

Originally published as:

Huang, S. Y., Sahraoui, F., Yuan, Z. G., Contel, O. L., Breuillard, H., He, J. S., Zhao, J. S., Fu, H. S., Zhou, M., Deng, X. H., Wang, D., Du, J. W., Yu, X. D., Wang, D. D., Pollock, C. J., Torbert, R. B., Burch, J. L. (2018): Observations of Whistler Waves Correlated with Electron-scale Coherent Structures in the Magnetosheath Turbulent Plasma. - *The Astrophysical Journal*, *861*, 1.

DOI: http://doi.org/10.3847/1538-4357/aac831

1	Observations of whistler waves correlated with electron-scale coherent structures in the
2	magnetosheath turbulent plasma
3	
4	S. Y. Huang ¹ , F. Sahraoui ² , Z. G. Yuan ¹ , O. Le Contel ² , H. Breuillard ² , J. S. He ³ , J. S. Zhao ⁴ , H.
5	S. Fu ⁵ , M. Zhou ⁶ , X. H. Deng ⁶ , X. Y. Wang ¹ , J. W. Du ¹ , X. D. Yu ¹ , D. D. Wang ^{1,7} , C. J. Pollock ⁸ ,
6	R. B. Torbert ⁹ , and J. L. Burch ¹⁰
7	
8	¹ School of Electronic Information, Wuhan University, Wuhan, China
9	² Laboratoire de Physique des Plasmas, CNRS-Ecole Polytechnique-UPMC, Palaiseau, France
10	³ School of Earth and Space Sciences, Peking University, Beijing, China
11	⁴ Key Laboratory of Planetary Sciences, Purple Mountain Observatory, Chinese Academy of
12	Sciences, Nanjing, China
13	⁵ School of Space and Environment, Beihang University, Beijing, China
14	⁶ Institute of Space Science and Technology, Nanchang University, Nanchang, China
15	⁷ Helmholtz Centre Potsdam GFZ German Research Centre for Geosciences, Germany
16	⁸ NASA, Goddard Space Flight Center, Greenbelt, Maryland, USA
17	⁹ Physics Department, University of New Hampshire, Durham, New Hampshire, USA
18	¹⁰ Southwest Research Institute, San Antonio TX, USA
19	
20	
21	Abstract
22	A new type of electron-scale coherent structure, referred to as electron vortex magnetic holes,
23	was identified recently in the Earth's magnetosheath turbulent plasma. These electron-scale
24	magnetic holes are characterized by magnetic field strength depression, electron density
25	enhancement, temperature and temperature anisotropy increase (significant increase in
26	perpendicular temperature and decrease in parallel temperature), and an electron vortex formed
27	by the trapped electrons. The strong increase of electron temperature indicates that these
28	magnetic holes have strong connection with the energization of electrons. Here, using high time

30 whistler waves coexist with electron-scale magnetic holes. These whistler waves were found

resolution in-situ measurements from the MMS mission, it is further shown that electron-scale

29

not propagating from remote regions, but generated locally due to electron temperature anisotropy $(T_{e\perp}/T_{e\parallel})$ inside the magnetic holes. This study provides new insights into the electron-scale plasma dynamics in turbulent plasmas.

- 34
- 35

36 1. Introduction

37 Magnetic holes, characterized by magnetic field strength depressions lasting from several 38 seconds to several minutes (corresponding to a spatial scale larger than ion gyro-radius), have 39 been detected in space plasmas such as the magnetosphere (e.g. Shi et al. 2009; Sundberg et al. 40 2015), magnetosheath (e.g. Tsurutani et al. 2011) or the solar wind (e.g. Tsurutani et al. 2002; 41 Zhang et al. 2008). These large-scale magnetic holes are thought to be generated by mirror 42 instability (Horbury et al. 2004; Sahraoui et al., 2004, 2006) or related to slow solitons 43 (Stasiewicz, 2004). Recently, thanks to the unprecedented high time resolution data of MMS 44 mission, a new type of kinetic-size magnetic holes (less than proton gyro-radius, or tens of 45 electron gyro-radius) was evidenced in the turbulent magnetosheath plasma (Huang et al. 2017a, 46 2017b; Yao et al. 2017). A statistical study showed that an increase in the electron density and 47 total temperature, significantly increase (resp. slightly decrease) in the electron perpendicular 48 (resp. parallel) temperature (Huang et al. 2017b). Furthermore, an electron vortex is formed by 49 the trapped electrons with enhanced fluxes at $\sim 90^{\circ}$ pitch angles inside the magnetic holes 50 (Huang et al. 2017b). The increase of electron temperature inside the magnetic holes suggests 51 the possible role the latter could play in energy dissipation in turbulent plasmas leading to 52 localized electron heating (Breuillard et al. 2016a; Huang et al. 2012a, 2014, 2017d; Fu et al. 53 2017). These magnetic holes can be best interpreted as electron vortex magnetic holes 54 according to Particle-In-Cell (PIC) simulations of Haynes et al. (2015).

