

Sinehbaghizadeh, S., Saptoro, A., Amjad-Iranagh, S., Naeiji, P., Tiong, A. N. T., Mohammadi, A. H. (2023): A comprehensive review on molecular dynamics simulation studies of phenomena and characteristics associated with clathrate hydrates. - Fuel, 338, 127201.

<https://doi.org/10.1016/j.fuel.2022.127201>

A comprehensive review on molecular dynamics simulation studies of phenomena and characteristics associated with clathrate hydrates

Saeid Sinehbaghizadeh ^a, Agus Saptoro ^{a*}, Sepideh Amjad-Iranagh ^b, Parisa Naeiji ^c, Angnes Ngieng
Tze Tiong ^a, Amir H. Mohammadi ^{d*}

^a Department of Chemical and Energy Engineering, Curtin University Malaysia, CDT 250 Miri,
Sarawak 98009, Malaysia

^b Department of Materials and Metallurgical Engineering, Amirkabir University of Technology, Tehran,
15875-4313, Iran

^c GFZ German Research Centre for Geosciences, Telegrafenberg, 14473 Potsdam, Germany

^d Discipline of Chemical Engineering, School of Engineering, University of KwaZulu-Natal, Howard
College Campus, King George V Avenue, Durban 4041, South Africa

* Corresponding authors, E-mail addresses:

Agus Saptoro, agus.sptoro@curtin.edu.my

Amir H. Mohammadi, amir_h_mohammadi@yahoo.com

Abstract

Clathrate hydrates or gas hydrates have received worldwide attention due to their potential to be utilized in various sustainable technologies. The hydrate-based industrial applications as well as developing green technologies or safely extracting natural gases stored in the nature require profound comprehension of the phenomena associated with gas hydrates. On the flip side, identifying the characteristics of different hydrate formers and the effects of a wide range of introduced additives to these technologies is the critical

objective, so that needs to be deeply investigated at both macroscopic and microscopic scales. The expensive experiments and limited availability of facilities at the nanoscale encourage researchers to apply novel computational methods and simulation approaches. For three decades, molecular dynamics (MD) simulations in the field of gas hydrates have been widely used to mathematically analyse the physical movements of molecules and the evolution of atomic positions in time. In this work, the mechanisms involved in the pure, binary, and mixed gas hydrates, and the impressions of promoters/inhibitors/minerals on gas hydrates were briefly reviewed. Also, the phenomena and properties associated with gas hydrates such as nucleation, growth, stability, dissociation, cage occupancy, storage capacity, morphology analysis, guest role, thermo-physical and mechanical properties, and dynamical and vibrational behaviour of gas hydrates were reviewed. This work aims to provide readers with an extensive overview of MD simulations of gas hydrates to stimulate further research on this riveting field.

Keywords: Clathrate Hydrate; Gas Hydrate; Promoters/ Inhibitors; Molecular Dynamics Simulation; Gas Hydrates Characteristics; Gas Hydrates Microscopic Phenomena

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1 Introduction

Gas hydrates are ice-like compounds that can be inherently formed under certain thermodynamic conditions within a gas/water mixture where hydrogen-bonded water molecules create a crystalline lattice around the gas molecules. Clathrate hydrates are categorized into three types of structures I, II, and H. Gas species are known as small guests with molecular diameters between 0.4-0.55 nm are able to form structure I; components like hydrocarbons with larger sizes mostly between 0.6-0.7 nm can generate structure II of clathrate hydrate; while for the formation of structure H, simultaneous presence of both small guests and large molecular guests (LMGs) normally liquid hydrocarbons with a molecular diameter of 0.75-0.9 nm are needed [1]. Aside from the clathrate hydrates, there are some substances such as tetra-alkylammonium salts/halides e.g. tetra-nbutylammonium bromide (TBAB), with which the guest gases like CH₄, CO₂, N₂, H₂, etc. can generate semiclathrate hydrates [2,3]. Due to the innovative applications of hydrates, they can be at the centre of research within sustainable chemistry. The early era of hydrate explorations was dominated by pipeline blockage and flow assurance, however, the idea of using this phenomenon for sustainable development, supporting the economy, and make to a cleaner atmosphere has recently received widespread attention. So that the processes such as hydrate-based CO₂ capture and sequestration, gas storage and transport, secondary refrigeration, water desalination, gas separation, and energy recovery have been the upsurge of research in this field. Figure 1 shows some suggested industrial applications of hydrate-based methods. The scope of these technologies in the context of a new master plan with concerted collaborative endeavour between various disciplines e.g. engineering, energy, chemistry, and physics have recently been developed.

