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## Effect of cationic substitution on the pressure-induced phase transitions in calcium carbonate

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## Abstract

The high-pressure  $CaCO<sub>3</sub>$  phase diagram has been the most extensively studied within the carbon-ates group. However, both the diverse mineralogy of carbonates and the abundance of solid solutions in natural samples require the investigation of multicomponent systems at high pressures  $(P)$  and temperatures  $(T)$ . Here we studied a member of the  $CaCO<sub>3</sub>-SrCO<sub>3</sub>$  solid-solution series and revealed the effect of cationic substitution on the pressure-induced phase transitions in calcium carbonate.

A synthetic solid solution  $Ca_{0.82}Sr_{0.18}CO<sub>3</sub>$  was studied in situ by Raman spectroscopy in a diamond-anvil cell (DAC) up to 55 GPa and 800 K. The results of this work show significant differences in the high-pressure structural and vibrational behavior of the  $(Ca, Sr)CO<sub>3</sub>$  solid solution compared to that of pure  $CaCO<sub>3</sub>$ . The monoclinic  $CaCO<sub>3</sub>$ -II-type structure (Sr-calcite-II) was observed already at ambient conditions instead of the "expected" rhombohedral calcite. The stress-induced phase transition to a new highpressure modification, termed here as Sr-calcite-IIIc, was detected at 7 GPa. Sr-calcite-VII formed already at 16 GPa and room T, which is 14 GPa lower compared to  $CaCO<sub>3</sub>$ -VII. Finally, crys-tallization of Sr-aragonite was detected at 540 K and 9 GPa, at 200 K lower T than pure aragonite. Our results indicate that substitution of  $Ca<sup>2+</sup>$  by bigger cations, such as  $Sr<sup>2+</sup>$ , in CaCO<sub>3</sub> structures can stabilize phases with larger cation coordination sites (e.g., aragonite,  $CaCO<sub>3</sub>$ -VII, and post-aragonite) at lower  $P$ -T conditions compared to pure  $CaCO<sub>3</sub>$ . The present study shows that the role of cationic composition in the phase behavior of carbonates at high pressures should be carefully considered when modeling the deep carbon cycle and mantle processes involving carbonates, such as metasomatism, deep mantle melting, and diamond formation.

Keywords: Deep carbon cycle, calcium carbonate, solid solution, phase diagram, phase transition, high pressure, vibrational spectroscopy

## Introduction

Investigations of natural samples such as mantle xenoliths, carbonatites and kimberlites, inclusions in diamonds, and car-bonate-bearing rocks from the ultrahigh-pressure metamorphic terranes provide direct evidence of the existence of carbonates and carbonatite melts in the Earth's deep upper mantle, transi-tion zone, and lower mantle (Stachel et al. 2000; Brenker et al. 2007; Logvinova et al. 2008, 2011, 2019a, 2019b; Kaminsky et al. 2009, 2013; Korsakov et al. 2009; Kamenetsky and Yaxley 2015; Sharygin et al. 2015). Syngenetic inclusions in diamonds and other mantle minerals frequently contain calcium carbonates, dolomite, magnesite, and ankerite (Stachel et al. 2000; Brenker et al. 2007; Logvinova et al. 2008, 2011, 2019a, 2019b; Kaminsky et al. 2009, 2013; Korsakov et al. 2010). The abundance of the calcium carbonate at near surface and mantle conditions, and its overall geological significance have triggered an enormous interest in the CaCO<sub>3</sub> phase diagram (Suito et al. 2001; Ono et al. 2005, 2007; Merlini et al. 2012, 2018; Pippinger et al. 2014; Koch-Müller et al. 2016; Gavryushkin et al. 2017; Lobanov et 1

## al. 2017; Bayarjargal et al. 2018).

At atmospheric pressure,  $CaCO<sub>3</sub>$  crystallizes predominantly in the rhombohedral calcite structure [space group  $R3c$ ] with coplanar CO $3$ <sup>-</sup> groups and Ca<sup>2+</sup> in sixfold coordination (Ca<sup>[6]</sup>). Compression of calcite leads to a complex polymorphism with several phase transitions. Previous experimental and computa-tional studies revealed four stable high-pressure polymorphs of calcite: aragonite (Pnma, Ca<sup>[9]</sup>), CaCO<sub>3</sub>-VII (P2<sub>1</sub>/c, Ca<sup>[10]</sup>), post-aragonite (*Pmmn*, Ca<sup>[12]</sup>), and sp<sup>3</sup>-CaCO<sub>3</sub> (P2<sub>1</sub>/  $c$ , Ca<sup>[12]</sup>, C<sup>[4]</sup>) (Fig. 1) (Ono et al. 2005, 2007; Gavryushkin et al. 2017; Lobanov et al. 2017; Bayarjargal et al. 2018). The formation of these high-pressure phases is kinetically hindered at ambient temperature and occurs only above 773 K (500 °C) (Bayarjargal et al. 2018). Cold compression of calcite, on the contrary, leads to a different sequence of high-pressure polymorphs:  $CaCO<sub>3</sub>$ -II  $(P2<sub>1</sub>/c, Ca<sup>[6]</sup>)$ ; CaCO<sub>3</sub>-III and -IIIb  $(P1, Ca<sup>[7]</sup>$  and Ca<sup>[9]</sup>); CaCO<sub>3</sub>-VI (P1, Ca<sup>[7+2]</sup>) (Fig. 1) (Suito et al. 2001; Merlini et al. 2012, 2018; Koch-Müller et al. 2016; Bayarjargal et al. 2018). All of these high-pressure modifications are considered metastable (Suito et al. 2001; Merlini et al. 2012, 2018; Koch-Müller et al. 2016; Bayarjargal et al. 2018). Such consideration arises from the fact that those polymorphs appear within the stability fields of



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¢Ôæ£Ü¦DÔUHÁ [|ãáÁ [|ˇqã] Á æ Áˆ} d@ •ã ^åÁæókGhŐÚæÁe}åÁFGÏHÁSÁ  $\mathfrak{E}$ = CEE ÁO DÓ DÓ DÓ BHÓ DÓ BH Ú[oråæ{ÁQYæn∖^¦ÁroÁ<del>æpLÁK</del>JJ€DÓ¥Vo2ÁY}Á duration was 8 h. The starting material was prepared from analytical-grade synthetic powders of CaCO3 and SrCO3 (99.999% purity, Sigma Aldrich Chemical Company), with 20 mol% SrCO3 in the initial mixture. The blended mixture was placed into a Pt capsule. A Cr-doped MgO octahedron with an edge length of 18 mm and other MgObased parts were used as pressure transmitting media. Heat was generated with a stepped graphite heater. Temperature was controlled using a W5%Re-W26%Re type C thermocouple. Eight tungsten carbide cubes with the truncation edge length of 11 mm were used as Kawai-type anvils. The sample was analyzed by electron microprobe (EMP), transmission electron microscopy (TEM), powder X-ray diffraction (XRD), Fourier transform infrared (FTIR), and Raman spectroscopy.