55

Whistler waves are frequently reported in space and laboratory plasmas (Stenzel, 1999), such
as the inner magnetosphere (e.g. Li et al. 2011; Fu et al. 2012, 2014), the magnetotail (e.g.
Huang et al. 2010, 2012b, 2016a, 2017c; Zhima et al. 2015; Breuillard et al. 2016b), the
magnetosheath (e.g. Moullard et al. 2002; Breuillard et al. 2017, Huang et al. 2016b; Cao et al.

60 2017), and the solar wind (e.g. Lacombe et al. 2014). They can be generated by various 61 mechanisms that include electron temperature anisotropies (e.g. Gary and Madland, 1985) and 62 current or beam-driven instabilities (e.g., Zhang et al. 1999; Fujimoto et al. 2014). Whistler 63 waves propagating in the magnetosphere (Santolik et al. 2009; Breuillard et al. 2012, 2014) can 64 be guided by density enhancements or depletions (e.g., Smith et al. 1961; Moullard et al. 2002; 65 Streltsov et al. 2006; Li et al. 2011) and modulated by ULF waves (e.g. Xia et al. 2016). The 66 occurrence of whistler waves is generally well correlated with the existence of magnetic 67 structures with magnetic field strength minimum (Smith and Tsurutani, 1976; Dubinin et al. 68 2007) or dipolarization fronts (Le Contel et al. 2009; Deng et al. 2010; Huang et al. 2012b; Li 69 et al. 2015). Recently, Tenerani et al. (2012, 2013) have reported the connection between 70 whistler waves and ion-scale magnetic structures using Cluster measurements in the 71 magnetotail.

72

73 However, whether whistler waves can correlate with electron-scale coherent magnetic 74 structures are still unsettled. In this Letter, we report observations of intense whistler waves 75 correlated with electron-scale magnetic holes (i.e., electron vortex magnetic holes) in the 76 turbulent magnetosheath plasma from the MMS spacecraft. The magnetic field data come from 77 the Fluxgate Magnetometer (FGM) (Russell et al. 2016) and the Search-Coil Magnetometer 78 (SCM) (Le Contel et al. 2016), the electric field data from the Electric Double Probe (EDP) 79 (Ergun et al. 2016; Lindqvist et al. 2016), and the 3D particle velocity distribution functions 80 and the plasma moments from the Fast Plasma Investigation (FPI) (Pollock et al. 2016).

81

82 2. MMS Observations

Figure 1 presents one example of an electron vortex magnetic hole in the magnetosheath. The local boundary normal coordinates (LMN) determined by minimum variance analysis (MVA) of the magnetic fields (Sonnerup and Scheible, 1998) were used. In this coordinate system, L is the maximum variation direction that can be considered as the axis of the magnetic hole, M is the intermediate variation direction, and N is the minimum variation direction. Magnetic field strength depression from ~ 19 nT to ~ 12 nT (Figure 1a) with small field shear angle (<15°) implies that a magnetic hole is detected by MMS. The electron density (Figure 1e) and electron 90 temperature (Figure 1d) increase while the electron parallel temperature decreases (Figure 1d) 91 and the electron perpendicular temperature strongly increases inside the magnetic hole (Figure 92 1h). Here, parallel and perpendicular directions are with respect to the ambient magnetic field. 93 There is a clear bipolar signature (from positive to negative) in V_{eN} component (Figure 1c) and 94 a strong enhancement in electron perpendicular temperature (Figure 1d) inside the magnetic 95 hole, which are typical features of an electron vortex (Haynes et al. 2015; Huang et al. 2017a, 96 2017b). The estimated width of magnetic hole (based on the frozen-in assumption) is about 27 97 km (0.15 ρ_i or ~ 19 ρ_e , where $\rho_i \sim 175$ km and $\rho_e \sim 1.4$ km are the proton and electron 98 gyroradius with $|B| \sim 24$ nT, $n \sim 9$ cm⁻³, $T_i \sim 850$ eV and $T_e \sim 100$ eV). Thus, MMS detected 99 clearly an electron-scale electron vortex magnetic hole in the magnetosheath.