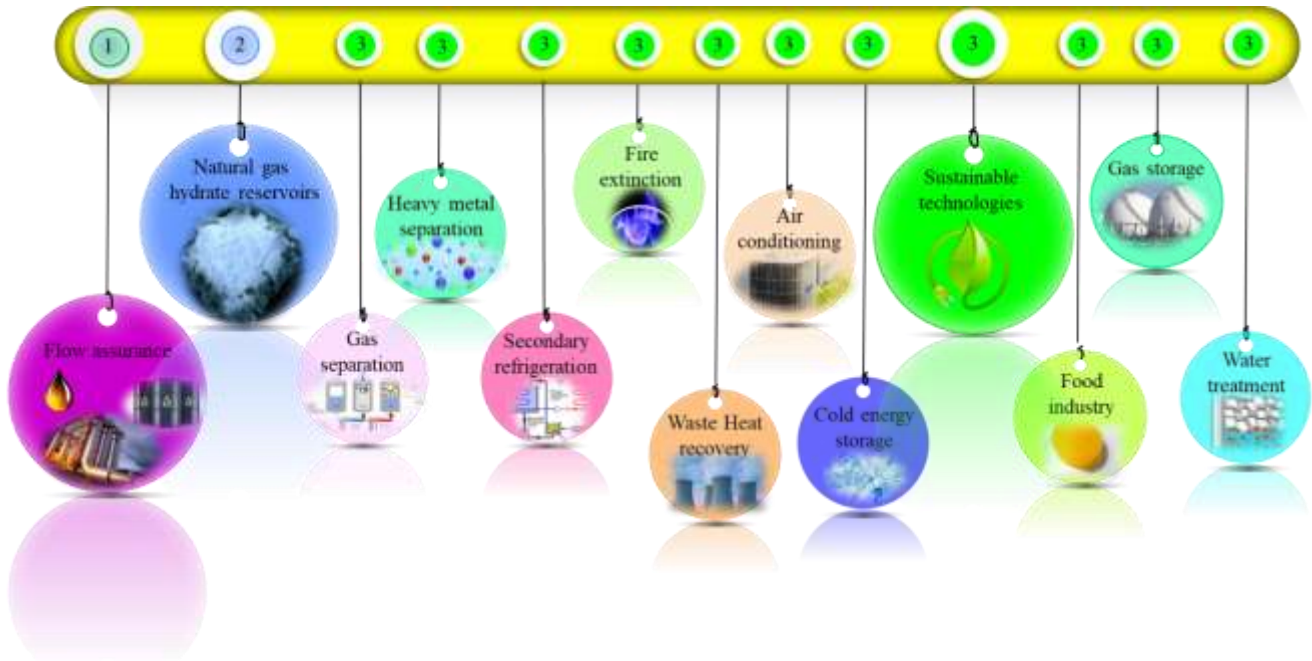


Figure 1: Gas hydrate industrial applications.

Recent experimental results have revealed the variety of properties associated with hydrate-based technologies. Although such investigations can cover the required scientific data and reveal the mechanisms at macroscopic and mesoscopic levels, the microscopic mechanisms and molecular characteristics of these systems cannot mostly be explored in the laboratory. In addition, investigating under the harsh operating pressure-temperature conditions may be the other limitation. In this regard, the utilization of molecular dynamics simulations would be the desirable option by which the vast majority of such characteristics can be probed. These new computational techniques may respond to many questionable issues in the engineering processes involved. The number of reviews on the computational studies of gas hydrates is not substantial. Liang and Kusalik provided an overview of explorations of gas hydrate crystal growth [4]. English and MacElroy conducted a review on thermodynamics, equilibrium properties, thermal conductivity, nucleation, growth, dissociation, electromagnetic fields, energy storage, and CO₂ sequestration [5]. Also, English and Waldron discussed the prospects and challenges of external