The chemical composition of the synthesized sample (Online Materials<sup>1</sup> Table OM1 and Fig. OM2) was determined with EMP analyses using a JEOL Hyperprobe JXA-8500F EMP with a field emission cathode (GFZ, Potsdam). Analysis was conducted at a 15 kV acceleration voltage, a 10 nA beam current, and a 0.5–10 μm beam size. Dolomite and strontianite were used as standards.

TEM with a high-resolution energy-dispersive spectrometer was used for the microtexture observation and preliminary structural analysis (Fig. 2; Online Materials<sup>1</sup> Fig. OM3). Thin sections of approximately 15 x 10 x 0.15  $\mu$ m<sup>3</sup> size were prepared with a<br>focused Ga-ion beam (FIB) system (FEI FIB 200 TEM) (Wirth 2009) directly from the crystals in the microprobe mounts, and analyzed using a FEI Tecnai G2 F20

#### X-Twin transmission electron microscope (GFZ, Potsdam).

The structural analysis was performed with XRD. The XRD patterns were measured at atmospheric pressure using a STOE Stadi P diffractometer equipped with a curved Germanium (111) primary monochromator, a high-resolution MYTHEN-detector and a normal focus Cu X-ray tube (Cu<sup>K</sup>α1 radiation) (GFZ, Potsdam). The XRD data were processed with the GSAS II software package (Toby and Von Dreele 2013). Unit-cell parameters were obtained by Le Bail/Pawley refinements (Fig. 3).

FTIR measurements in the mid-infrared region were conducted with the Ver-tex 80v FTIR spectrometer combined with a Hyperion 2000 microscope (GFZ, Potsdam) within a 500-2000  $cm^{-1}$  spectral range (Fig. 4; Online Materials<sup>1</sup> Fig. OM4). We used a KBr beamsplitter and a mercury-cadmium-telluride detector. The spectra were averaged over 1536 scans with a spectral resolution of 2 cm<sup>-1</sup>. The description of the Raman measurements is given in the section below.

#### High pressure in situ experiments

Pressure-induced phase transitions in the Ca0.82Sr0.18CO3 solid solution were studied in situ using Raman spectroscopy at room and at high temperatures up to 800 K (Fig. 1; Online Materials<sup>1</sup> Fig. OM1). Two different types of diamond-anvil

cells (DACs) were employed, depending on the target temperature.

Symmetric piston-cylinder DACs with 250 μm culet size diamonds were used for the room-temperature experiments. Rhenium gaskets were indented to a thickness of 30–40 μm and drilled in the center of the indentation to form a sample chamber. The crystals were carefully selected, cleaned with ethanol, and placed into the sample chamber together with ruby spheres acting as pressure sensors. Liquid argon (Ar) was loaded cryogenically as a pressure-transmitting medium, follow-ing the same procedure as in Koch-Müller et al. (2016). Pressure was measured using laser-induced fluorescence spectroscopy of ruby  $(Al_2O_3:Cr^{3+})$  (Dewaele et al. 2008). Previous studies have shown that the ruby scale is accurate within 2–5% up to 55 GPa (Dewaele et al. 2008). Four different experimental high-pressure runs were conducted at room temperature (Fig. 1; Online Materials<sup>1</sup> Fig. OM1). Raman spectra were measured in 1–2 GPa steps up to 55 GPa, both on compression and decompression cycles (Online Materials<sup>1</sup> Fig. OM1).

The second set of experiments was performed using an internally heated membrane-driven DAC (DAC-HT) [Diacell μScopeDAC HT (G), EasyLab, U.K.], equipped with diamonds of 300 μm culet size. The measurements were carried out in the 9–10 GPa pressure range and at high temperatures from 540 to 800 K (Fig. 1). The sample, together with  $SrB<sub>4</sub>O<sub>7</sub>:Sm<sup>2+</sup>$  powder and ruby spheres used as P-T sensors, was sandwiched between two NaCl layers and placed inside the sample chamber. NaCl served both as a pressure medium and as a thermal insula-tor. The temperature in the DAC-HT experiments was measured both outside the sample chamber with a K-type thermocouple placed close to the lower diamond, as well as inside the sample chamber using the fluorescence lines of  $SrB_4O_7$ :  $Sm^{2+}$  and ruby (Datchi et al. 2007). Details of the temperature calibrations are given in Online materials. The  $P$ -T calibration method allows to determine pressure with a 0.4–1.2% and temperature with a 1.0–1.3% accuracy up to 20 GPa and 900 K (Datchi et al. 2007; Romanenko et al. 2018). The temperature outside the gasket hole, measured by the thermocouple, was always lower than that estimated inside the sample chamber, with the difference reaching about 50 K at the maximum temperature of 800 K.

The Raman spectra in both room and high- $T$  experiments were measured with a HORIBA Jobin Yvon LabRAM HR800 VIS spectrometer (GFZ, Potsdam) equipped with a blue 473 nm diode-pumped solid-state laser. The spectral range for the measurements was 140–1200 cm<sup>-1</sup> and the data acquisition time was 120 s. The spectra were collected at ambient conditions before and after the pressure increase and at high pressures both upon compression and decompression (Online Materials<sup>1</sup> Fig. OM1). The measurements at high temperature were performed every 2–5 min. The software Fityk (Wojdyr 2010) was used for data analysis.



Figure 2. Results of the TEM analysis performed on the Ca0.82Sr0.18CO3 solid solution: (a) the lamellar texture of the sample, (b) the electron diffraction pattern. The lines in b represent the unit cell of calcite with R3c space group with  $a = 5.01$  Å and  $c =$ 17.25 Å. The EDX spectra are given in Online Materials<sup>1</sup> Figure OM3.