100

101 Figure 2 illustrates the statistical results of electron temperature for all electron vortex magnetic 102 holes (66 in total) in the turbulent magnetosheath from September 2015 to January 2016 (Huang 103 et al. 2017b). It shows that the electron perpendicular temperature significantly increases inside 104 the magnetic holes $(T_{e\perp(inside)}/T_{e\perp(outside)} > 1)$ for all events (Figure 2a), while the electron parallel 105 temperature decreases $(T_{e\parallel (inside)}/T_{e\parallel (outside)} < 1)$ for most of the events (Figure 2b), with a net 106 increase in electron total temperature inside the magnetic holes for most of the events (~93%, 107 Figure 2c). The increase of electron temperature indicates that these magnetic holes have strong 108 correlation with the heating of electrons. The statistics of the ratio of $T_{e\perp}/T_{e\parallel}$ inside the magnetic 109 holes (Figure 2d) show strong electron temperature anisotropies (up to 1.7) in the electron 110 vortex magnetic holes for ~90% events. Such electron temperature anisotropy is favorable to 111 the excitation of high frequency waves inside the magnetic holes.

112

113 After surveying the wave observations for all magnetic holes, we successfully identified the 114 whistler waves inside 11 (17%) electron vortex magnetic holes. Figure 3 shows an example of 115 such identification. Magnetic field strength, the ratio of $T_{e\perp}/T_{e\parallel}$, and high-pass filtered magnetic 116 field fluctuations with cut-off frequency of 60 Hz are presented in Figure 3a-3c respectively. 117 The occurrence of strong electron temperature anisotropy ($T_{e\perp}/T_{e\parallel}$) inside the magnetic hole is 118 accompanied by large-amplitude magnetic fluctuations (up to 0.1 nT). Figures 3d-3h show the 119 power spectral densities and the properties of wave emissions derived from singular value 120 decomposition (SVD) method (Santolik et al. 2003) during this period. Intense electromagnetic 121 waves (red or yellow in Figure 3d-3e) are observed inside the magnetic hole with their 122 frequencies between $0.1f_{ce}$ and $0.5f_{ce}$, which corresponds to the frequency band where large-123 amplitude magnetic fluctuations are observed (Figure 3c). These waves have large polarization 124 degree (>0.9 in Figure 3h) and positive ellipticity (close 1 as shown in Figure 3f), and relatively 125 small propagation angles ($<30^{\circ}$ in Figure 3g). All these wave properties indicate that the 126 observed fluctuations are carried out by right-hand polarized and quasi-parallel propagating 127 whistler mode. This scenario is confirmed using the theoretical results of the linear solver of 128 the Maxwell-Vlasov equations (WHAMP, Ronmark, 1982) shown in Figure 4. The linear dispersion relations were obtained using the local plasma parameters ($|B| \sim 12$ nT, $n \sim 10$ cm⁻³, 129 130 $T_{\rm i} \sim 800 \text{ eV}$ and $T_{\rm e} \sim 100 \text{ eV}$, $T_{\rm e\perp}/T_{\rm e\parallel} = 1.4$). One can see a positive growth rate corresponding 131 to $\omega/\omega_{ce} = 0.07 \sim 0.34$ due to the electron temperature anisotropy, with maximum growth rate 132 $\gamma/\omega_{ce} = 0.04$ occurring at $\omega/\omega_{ce} = 0.23$, which is very close to ~ 0.29 ω/ω_{ce} of the peak frequency 133 of largest power spectral densities from our observations (see in Figure 3d-3e). These 134 theoretical results imply that the whistler waves inside the electron-scale magnetic holes can be 135 locally generated by electron temperature anisotropy.

