plotted as function of
normalized model roughness. The shaded area shows the smoothing factors near the corner of

- the trade-off curve.
- 40
- 41
- 42





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Figure S4. Normalized misfit variance of the pure viscoelastic relaxation model depending on the viscosity of the upper mantle beneath the Tibet Plateau (η_m), where the other viscosities are fixed at their optimal values (e.g., $\eta_T = 1.0 \times 10^{18}$ and $\eta_S = 1.0 \times 10^{20}$).



Figure S5. Afterslip distribution and data fitting for Model PAS-1 with the deep horizontal
décollement limited at a length of 90 km.



Figure S6. Slip distribution of the pure afterslip model with arrows showing the direction ofslip on each fault patch.



Figure S7. Slip distribution of the coseismic slip model, stress-driven afterslip model and free
afterslip distribution derived from the residual data of the PVR model with arrows showing
the direction of slip on each fault patch.



Figure S8. Comparison between two models for the time history of afterslip rate. The characteristic decay time τ used is 1.2 and 0.18 year for the blue and red curves, respectively.





Figure S9. (a) CFS change on fault plane induced by coseismic rupture. (b) Temporal evolution of CFS changes caused by coseismic slip and viscoelastic relaxation on four typical points shown in (a). The gray line shows the time of the event. The start point of each line shows the coseismic CFS change.

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Figure S10. GPS displacement time series of Station QLAI before and after the Wenchuan earthquake (Wu et al., 2013). The two solid (or dashed in parallel) red lines show the pre-event linear trends for the two horizontal components, respectively, in comparison with the green lines for the post-event linear trends. The location of QLAI is shown in Fig. 1b in the main text.

Station	Longitude	Latitude	De	Dn	σDe	σDn
	(deg)	(deg)	(cm)	(cm)	(cm)	(cm)
H007	106.155	33.340	1.34	0.35	0.09	0.13
H009	106.023	32.962	-0.66	-0.46	0.22	0.23
H010	105.226	32.571	6.31	5.16	0.11	0.15
H011	105.830	32.448	-0.20	0.20	0.17	0.11
H012	105.457	32.018	-1.18	-0.31	0.09	0.14
H021	104.824	33.423	0.99	1.41	0.30	0.17
H022	104.625	33.001	2.68	1.24	0.06	0.13
H024	104.226	33.228	2.50	-0.34	0.13	0.18
H025	103.435	32.931	4.32	-2.65	0.12	0.08
H027	101.482	33.429	1.11	-1.28	0.15	0.20
H028	101.706	32.902	3.56	-1.50	0.16	0.44
H029	100.595	33.093	1.25	-0.94	0.11	0.31
H030	103.613	32.591	5.57	-2.47	0.15	0.12
H031	102.500	32.786	4.12	-3.85	0.23	0.31
H032	104.571	32.405	7.90	3.81	0.27	0.09
H033	104.831	32.182	6.01	1.84	0.19	0.16
H034	103.732	32.361	7.89	-0.98	0.22	0.07
H035	104.444	31.802	5.98	0.55	0.18	0.12
H037	103.166	32.075	9.71	-3.31	0.15	0.10
H040	101.614	31.770	6.07	-0.37	0.35	0.15
H041	101.071	32.319	3.13	-0.54	0.17	0.26
H042	100.334	32.275	1.05	-0.65	0.12	0.19
H043	104.782	31.486	-0.97	-0.51	0.12	0.11
H044	104.187	31.353	2.94	-0.90	1.28	0.23
H045	103.612	31.474	14.87	-0.49	0.31	0.31
H046	102.670	31.850	11.04	-3.21	0.18	0.15
H047	102.096	31.466	7.77	-0.44	0.17	0.08
H048	104.441	31.157	-0.46	1.11	0.09	0.12
H049	103.692	31.060	1.96	-1.62	0.13	0.11
H050	103.145	31.008	8.21	-5.00	0.17	0.39
H051	102.775	30.992	7.59	-1.63	0.17	0.22
H052	101.866	30.949	5.31	0.82	0.21	0.10
H053	101.163	30.955	4.16	0.78	0.09	0.12
H054	100.750	31.297	2.84	0.97	0.06	0.13
H056	100.297	31.646	2.08	-0.46	0.21	0.20
H060	103.410	30.415	0.35	-0.37	0.16	0.25
H063	100.307	30.925	0.87	1.13	0.13	0.24
JB24	106.034	30.804	-0.92	-0.66	0.16	0.15
JB33	103.889	33.276	0.94	-0.71	0.11	0.15

Table S1. Postseismic displacements accumulated in the first 7 years after the Wenchuan earthquake.