Figure 3. The results of the Pawley refinement for the Ca0.82Sr0.18CO3 solid solution (Sr-calcite-II) using (a and b) calcite ( $R3c$ ) and (c and d) CaCO<sub>3</sub>-II ( $P21/c$ ) as a model. The XRD patterns were obtained at ambient conditions. Insets on top (b and c) show magnified region with splitted peaks. The calculated Bragg reflection positions of CaCO3-II and calcite are marked with red vertical bars; experimental patterns  $=$  dark blue crosses, calculated patterns = green line; difference profiles = blue line; and background profiles = red line. (Color online.)

#### Experimental results

#### Characterization of the synthesized (Ca,Sr)CO3 solid solution

Backscattered electron images of the synthesized samples are presented in Online Materials<sup>1</sup> Figure OM2. The recovered sample consists of large (≥30–100 μm) grains of (Ca,Sr)CO<sub>3</sub> solid solution with homogeneous composition and Sr<sup>2+</sup> content of 18 mol% according to the EMP analyses (Online Materials<sup>1</sup> Table OM1). Some of the grains contain small inclusions ≤1–2 μm in size, located primarily in voids and cracks (Online

Materials<sup>1</sup> Fig. OM2b).<br>Two thin films from different grains were prepared for the TEM analysis by the focused ion beam method (Fig. 2; Online Materials<sup>1</sup> Fig. OM3). The samples have a high density of dislocations

and a lamellar texture with clear interfaces caused by complex twining (Fig. 2a). Analytical TEM confirmed that both samples have homogeneous compositions and consist predominantly of crystalline Ca-rich (Ca~80 mol%) (Ca,Sr)CO3 solid solution (Online Materials<sup>1</sup> Figs. OM3a and OM3c). The electron diffraction pattern of the main phase was refined in the rhombohedral calcite unit cell; however, we see several additional low-intensity reflections (Fig. 2b). The few grains of the secondary Sr-rich (Sr, Ca)CO3 phase, which constitutes less than 1 vol% of the thin film, were detected as small inclusions along the dislocations and cracks (Online Materials<sup>1</sup> Figs.<br>OM3a and OM3b). The electron diffraction pattern of the secondary phase could be refined with orthorhombic symmetry.

Most of the reflections in the XRD pattern of the Ca0.82Sr0.18CO3 solid solution, including the one with the highest intensity, can be refined in the calcite structure (Figs. 3a and 3b) with lattice parameters  $a = 5.01$  Å,  $c = 17.25$  Å (residual factor  $wP_{\text{D}} = 19\%$ ). However, several low-intensity peaks located at  $2\theta = 23$ , 36, 47, and 48° show a splitting incompatible with the calcite structure (Figs. 3a and 3b). The attempts to include an orthorhombic phase, as observed in the TEM analyses, did not improve the XRD refinement. It is likely that due to the low concentration, the secondary phase does not appear in the XRD pattern. The best fit for all of the observed Bragg peaks was achieved using a monoclinic unit cell similar to that of CaCO<sub>3</sub>-II ( $wR<sub>D</sub>$  = 8.3%, Figs. 3c and 3d) (Merrill and Bassett 1975). The refined unit cell of Can  $82$ Srn  $18$ CO<sub>3</sub> solid solution has  $P21/c$  space group with the following lattice parameters:  $a = 6.44 \text{ Å}$ ;  $b = 5.02 \text{ Å}$ ;  $c = 8.13$ Å;  $β = 108.13°$ . For comparison, the lattice parameters of pure CaCO3-II are  $a = 6.33$  Å;  $b = 4.95$  Å;  $c = 8.03$  Å;  $\beta = 107.9^{\circ}$  at 1.5 GPa (Merrill and Bassett 1975).

The difference in the structure of the synthesized solid solution from that of calcite was also indicated by the results of the mid-infrared (MIR) and Raman measurements (Figs. 4 and 5; Online Materials<sup>1</sup> Fig. OM4–OM6). The vibrational spectrum of calcite contains 4 fundamental internal modes, assigned to CO3<sup>2</sup> vibrations: symmetric stretching (ν<sub>1</sub>), out-of-plane bend-<br>ing (ν2), asymmetric stretching (ν3), and in-plane bending (ν4) (White 1974) (Fig. 4; Online Materials<sup>1</sup> Fig. OM5). The v<sub>1</sub> and  $v_2$  modes are Raman- (v<sub>1</sub>) or infrared- ( $v_2$ ) active only, while the ν3 and ν4 bands can be detected in both Raman and IR spectra (White 1974).

The fundamental modes in the MIR-spectrum of pristine CaCO<sub>3</sub> calcite are located at 872 cm<sup>-1</sup> (v<sub>2</sub>), 1407 cm<sup>-1</sup> (v<sub>3</sub>), and 712 cm<sup>-1</sup> (v<sub>4</sub>) (Fig. 4) (White 1974; Vahur et al. 2016). Addi-tionally, spectra can exhibit a satellite mode at 1660  $cm^{-1}$ and a combination band at 1750 cm<sup>-1</sup> (ν<sub>1</sub>+ν4). The MIR spectrum of the studied Ca0.82Sr0.18CO3 solid solution differs from that of calcite, as both of the v2 and v4 bending modes split into doublets separated by  $\sim$ 10  $\rm cm^{-1}$ , and additional modes appear at 1085 cm<sup>-1</sup> (v<sub>1</sub>), 859 cm<sup>-1</sup>, and at 1740 cm<sup>-1</sup> (Fig. 4; Online Materials<sup>1</sup> Fig. OM4). The appearance of these new modes, in particular the IR-forbidden v<sub>1</sub> band and the apparent splitting of the bending modes, clearly indicate a lower crystalline symmetry of the solid solution compared to pure calcite, at least on the local scale (White 1974). Similar IR features were previously observed in the infrared spectra of CaCO3-II, collected in situ at 1.9 GPa (Koch-Müller et al. 2016), and in barytocalcite CaBa(CO3)2 (P21/m) (Scheetz and White 1977). Both of these phases have structures derivative of calcite.