136

137 **3. Discussions and Conclusions**

138 The simultaneous occurrence of high frequency whistler waves and ion-scale magnetic 139 structures (characterized by a magnetic field strength depression) was reported in the plasma 140 sheet (Tenerani et al. 2013), but the generation mechanism of the whistler waves was not 141 addressed in that work. Intense lion roars were often observed in mirror structures with scale-142 size (much/several times) larger than the ion larmor scale in the magnetosheath (Smith and 143 Tsurutani, 1976; Thore and Tsurutani, 1981; Breuillard et al. 2017). Masood et al. (2006) have 144 shown that some instances of lion roars could be locally generated, while the source of others 145 must be in remote regions of the magnetosheath. Recently, a new type of electron-scale coherent 146 structures, referred to as electron vortex magnetic hole, was identified in the turbulent 147 magnetosheath plasma (Huang et al. 2017a, 2017b; Yao et al. 2017). The present statistical 148 study showed an increase in electron temperature inside those magnetic holes (Figure 2), 149 indicating that the latter may be able to heat electrons. While the mechanism of such heating 150 remains to be elucidated, the present observations provide nevertheless a new possible way of 151 (coherent) heating of electrons in turbulent plasmas. This mechanism can be complementary to 152 those already proposed to explain energy dissipation at electron scales in solar wind turbulence 153 (e.g., Sahraoui et al. 2009). Our results show furthermore that the subsequent electron 154 temperature anisotropy instability can give rise to intense high frequency fluctuations identified 155 as quasi-parallel propagating whistler waves, which, in turn, may accelerate electrons. Indeed, 156 whistler waves were proposed to generate energetic electrons in the direction parallel to the 157 ambient magnetic field through pitch-angle scattering (e.g., Kennel and Petschek, 1966). They 158 were also shown to lead to the formation of both halo and strahl electrons in the solar wind, 159 which could explain the generation of gamma ray bursts in interplanetary space (Vocks et al. 160 2005). The first evidence of simultaneous occurrence of electron-scale whistler waves and 161 electron-scale coherent structures (i.e. electron vortex magnetic holes) reported here would help 162 to better understand electron-scale dynamics in the magnetosheath and in other distant turbulent 163 plasmas.

164

165

166 Acknowledgement

We thank the entire MMS team and instrument leaders for data access and support. This work
was supported by the National Natural Science Foundation of China (41674161, 41404132,
41374168, and 41331070), and China Postdoctoral Science Foundation Funded Project
(2015T80830). The French participation in the MMS mission was funded by CNES. Data is
publicly available from the MMS Science Data Center at http://lasp.colorado.edu/mms/sdc/.

173

174 **References**

175 Breuillard, H. Zaliznyak, Y. Krasnoselskikh, V. Agapitov, et al. 2012 Ann. Geophys. 30, 1223-

176 1233, https://doi.org/10.5194/angeo-30-1223-2012.

- 177 Breuillard, H. Agapitov, O. Artemyev, A. Krasnoselskikh, et al. 2014 Ann. Geophys. 32, 1477-
- 178 1485, https://doi.org/10.5194/angeo-32-1477-2014
- 179 Breuillard, H., E. Yordanova, A. Vaivads, et al. 2016a, Astrophys. J., 829, 54
- 180 Breuillard, H. et al. 2016b Geophys. Res. Lett. 43, 7279–7286, doi:10.1002/2016GL069188.
- 181 Breuillard, H. et al. 2017 J. Geophys. Res. Space Physics, under review.
- 182 Cao, D. et al. 2017 Geophys. Res. Lett., 44, 3954–3962
- 183 Deng, X. M. Ashour- Abdalla, M. Zhou, R. Walker, et al. 2010 J. Geophys. Res. 115, A09225
- 184 Gary, S. P. and Madland, C. D. 1985 J. Geophys. Res. 90 (A8), 7607–7610
- 185 Ergun, R. E. et al. 2016 Space Science Reviews, 199(1), 167–188
- 186 Fu, H. S. J. B. Cao, F. S. Mozer, et al. 2012 J. Geophys. Res. 117, A01203
- 187 Fu, H. S. et al. 2014 J. Geophys. Res. Space Physics, 119, 9089–9100
- 188 Fu, H. S. A. Vaivads, Y. V. Khotyaintsev, M. André, et al. 2017 Geophys. Res. Lett. 44, 37–43,
- 189 doi:10.1002/2016GL071787.
- 190 Fujimoto, K. 2014 Geophys. Res. Lett. 41, 2721–2728, doi:10.1002/2014GL059893.
- Haynes, C. T. D. Burgess, E. Camporeale, and T. Sundberg 2015 Phys. Plasmas, 22, 012309,
 doi:10.1063/1.4906356.
- Horbury, T. S. E. A. Lucek, A. Balogh, I. Dandouras et al. 2004 J. Geophys. Res. 109, A09209,
- 194 doi:10.1029/2003JA010237.
- 195 Huang, S. Y. et al. 2010 J. Geophys. Res. 115, A12211, doi:10.1029/2010JA015335.
- 196 Huang, S. Y. et al. 2012a Geophys. Res. Lett. 39, L11104, doi:10.1029/2012 GL052210.
- 197 Huang, S. Y. M. Zhou, X. H. Deng, Z. G. Yuan et al. 2012b Ann. Geophys. 30, 97–107,
- 198 doi:10.5194/angeo-30-97-2012.
- Huang, S. Y. F. Sahraoui, X. H. Deng, J. S. He et al 2014 Astrophys. J. Lett. 789, L28,
- 200 doi:10.1088/2041-8205/789/2/L28.
- 201 Huang, S. Y. et al. 2016a J. Geophys. Res. Space Physics, 121, 6639–6646
- 202 Huang, S. Y. et al. 2016b Geophys. Res. Lett. 43, 7850–7858, doi:10.1002/2016GL070033.
- 203 Huang, S. Y. et al. 2017a Astrophys. J. Lett. 836, L27, doi:10.3847/20418213/aa5f50.
- 204 Huang, S. Y. et al. 2017b J. Geophys. Res. Space Physics, 122, 8577–8588
- 205 Huang, S. Y. et al. 2017c J. Geophys. Res. Space Physics, 122, 7188–7196
- Huang, S. Y. L. Z. Hadid, F. Sahraoui, Z. G. Yuan et al. 2017d Astrophys. J. Lett. 836,