electric fields in molecular simulations [6]. Ripmeester and Alavi overviewed the nucleation, decomposition, and memory effect of clathrate hydrates [7]. Somewhat more recently, Alavi and Ripmeester conducted a review focusing on computational studies of H₂ storage in clathrate hydrate phases [8]. Recently, we reviewed the new research findings of CO₂ clathrate hydrates which have been revealed from the molecular dynamics (MD) simulations [9]. Nonetheless, there is still a lack of comprehensive review discussion in the literature on the investigated gas hydrates from the molecular perspective which can be useful for further progress in hydrate-based industrial applications. Therefore, in this work, the phenomena and characteristics of all clathrates (sI, sII, and sH) and semiclathrate hydrates at the molecular level ranging from pure, binary, mixed hydrate systems, and gas hydrates in the presence of promoters, inhibitors, and minerals will be reviewed. Hence the following sections are organized to review the performed MD investigations of various gas hydrate systems (in section 2); explored hydrate phenomena and properties at microscopic scale (in section 3); and the proportions of conducted MD studies of gas hydrates and the future research suggestions (in section 4).

2 Molecular dynamics (MD) simulation theory

Due to the power in calculating the motion and equilibrium of each molecule or atom, molecular dynamics (MD) simulations in diverse fields of science and engineering have received worldwide attention. The knowledge concerning the structural as well as dynamic properties of substances in either gas, liquid, and solid states can be achieved by MD simulations at molecular or atomic scales [10]. The common applications of MD are either to provide explanations by determining the mechanisms involved or to predict the properties of materials. In classical MD, to predict the energies of molecules and equation of motions at thermodynamic conditions, the laws of mechanics are applied. The positions and velocities in molecular systems are dependent on P-T condition and the chemical structure of the simulated system. Empirical interatomic potentials including a long-range Coulombic force and short-range repulsive/ attractive force are usually used to describe the atomic interactions. Generally, two approaches can be

used to conduct the MD simulations; the first is the non-equilibrium mode in which the system away from the equilibrium is stimulated and the system response is followed while in the equilibrium model, the macroscopic property of interest during the simulation is calculated from the time average of that property [11–13]. The theories, approaches, and outputs of MD simulations are displayed in Figure 2.