Raman spectra of Ca0.82Sr0.18CO3 were collected at atmo-spheric pressure in the range of 100–1250  $\text{cm}^{-1}$  (Fig. 5; Online Materials<sup>1</sup> Figs. OM5-OM6). The following modes are assigned to internal vibrations of the  $(\text{CO}_3)^{2-}$  group: 1087 cm<sup>-1</sup>  $\frac{1}{2}$  with a sat-ellite mode at 1066 cm<sup>-1</sup> (v<sub>1</sub>), 711 cm<sup>-1</sup> with a shoulder band at



Figure 4. Mid IR spectrum of the Ca0.82Sr0.18CO3,Srcalcite-II (Sr-CC-II), in comparison with spectrum of CaCO3 calcite [from IR database (Vahur et al. 2016)]. The spectrum of calcite contains 4 modes (shown by gray dashed lines): 712 cm– <sup>1</sup> (ν<sub>4</sub>); 872 cm<sup>-1</sup> (ν<sub>2</sub>); 1407 cm<sup>-1</sup> (ν<sub>3</sub>); and 1750 cm<sup>-1</sup> (ν<sub>1</sub>+ν<sub>4</sub>).<br>In the MIR spectrum of the Sr-CC-II the ν<sub>2</sub> and ν<sub>4</sub> bending modes are splitted, additional modes appear at 1085  $cm^{-1}$  (v<sub>1</sub>), 859 cm<sup>-1</sup>, and at 1740 cm<sup>-1</sup> (all marked by arrows). The fitting of the bands of the Ca0.82Sr0.18CO3 spectrum is illustrated in Online Materials<sup>1</sup> Figure OM4. (Color online.)

706 cm<sup>-1</sup> (v<sub>4</sub>), and a third peak at 875 cm<sup>-1</sup> (v<sub>2</sub>) (Online Materials<sup>1</sup> Figs. OM6 and OM7). Three main low-frequency bands located at 148, 190, and  $275 \text{ cm}^{-1}$  correspond to lattice vibrations (Online Materials<sup>1</sup> Fig. OM6). All bands, with the exception of the 706, 190, and 875  $\text{cm}^{-1}$  can be assigned to calcite, yet shifted to lower wavenumbers due to the expansion of the unit cell caused by the presence of the larger  $Sr<sup>2+</sup>$  cations (Online Materials<sup>1</sup> Fig. OM5). The low-intensity mode at 875  $\text{cm}^{-1}$  (Online Materials<sup>1</sup>



Figure 5. Raman spectra of Ca0.82Sr0.18CO3 solid solution collected upon compression up to 7 GPa. Four high-P phases were observed: Sr-calcite-II (Sr-CC-II), Sr-calcite-IIIb (Sr-CC-IIIb), Sr-calcite-III (Sr-CC-III), and Sr-calcite-IIIc (Sr-CC-IIIc). Raman spectrum collected before the experiment is shown in black. For more details see run 3 in Online Materials<sup>1</sup> Figure OM1. (Color online.)

Fig. OM6) corresponds to the out-of-plane bending vibration (ν2), a Raman-inactive vibration in the calcite structure (White 1974). As in the MIR case, the observation of additional modes indicates a lower symmetry of Ca0.82Sr0.18CO3 compared to that of rhombohedral calcite.

Previously proposed as a main feature of the CaCO<sub>3</sub>-II polymorph (Pippinger et al. 2014) was the appearance of the additional lattice vibration mode at 190 cm–1 and the clear splitting of the ν4 band in the Raman spectra. It should be noted, however, that while 190 cm<sup>-1</sup> could be unequivocally distinguished in the spectra of Ca0.82Sr0.18CO3, the splitting of ν4 band is not appar-ent, and the "extra" component appears as a shoulder (Online Materials<sup>1</sup> Fig. OM6). Possible reasons behind the absence of a clear ν4 band splitting might be the spectral resolution, as well as the structural differences between Ca0.82Sr0.18CO3 and CaCO3-II at a microstructural scale.

Overall, all of our collected data allow us to conclude that the structure of the synthesized Ca0.82Sr0.18CO3 phase at ambient conditions has a lot of similarities with that of CaCO<sub>3</sub>-II and can be viewed as a monoclinic distortion of the rhombohedral calcite structure. Therefore, we will call this new phase Srcalcite-II.

### High-pressure Raman study

The results of the in situ high-pressure study of the Ca0.82Sr0.18CO3 solid solution are summarized in Figures 5–7

# and in Online Materials1 Figs. OM7–OM13. Since the Raman spectra of the detected high-pressure polymorphs are similar to

those of CaCO3 high-pressure modifications, the same nomenclature is used. The Raman spectra collected in the compression experiments at room temperature are shown in Figures 5 and 6. The first phase transition was detected at 1.8 GPa. New bands

appear at 1080  $cm^{-1}$  (Fig. 5; Online Materials<sup>1</sup> Fig. OM7) and below 900 cm–1, where we observe at least 10 new bands

 $(Fig, 5;$  Online Materials<sup>1</sup> Fig. OM9). The arising features of the Raman spectra are in good agreement with previous studies on the pure CaCO3 system (Pippinger et al. 2014; Koch-Müller et al. 2016), resembling the CaCO3-IIIb Raman pattern. Thus, the Ca0.82Sr0.18CO3 solid solu-



Figure 6. Raman spectra collected at 14-50 GPa (see runs 3 and 4 in Online Materials<sup>1</sup> Fig. OM1) showing formation of Sr-calcite-VII (Sr-CC-VII) and Sr-post-aragonite (Post-Arag) in the room-temperature experiments. The Sr-post-aragonite band appears around 44.6 GPa (marked by arrows). (Color online.)



Figure 7. Transformation of Sr-calcite-IIIc (Sr-CC-IIIc) to Sr-calcite-IIIb (Sr-CC-IIIb) and Sr-aragonite (Arag) at high temperature and  $9-10$  GPa.  $P$ -T conditions and time from the beginning of heating  $(t)$  at which spectrum was collected are shown on the figures. The characteristic Sr-aragonite peak is marked by the arrows. (Color online.)

tion with CaCO3-IIIb-like structure will be called Sr-calcite-IIIb. Srcalcite-IIIb was detected in the pressure range of 1.8–4 GPa

(Fig. 1; Online Materials1 Fig. OM1). Increase of pressure above 4 GPa leads to a clear and abrupt splitting of the 1100  $\text{cm}^{-1}$  band in the Raman spectra (Fig. 5; Online Materials1 Figs. OM7 and OM8a). Previous studies on the CaCO3 system attributed this splitting to the formation of CaCO3-III (Pippinger et al. 2014; Koch-Müller et al. 2016; Bayarjargal et al. 2018). Consequently, we conclude that Srcalcite-III has appeared above 4 GPa. This phase was detected in a relatively narrow pres-sure range from 4 to 7 GPa (Fig. 1; Online Materials<sup>1</sup> Fig. OM1).