- 207 L10, doi.org/10.3847/2041-8213/836/1/L10
- 208 Kennel, C. F., and H. E. Petschek, 1966 J. Geophys. Res., 71(1), 1–28
- Lacombe, C., Alexandrova, O. Matteini, L. Santolik, O. et al. 2014 Astrophys. J. 796, 5.
- 210 doi:10.1088/0004-637X/796/1/5
- 211 Le Contel, O. et al. 2009 Ann. Geophys. 27, 2259–2275, doi:10.5194/angeo-27-2259-2009.
- 212 Le Contel, O. et al. 2016 Space Sci. Rev. 199(1), 257–282, doi:10.1007/s11214-014-0096-9.
- 213 Li, H. M. Zhou, X. Deng, Z. Yuan, et al. 2015 J. Geophys. Res. Space Physics, 120, 1086–1095,
- 214 doi:10.1002/2014JA020474.
- Li, W. J. Bortnik, R. M. Thorne, Y. Nishimura et al. 2011 J. Geophys. Res. 116, A06206,
- 216 doi:10.1029/2010JA016313.
- 217 Lindqvist, P.-A. et al. 2016 Space Sci. Rev. 199, 137–165, doi:10.1007/s11214-014-0116-9.
- 218 Masood, W. S. J. Schwartz, M. Maksimovic, and A. N. Fazakerley 2006 Ann. Geophys. 24,
- 219 1725–1735, doi:10.5194/angeo-24-1725-2006.
- 220 Moullard, O. A. Masson, H. Laakso, M. Parrot et al. 2002 Geophys. Res. Lett. 29(20), 1975,
- 221 doi:10.1029/2002GL015101.
- 222 Pollock, C. et al. 2016 Space Sci. Rev. 199, 331–406, doi:10.1007/s11214-016-0245-4.
- 223 Rönnmark, K. 1982, Computation of the dielectric tensor of a Maxwellian plasma, Tech. rep.
- 224 Russell, C. T. et al. 2016 Space Sci. Rev. 199, 189–256, doi:10.1007/s11214-014-0057-3.
- 225 Santolik O. M. Parrot, and F. Lefeuvre 2003 Radio Sci. 38 (1), 1010
- 226 Santolik, O. M. Parrot, U. S. Inan, D. Buresova et al. 2009 J. Geophys. Res. 114, A03212,
- doi:10.1029/2008JA013776.
- 228 Sahraoui, F., Belmont, G., Pin, con, J., et al. 2004, AnGeo, 22, 2283
- Sahraoui, F. G. Belmont, L. Rezeau, and N. Cornilleau-Wehrlin 2006 Phys. Rev. Lett. 96,075002.
- 231 Sahraoui, F., Goldstein, M. L., Robert, P., & Khotyaintsev, Y. V. 2009, PhRvL, 102, 231102
- 232 Shi, Q. Q. et al. 2009 J. Geophys. Res. 114, A10202, doi:10.1029/2009JA014283.
- 233 Stasiewicz, K. 2004 Phys. Rev. Lett. 93, 125,004, doi:10.1103/ PhysRevLett.93.125004.
- 234 Stenzel, R. L. 1999 J. Geophys. Res. 104(A7), 14379–14395, doi:10.1029/1998JA900120.
- 235 Sonnerup, B. U. O. and M. Scheible, Minimum and maximum variance analysis, in Analysis
- methods for Multi-spacecraft data, edited by G. Paschmann and P. W. Daly, no. SR-001 in ISSI