JB34	102.306	31.706	9.22	-1.76	0.32	0.26
JB35	101.497	30.495	2.99	0.95	0.14	0.11

88 Reference:

89	Wu Y O Z S Jiang M Wang S Che H Liao O Li P Li Y Yang H Xiang Z Shao W
05	
90	Wang, W. Wei, and X. Liu (2013), Preliminary results pertaining to coseismic
91	displacement and preseismic strain accumulation of the Lushan M(S) 7.0 earthquake,
92	as reflected by GPS surveying, Chin. Sci. Bull., 58, 3460-3466, doi:
93	10.1007/s11434-013-5998-5.
94	
95	
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Figure 1. (a) Tectonic setting of the study area. The black and blue contours show the coseismic slip distributions of the 2008 Wenchuan (this study) and 2013 Lushan earthquakes (Jiang et al., 2014) projected on the surface using intervals of 2 m and 0.2 m, respectively. The thick red lines indicate the surface rupture trace

of the Wenchuan earthquake (Xu et al., 2009). The gray dots denote aftershocks of the two earthquakes that occurred within the first year (Fang et al., 2015). The dashed black box represents the seismic gap between the ruptures of the two earthquakes. (b) Observed cumulative postseismic surface displacements during the first 7 years after the Wenchuan earthquake. The gray lines represent the active faults in the study area. The blue vectors denote the observed postseismic displacements, and the ellipses mark the uncertainties with 95% confidence. The inset shows the normalized coseismic (black) and postseismic (red) displacement profiles from the area marked by the dashed violet rectangle. Abbreviations: GJF = Guanxian-Jiangyou Fault, GLF = Gengda-Longdong Fault, MWF = Maoxian-Wenchuan Fault, LMSF: Longmenshan Fault, BYF = Beichuan-Yingxiu Fault, QF = Qingchuan Fault.



Figure 2. (a) Structure and mesh configuration of the finite element model. The element size increases from about 1 km for the near-fault region to about 50 km for the far-field region. (b) Sketch map and the 2D average structure across the Longmenshan fault.



Figure 3. Results of the pure viscoelastic relaxation model (PVR). (a) Misfit variance (normalized by the data variance) as a function of the lower crustal viscosities of the Tibet Plateau (η_T) and the Sichuan Basin (η_S) used in the PVR model. (b) between the observed and predicted postseismic displacements in the first 7 years the Wenchuan earthquake.



Figure 4. Comparison of the two pure afterslip models (PAS-1 and PAS-2). (a) and (b) Afterslip distributions and data fittings of Model PAS-1 and Model PAS-2, respectively. (c) Fault geometry of the two afterslip models on the profile P-P' marked in (a), where BYF and GJF represent the Beichuan-Yingxiu Fault and the Guanxian-Jiangyou Fault, as shown in Fig. 1. (d) Comparison of the aftershock seismicity (gray histograms with the right axis) and the two afterslip distributions with depth.



Figure 5. (a) Surface projection of the coseismic slip distribution. (b) Coseismic CFS changes on the fault plane. (c) Stress-driven afterslip distribution in the combined model. (d) Free afterslip distribution derived from the residual data of the PVR model. (e) and (f) Comparison of the aftershock (gray square) with the spatial distributions of the stress-driven afterslip and the free afterslip, respectively.



Figure 6. Results of the combined model (VR+SAS). (a)-(c) 3D map of the misfit variance (normalized by the data variance) projected on the (η_T, τ) , (η_T, α) and (τ, α) planes, respectively, where η_T is the lower crustal viscosity of the Tibet Plateau, τ is the characteristic decay time of the afterslip, and α is the scaling factor between the afterslip and the coseismically induced positive Coulomb stress change. The red star marks the minimum of the misfit variance, which defines the location of the 3 optimal model parameters. (d) Comparison between the observed and predicted postseismic displacements in the first 7 years after the Wenchuan earthquake.



Figure 7. Contributions of different deformation processes to the observed postseismic surface displacement. (a) East component of the 7-year postseismic displacement along the profile P-P' shown in Fig. 3a, predicted by the combined model. D_{cv} is the postseismic viscoelastic effect induced by the coseismic slip, D_{ae} and D_{av} are the elastic and viscoelastic effects induced by the aseismic afterslip, respectively. (b) The different effects in the time domain on a selected station (H037). The colour coding is the same as used in (a).



Figure 8. Effective lower crustal viscosities beneath the Tibet Plateau inferred by using postseismic geodetic data covering different time periods and for different earthquakes published in recent years (Table 1).