Another clear change in the Raman spectra takes place above 7 GPa (Fig. 5; Online Materials<sup>1</sup> Figs. OM7–OM9). Even though the spectra above 7 GPa contain a similar set of modes as that of Sr-calcite-III, the bands above  $1090 \text{ cm}^{-1}$  broaden and shift abruptly to higher wavenumbers (Fig. 5; Online Materials<sup>1</sup> Fig. OM7). The observed changes indicate the appearance of another structural modification. However, the Raman spectra cannot be explained by any known highpressure polymorphs of CaCO3 (Koch-Müller et al. 2016; Bayarjargal et al. 2018). Based on the similarities in the Raman spectra (Fig. 5; Online Materials<sup>1</sup> Figs. OM7–OM9), we assume that this new high-pressure modification has a closely related structure to that of Sr-calcite-III. Hence, we will refer to it as Srcalcite-IIIc.

Sr-calcite-IIIc was detected as the only phase in the Raman spectra at room temperature from 7 to 14 GPa (Fig. 6). Further compression of Sr-calcite-IIIc leads to the formation of another high-pressure polymorph between 14–18 GPa (Fig. 6). The transition is sluggish and the sample transforms fully to the highpressure phase at 20 GPa (Fig. 6). The Raman spectra collected at 20 GPa contain 13 intense bands in the range of 150–1250 cm<sup>-1</sup> (Fig. 6; Online Materials<sup>1</sup> Figs. OM10 and OM11). The external vibrational modes located below 500 cm–1 include at least 7 bands (Fig. 6; Online Materials<sup>1</sup> Fig. OM11). The internal vibrational modes are represented by a single band at 1145 cm–  $1$  (v<sub>1</sub>) with a

shoulder on the high-frequency side, and two doublets at 733, 745 cm<sup>-1</sup> (v4), and 862, 868 cm<sup>-1</sup> (v<sub>2</sub>) (Fig. 6; Online Materials<sup>1</sup> Figs. OM10 and OM11).

In pure CaCO3 all previous studies showed the formation of CaCO3-VI above 16 GPa at room temperature (Koch-Müller et al. 2016; Bayarjargal et al. 2018). However, the Raman spectra of the Ca<sub>0.82</sub>Sr<sub>0.18</sub>CO<sub>3</sub> solid solution show major differences from those of CaCO3-VI (Online Materials<sup>1</sup> Fig. OM12). Raman spectra of CaCO3-VI contain a single v2 band (850 cm<sup>-1</sup> at 20<br>GPa) (Bayar-jargal et al. 2018), while the appearance of the two distinct bands at 850–890  $cm^{-1}$ , as observed in the present study, was reported as the main feature of the CaCO3-VI to -VII transition (Bayarjargal et al. 2018) (Online Materials<sup>1</sup> Fig. OM12). Based on our observa-tions, we conclude that in the Ca0.82Sr0.18CO3 solid solution, we observe the formation of a polymorph with a structure similar to CaCO3-VII (Sr-calcite-VII).

The Raman spectra of Ca0.82Sr0.18CO3 solid solution measured between 20–55 GPa showed the persistence of Srcalcite-VII up to the highest experimental pressure (Figs. 1 and 6; Online Ma-terials<sup>1</sup> Fig. OM1). A new feature arising at 810 cm<sup>-1</sup> at 44.6 GPa (Fig. 6; Online Materials<sup>1</sup> Fig. OM10) may indicate the onset of a transition to a post-aragonite phase (Srpost-aragonite further below) (Bayarjargal et al. 2018). Srcalcite-VII and Sr-post-aragonite coexist between 44.6–55 GPa at room temperature (Fig. 6; Online Materials<sup>1</sup> Fig.<br>OM10). Formation of both phases at ambient temperature contrasts previous data on pure CaCO3, where the respective polymorphs were detected only after heating (Gavryushkin et al. 2017; Bayarjargal et al. 2018).

In the experimental runs 1 and 2 (Fig. Online Materials<sup>1</sup> Fig. OM1), Raman spectra were also measured on decompression. The decompression of Sr-calcite-VII led to the formation of Sr-calcite-IIIc together with Sr-calcite-IIIb below 16 GPa (Online Materials<sup>1</sup> Figs. OM1, OM8b, and OM8c). Thus,Sr-calcite-IIIb could be observed in two pressure ranges from 1.8 to 4 GPa, and from 7 to 14 GPa (Online Materials<sup>1</sup> Fig. OM1), same as reported by Koch-Müller et al. (2016) for the CaCO3-IIIb. However, the "pressure stability window" of Sr-calcite-IIIb depends on the experimental path.

For better understanding of the Sr-calcite-IIIc and -IIIb behavior, additional high-temperature Raman experiments were carried out at 9–10 GPa and 540–800 K (Fig. 1). First, pressure was increased to 9.4 GPa and Raman spectra were measured before heating (Fig. 7; Online Materials<sup>1</sup> Fig. OM13). Sr-calcite-IIIc<br>was detected at room temperature, confirming the previous observa-tions. After <1 min of heating, Sr-calcite-IIIc fully transformed to Sr-calcite-IIIb (Fig. 7; Online Materials<sup>1</sup> Fig.<br>OM13). Further heating led to the formation of an aragonite-type polymorph (Sr-aragonite) already at 540 K and 9 GPa (Fig. 7; Online Materials<sup>1</sup> Fig. OM13). Sr-aragonite is quenchable to atmospheric pressure and was detected in the Raman spectra collected after experiments (Online Materials<sup>1</sup> Fig. OM13).

#### **Discussion**

In this paper, we present the results of the first in situ experimental study on the CaCO3-SrCO3 system, at pressures up to 55 GPa (Fig. 1; Online Materials<sup>1</sup> Fig. OM1). The results of the present study show that incorporation of  $Sr<sup>2+</sup>$  in CaCO<sub>3</sub> has a significant effect on the phase behavior of calcium carbonate at high pressures both at room and at high temperature. Below

we discuss the findings of this work in detail, as well as their implications for understanding the structural behavior of carbonates at mantle conditions.