- 237 Scientific Reports, Chap. 1, pp. 185-220, ESA Publ. Div. Noordwijk, Netherlands, 1998.
- 238 Smith, R. L. 1961 J. Geophys. Res. 66, 3699–3707, doi:10.1029/JZ066i011p03699.
- 239 Streltsov, A. V. M. Lampe, W. Manheimer, G. Ganguli et al. 2006 J. Geophys. Res. 111, A03216,
- doi:10.1029/2005JA011357.
- 241 Sundberg, T., Burgess, D., & Haynes, C. T. 2015, JGRA, 120, 2600
- 242 Tenerani, A. et al. 2012 Phys. Rev. Lett. 109, 155,005.
- 243 Tenerani, A. O. Le Contel, F. Califano, P. Robert et al. 2013 J. Geophys. Res. Space Physics,
- 244 118, doi:10.1002/jgra.50562.
- Tsurutani, B. T. B. Dasgupta, C. Galvan, M. Neugebauer et al. 2002 Geophys. Res. Lett. 29(24),
- 246 2233, doi:10.1029/2002GL015652.
- 247 Thorne, R. M. and B. T. Tsurutani 1981 *Nature*, 293, 384–386, doi:10.1038/293384a0.
- 248 Vocks, C., Salem, C., Lin, R. P., & Mann, G. 2005, ApJ, 627, 540
- 249 Xia, Z. L. Chen, L. Dai, S. G. C laudepierre, et al. 2016 Geophys. Res. Lett. 43, 9444–9452,
- 250 doi:10.1002/2016GL070280.
- 251 Yao, S. T. et al. 2017 J. Geophys. Res. Space Physics, 122, doi:10.1002/2016JA023858.
- 252 Zhima, Z. J. Cao, H. Fu, W. Liu 2015 et al. J. Geophys. Res. Space Physics, 120,
- 253 doi:10.1002/2014JA020786.
- 254 Zhang, Y. H. Matsumoto, and H. Kojima 1999 J. Geophys. Res. 104(A12),28633–28644,
- 255 doi:10.1029/1999JA900301.
- 256
- 257
- 258
- 259 Figure Captions



260

Figure 1. Example of electron-scale electron vortex magnetic hole observed on Nov 03 2015 in
the magnetosheath. (a) Magnetic field strength; (b) Three components of magnetic field in
LMN coordinates; (c) Three components of electron velocity in the LMN coordinates; (d)
Electron temperature (black), parallel temperature (red) and perpendicular temperature (green);
and (e) Electron density.



Figure 2. Electron temperatures inside the electron vortex magnetic holes and outside these magnetic holes observed by MMS from September 2015 to January 2016. The ratios of (a) perpendicular temperatures, (b) parallel temperatures and (c) total temperatures inside the holes to these outside the holes; (d) the ratios of perpendicular temperatures to parallel temperature inside the holes.



279

Figure 3. Example of the observations of whistler waves inside electron vortex magnetic holes on 03 November 2015. (a) Magnetic field strength; (b) The ratios between electron perpendicular temperature and parallel temperature; (c) Filtered magnetic field fluctuations from SCM instruments; (d)-(e) Power spectra densities of the electric and magnetic components of the waves; (f) Ellipticity (positive or red: right hand polarization; negative or blue: left hand polarization); (g) Propagation angle; (h) Polarization degree. The black, cyan, red curves in (dh) indicate the frequencies $0.1 f_{ce}$, $0.5 f_{ce}$, and $1.0 f_{ce}$, respectively.



287

288

Figure 4. The dispersion relations (black) and growth rates (red) of the theoretical whistlers based on the plasma parameters from local observations inside the electron vortex magnetic hole. The local plasma parameters used to calculate the growth rate of whistler waves: $|B| \sim 12$ nT, $n \sim 10$ cm⁻³, $T_i \sim 800$ eV and $T_e \sim 100$ eV, $T_{e\perp}/T_{e\parallel} = 1.4$.

293