The first synthesis experiment was performed at 2 GPa and 1273 K (1000 °C) (Fig. 1), which according to the CaCO3 phase diagram, corresponds to the CaCO3-IV stability field, close to the calcite-aragonite transition curve (~2.5 GPa, 1273 K) (Ter Heege and Renner 2007). The choice of the P-T synthesis parameters was influenced by the fact that a miscibility gap exists in the CaCO3-SrCO3 system at low pressures (Chang and Brice 1972; Carlson 1980). Due to the difference in the ionic radii of  $Ca<sup>2+</sup>$  (1.00 Å) and Sr<sup>2+</sup> (1.18 Å) (Shannon 1976), miscibility of  $Sr<sup>2+</sup>$  in the calcite structure is limited to 15 mol% at room temperature (Chang and Brice 1972; Carlson 1980; Matsunuma et al. 2014) and  $~40$  mol% at 973 K (700 °C) (Chang and Brice 1972; Carlson 1980). The coexisting phases in the miscibility gap have orthorhombic aragonite-type and rhombohedral calcite-type structures on the Sr-rich and Ca-rich sides, respectively (Chang and Brice 1972; Carlson 1980). The two-phase field exists up to the transition to the CaCO<sub>3</sub>-IV anion-disordered phase  $(R3m,$ disordered calcite structure) at 973–1173 K (700–900 °C) (Chang and Brice 1972; Carlson 1980).

In this study we were able to synthesize a Ca0.82Sr0.18CO3 solid solution at 2 GPa and 1273 K (1000 ° C). The results of TEM (Fig. 2; Online Materials<sup>1</sup> Fig. OM3) and EDS (Online Materials<sup>1</sup> Fig. OM2) studies show textural and compositional homogeneity of the solid solution. The minor Sr-rich secondary phase was detected as small inclusions along the dislocations and cracks, which allows us to conclude that it is likely a relic of the growth media.

According to the XRD results, Ca<sub>0.82</sub>Sr<sub>0.18</sub>CO<sub>3</sub> is monoclinic at 2 GPa (Sr-calcite-II, space group  $P21/c$ ). MIRand Raman spectroscopy results show that Sr-calcite-II has a structure similar to that of CaCO<sub>3</sub>-II, which has been described as a monoclinic distortion of the calcite structure. CaCO3-II is a metastable phase that crystallizes in a narrow pressure range between 1.7 and 2.5 GPa (Pippinger et al. 2014). The majority of the experimental studies show that it is not quenchable to atmospheric pressure and transforms to aragonite upon heating (Merrill and Bassett 1975; Pippinger et al. 2014; Koch-Müller et al. 2016; Bayarjar-gal et al. 2018). However, based on our

observations, we can conclude that incorporation of  $Sr<sup>2+</sup>$  in the solid solution expands the stability field of the CaCO3-II-like structure (Sr-calcite-II) to higher temperatures and lower pressures. Moreover, Sr-calcite-II is quenchable to ambient conditions.

As for the effect of cold compression, Sr-calcite-II was observed at room temperature up to 1.8 GPa. Further increase of pressure led to the formation of the following phases: (1) Sr-calcite-IIIb at 1.8–4 GPa; (2) Sr-calcite-III at 4–7 GPa;(3) Sr-calcite-IIIc and/or IIIb at 7–16 GPa; (4) Sr-calcite-VII 16–55 GPa; and (5) Sr-post-aragonite above 45 GPa (Fig. 1; Online Materials<sup>1</sup> Fig. OM1).

The Raman spectra of the Sr-calcite-IIIb, -III, and -IIIc share a lot of similarities (Online Materials<sup>1</sup> Figs. OM8 and OM9), and differ mostly in the symmetric stretching vibration regions (900–  $1200 \text{ cm}^{-1}$ ), composed of a single band with a shoulder for the Srcalcite-IIIb modification, and doublets for the Sr-calcite-III and -IIIc (Online Materials<sup>1</sup> Figs. OM8 and OM9). Comparison of

our spectroscopic results with previous publications (Koch-Müller et al. 2016; Bayarjargal et al. 2018) shows that all phases have structures closely related to CaCO3-III and -IIIb. Moreover, the formation of Sr-calcite-IIIb and -III phases in the solid solution occurs at the same  $P$ - $T$  conditions where the respective CaCO<sub>3</sub> polymorphs, CaCO<sub>3</sub>-IIIb and -III also appear (Fig. 1; Online Materials<sup>1</sup> Fig. OM1) (Pippinger et al. 2014; Koch-Müller<br>et al. 2016; Bayarjargal et al. 2018). The third phase, Sr-calcite-IIIc, is a unique modification, not observed in CaCO3.

CaCO3-IIIb and -III structures are topologically similar and contain non-co-planar  $CO3^{2-}$  groups and two non-equivalent Ca coordination polyhedra (Ca[7] and Ca[9]) (Merlini et al. 2012). The main difference between them is in the arrangement of the structural blocks, which results in 10 (CaCO3-III) and 4 (CaCO3-IIIb) formula units, respectively, in the primitive unit cell (Mer-lini et al. 2012). According to DFT calculations, the increased number of atoms in the unit cell of CaCO3-III as compared to CaCO3-IIIb results in the splitting of the ν1, ν2, and ν4 bands in the CaCO3-III spectra (Koch-Müller et al. 2016). Taking into account the previous results on CaCO<sub>3</sub>-III and -IIIb, and the similarities of the Raman spectra of Sr-calcite-IIIb, -III, and -IIIc, we could speculate that those high-pressure modifications have as well topologically similar structures (Koch-Müller et al. 2016; Bayarjargal et al. 2018). However, as the Raman spectroscopy does not give definite structural solutions, further single-crystal XRD studies are necessary.

The similarity of Sr-calcite-IIIb and -IIIc structures could explain the observed behavior of those phases at high pressure and temperature. Formation of Sr-calcite-IIIc or/and reappearance of Sr-calcite-IIIb occurs between 7 and 16 GPa and depends on the experimental  $P$ -T path (Fig. 1; Online Materials<sup>1</sup> Fig. OM1). Sr-calcite-IIIc forms in the cold compression experiments, while heating transforms it to Srcalcite-IIIb. In the decompression experiments at room temperature, Sr-calcite-IIIc and IIIb coexist in the 7–16 GPa pressure range (Online Materials<sup>1</sup> Fig. OM1)

The observed dependence of the phase behavior on the experimental  $P$ -T path may result from the sensitivity of the Sr-calcite-IIIb and -IIIc formation to the influence of the hydrostatic conditions. Both pressure media, Ar and NaCl, used in this study are non-hydrostatic at high pressure and room temperature (Klotz et al. 2009). Non-hydrostaticity may lead to the stress-induced phase transition at 7 GPa from Sr-calcite-III to -IIIc. On the contrary, thermal annealing and/or pressure decrease can reduce the stress gradient (Angel et al. 2007) and promote formation of the Sr-calcite-IIIb polymorph at similar pressures. Srcalcite-IIIb under the quasi-hydrostatic conditions forms in two pressure ranges 1.8–4 and 7–16 GPa, showing similar behavior to CaCO3-IIIb (Koch-Müller et al. 2016).

The similar effect of the pressure medium and the experimental  $P$ -T path on the CaCO3-IIIb formation was previously observed in several experimental studies (i.e., Merlini et al. 2012, 2018; Yuan et al. 2018). In pure CaCO3, non-hydrostaticity suppresses the formation of CaCO3-IIIb and only CaCO3-III could be detected. The addition of  $Sr^{2+}$  to the system makes the effect more pronounced and leads to the new structural modification Sr-calcite-IIIc.

Sr-calcite-IIIb and -IIIc are most likely metastable phases in the range of 7–16 GPa. They transform to the aragonite-structured

polymorph, Sr-aragonite, already at 540 K and 9 GPa. The formation of Sr-aragonite is observed at ~200 K lower temperatures than in CaCO<sub>3</sub> (Fig. 1).

Another high-pressure polymorph, Sr-calcite-VII appears already at room temperature and at 16 GPa, and it is stable until the highest-pressure point investigated. Above 45 GPa Srcalcite-VII probably coexists with the Sr-post-aragonite phase. The stability fields of both Sr-calcite-VII and Sr-post-aragonite are located at ~10–15 GPa lower pressures than those of the respective CaCO3 polymorphs (Ono et al. 2005; Gavryushkin et al. 2017; Bayarjargal et al. 2018). Both the formation of Sraragonite at much lower temperatures, as well as the appearance of Sr-calcite-VII and Sr-post-aragonite in the cold compression experiments, reveal a significant effect of  $Sr<sup>2+</sup>$  on the kinetics of the pressure-induced structural transformations.

Even though the formation of the CaCO3 polymorphs aragonite, CaCO3-VII, and post-aragonite is thermodynamically favored in the pure CaCO3 system, it is kinetically hindered and occurs in experiments only above 773 K (500 °C) (Fig. 1) (Gavryushkin et al. 2017; Bayarjargal et al. 2018). On cold compression and at low temperatures, CaCO<sub>3</sub> transforms to metastable polymorphs, i.e., CaCO3-IIIb instead of aragonite at 7–15 GPa, and CaCO3-VI instead of CaCO3-VII and post-aragonite at 15–50 GPa (Koch-Müller et al. 2016; Bayarjargal et al. 2018).

The introduction of  $Sr<sup>2+</sup>$  to CaCO3 seems to lower the kinetic barrier that prevents the formation of aragonite, CaCO3-VII, and post-aragonite in the end-member CaCO3, and promotes their formation at lower  $P$ - $T$  conditions. The observed tendency is most likely related to the capacity of the different Ca-carbonate high-pressure polymorphs to accommodate large cations. The known structures of aragonite, CaCO3-VII, and post-aragonite have larger cationic coordination numbers-9, 10, and 12, respectively—than CaCO<sub>3</sub>-IIIb (Ca<sup>[7]</sup>-Ca<sup>[9]</sup>) and CaCO<sub>3</sub>-VI (Ca<sup>[7</sup> <sup>+2]</sup>) (Merlini et al. 2012; Gavryushkin et al. 2017). Ca-carbonate solid solutions with  $Sr<sup>2+</sup>$  favor likely structures with larger coordination sites at lower pressures compared to pure CaCO3.

This tendency is in good agreement with previous observations in both end-members (Wang et al. 2015) and the CaCO3- SrCO3 solid solution (Carlson 1980). Strontianite,SrCO3, tends to form structures with larger coordination numbers at a given pressure, compared to CaCO3: aragonite with Sr<sup>[9]</sup> at atmospheric pressure, and post-aragonite with Sr<sup>[12]</sup> already at 20 GPa (Wang et al. 2015). The small amounts of Sr<sup>2+</sup> in CaCO<sub>3</sub> aragonite extend the stability field of aragonite structure to slightly lower pressures than in the unary CaCO3 system (Carlson 1980). The previous high-pressure studies demonstrated the formation of aragonite structured Ca<sub>0.80</sub>Sr<sub>0.20</sub>CO<sub>3</sub> solid solution at 1.6 GPa and 923 K, which in pure CaCO3 system would still belong to the stability field of calcite (Carlson 1980).

Our experiments show a significant role of the cation in the high-pressure phase transitions of calcium carbonate. The relatively small amount of  $Sr^{2+}$  in the synthesized solid solution, 18 mol%, comparable with that found in some diamond inclusions (Logvinova et al. 2008, 2011, 2019a), leads to significant changes in the phase diagram of calcium carbonate: (1) stabilization at high temperatures of the CaCO3-II-type structure (Sr-calcite-II); (2) stress-induced phase transition to the new high-pressure modification, Sr-calcite-IIIc; (3) formation of

a polymorph with aragonite type structure at lower temperatures compared with that in CaCO3; and (4) phase transition to Srcalcite-VII and Sr-post-aragonite already at room temperature and at pressures 10-15 GPa lower than in CaCO3 (Fig. 1; Online Materials<sup>1</sup> Fig. OM1).

The formation of monoclinic Sr-calcite-II at high temperatures raises the important question on whether the metastable CaCO3 polymorphs can become thermodynamically stable in the other calcium carbonate solid solutions as well. For instance, previous studies reported that addition of 5 mol% of  $\rm Fe^{2+}$  or Ma $^{2+}$  to CaCO3 stabilizes the calcite structure at 6 GPa and 1273–1473 K (1000–1200 °C) in the  $P$ -T stability field of CaCO3 aragonite (Shatskiy et al. 2014, 2018; Müller et al. 2017). It is, however, unknown what structures Mg- or Fe-calcites adopt at high pressures and temperatures, as all previous studies examined the samples ex situ (Shatskiy et al. 2014, 2018; Müller et al. 2017).

One of the possibilities, proposed by Merlini et al. (2012), is the stabilization of CaCO<sub>3</sub>-III-like structures in solid solutions with smaller cations. Due to the different polyhedral site volumes  $(Ca^{[7]}-Ca^{[9]})$ , the structure of CaCO3-III could host cations with different ionic radii without a major elastic strain from non-ideal solid solution. The results of the present study show the fundamental possibility of the stabilization of metastable CaCO3 phases in calcium carbonate solid solutions. Since CaCO3-III transforms back to calcite with the release of pressure, it is necessary to study CaCO3-MgCO3 and CaCO3-FeCO3 in situ at high pressures.

#### Implications

The present study and previous observations (Shatskiy et al. 2014, 2018; Müller et al. 2017) show that the high-pressure phase behavior of the calcium carbonate-based solid solutions depends strongly on the specific chemical composition of the solid solution. Even substituting a small amount of  $Ca<sup>2+</sup>$  by different cations, such as smaller cations ( $Mq^{2+}$  or Fe<sup>2+</sup>), or bigger ones  $(Sr<sup>2+</sup>$  or  $Ba<sup>2+</sup>$ ), shifts phase transition pressures and temperatures and probably also leads to the formation of new crystal structures.

Due to the large coordination sites in aragonite, which is long believed to be the major calcium carbonate polymorph in the upper mantle, the solubility of Mg<sup>2+</sup> and Fe<sup>2+</sup> in the structure is limited to few molar percentages (Shatskiy et al. 2014, 2018; Müller et al. 2017), and higher concentrations will promote crys-tallization of phases with smaller cationic sites either calcite, CaCO<sub>3</sub>-III, or others that are currently unknown. Indeed, natural aragonites contain only a few molar percentages of Mg<sup>2+</sup>and Fe<sup>2+</sup>, while there are findings of Mg-rich and Fe-rich calcium carbon-ates with calcite structure in polycrystalline carbonate inclusions in garnets from diamond grade metamorphic rocks (Korsakov et al. 2010; Frezzotti et al. 2011). Thus, we could assume that  $Mg^{2+}$  and  $Fe^{2+}$  impurities in calcium carbonates will prevent the aragonite formation at the upper mantle conditions down to at least 200 km depth (6 GPa).

On the contrary, the present experimental data shows that structures with large cation sites (aragonite, CaCO3-VII, and post-aragonite) are favorable and will likely be observed only in the pure CaCO3, in carbonates with larger cations such as SrCO3 and BaCO3, and in their solid solutions. Our hightemperature experiments show that syngenetic (Sr,Ca) carbonate inclusions in diamonds should have aragonite-type structure.

Because of the generally very small size of inclusions (Klein-BenDavid et al. 2006, 2009), which, furthermore, contain multi-phase assemblages of solid phases (silicates, oxides, carbonates), brines (halides), and fluid bubbles, the analytical techniques for the accurate investigation of the structures and compositions of the minerals are largely confined to TEM (Klein-BenDavid et al. 2006, 2009; Logvinova et al. 2008, 2011, 2019a; Kaminsky et al. 2009), IR, and Raman spectroscopy (Logvinova et al. 2008, 2011, 2019a; Kaminsky et al. 2009). Most of the currently available studies on (Ba,Ca)CO3, (Sr,Ca)CO3, and (Ca,Ba,Sr)CO3 used the TEM method and IR spectroscopy and focused on the composition of the minerals rather than the structures (Klein-BenDavid et al. 2006, 2009; Logvinova et al. 2008, 2011, 2019a). However, in the paper by Kaminsky et al. (2009), Sr- and Ba-bearing calcium carbonate (Ca =  $99.26$ ) mol%), which was found in the diamond nano-inclusions with the low mantle minerals walstromite-structured CaSiO3, CaTiO3, and ferropericlase, was identified as rhombohedral calcite with the aid of Raman spectroscopy. According to the CaCO3 phase diagram (Bayarjargal et al. 2018), at the low mantle conditions, three carbonate phases could be stable: aragonite, CaCO3-VII, and post-aragonite. Aragonite is quenchable to ambient condi-tions, while post-aragonite and CaCO3-VII transform back to calcite on decompression. Given that the addition of  $Sr^{2+}$  and  $Ba^{2+}$  stabilizes both CaCO3-VII and postaragonite type structures at lower pressures, we could assume that the crystallization of the (Ca,Ba,Sr)CO3 in the inclusions should have happened above 20–25 GPa likely in CaCO3-VII or post-aragonite type structures.

Inclusions in natural diamonds provide direct samples of the diamond forming media and could be used to reconstruct the physical-chemical environments in which their host diamonds were formed. (Ba,Ca)CO<sub>3</sub>, (Sr,Ca)CO<sub>3</sub>, and (Ca,Ba,Sr)CO3, despite the scarcity of the findings in inclusions of typical kim-berlite and placer diamonds, are vivid indicators of metasomatic processes (Logvinova et al. 2008, 2011, 2019a). They appear in unique polyphase assemblages together with phosphates, halides, sheet silicates, and abundant fluid segregations and indicate the involvement of a carbonatitic high-density fluid/melt enriched in incompatible elements in the diamond formation (Cl, K, P, Ba, and Sr) (Klein-BenDavid et al. 2006, 2009; Logvinova et al. 2008, 2011, 2019a). The carbon isotopic composition of the host diamonds and predominant eclogitic paragenesis of mineral inclusions imply that this fluid/melt might have been supplied from the subducted rocks of the oceanic and continental litho-sphere (Ragozin et al. 2009).

Thus, the carbonate phase diagrams,  $P$ -T parameters of the phase transitions, and melting temperatures are a useful tool for better understanding of the mantle processes and the recon-struction of their physical-chemical conditions. In particular, CaCO3-SrCO3 and CaCO3-BaCO3 systems are important for the reconstructions of the processes related to the carbonatites metosomatic activity.

The present study shows that the high-pressure phase be-havior of the calcium carbonate-based solid solutions depends strongly on the specific chemical composition of the solid solu-tion. The structural changes derived from the cationic substitution can have an important impact on the phase diagrams of carbon-ates and, thus, affect the physical properties of these materials

such as equations of state, sound velocities, melting temperatures, and other factors. They should be considered in the modeling of the processes that involve carbonates. Overall, it is clear that the carbonate phase diagrams in the complex multi-component systems should be further studied with the use of in situ methods.

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