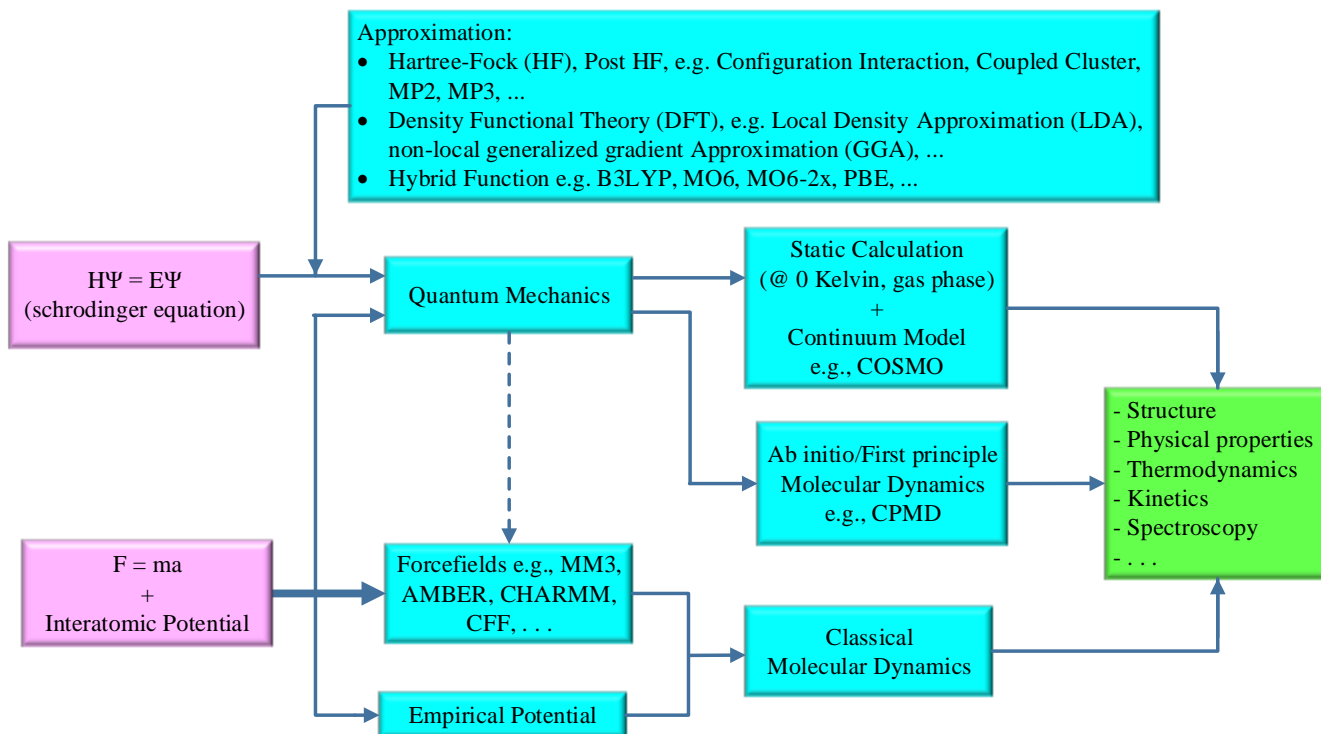


Figure 2: The theories, approaches, and outputs of molecular simulations.

As exhibited, solving the Schrödinger equation is the basis of the calculations of quantum mechanics; more details of this equation have been given elsewhere [14]. Density functional theory (DFT) is an exact theorem that was developed in the physics community. This method carries out energy calculations and electronic structure using the approximation of Kohn–Sham [15] which can be employed for many practical calculations. Besides, DFT can be combined with the HF method which is known as a Hybrid function [16–18]. Worth highlighting that, against semi-empirical methods, *ab initio* calculations can determine atomic interactions accurately in which the calculations are based on principles of quantum

mechanics. However, limitations to treating nanoscale and macroscale decisions are substantial. Many properties that are of critical importance to process performance cannot yet be integrated into process design and computer-aided molecular problems. In this context, the interplay between process design and the design of materials and molecules has been discussed elsewhere [19]. Using MD analysis, many different questions about the molecular mechanisms of phase separation, crystal growth, surface absorption, etc. can be addressed. In this regard, the combination of Monte Carlo and MD can also be employed [20–29]. It should be pointed out that Monte Carlo simulations can be regarded as classical mechanics methods.

2.1 MD Simulations of Gas Hydrate systems

Various aspects of gas hydrates through MD approaches e.g. Free-energy methods, classical MD, potential models, massively parallel MD, DFT, and *ab initio* calculations have been evaluated. It is conceivable that these studies can be rewarding for either hydrate-based gas separation or other utilization approaches of gas hydrates. The phenomena and properties associated with hydrates explored by applying MDs are shown in Figure 3. Generally, the workflow of MD simulations consists of three steps: pre-processing, MD simulations, and post-processing. The first step includes the preparation of simulation box setup and initial configuration, adjustments of the force field (FF) for each molecule, and energy minimization of the system. To provide the initial hydrate structure, the X-ray diffraction analysis for the initial positions of water molecules in the hydrate structure has been reported [30]. However, the coordinates of water molecules and cages for sI, sII, and sH clathrate hydrates specified using TIP4P potential are presented elsewhere [31]. These water molecular orientations of guest-free cages were determined based on a net-zero dipole moment as well as the lowest potential energy. The next step is a production run in which the system is equilibrated under specific thermodynamic ensembles to obtain the desired results. The final step is analysing the parameters along which the properties or phenomena can

be evaluated. To study the intermolecular behaviours, properties, phenomena, and microscopic mechanisms, there are many analysis parameters (target parameters) to employ. Figure 3 exhibits the list of these parameters along with MD steps and employed software in gas hydrate field.

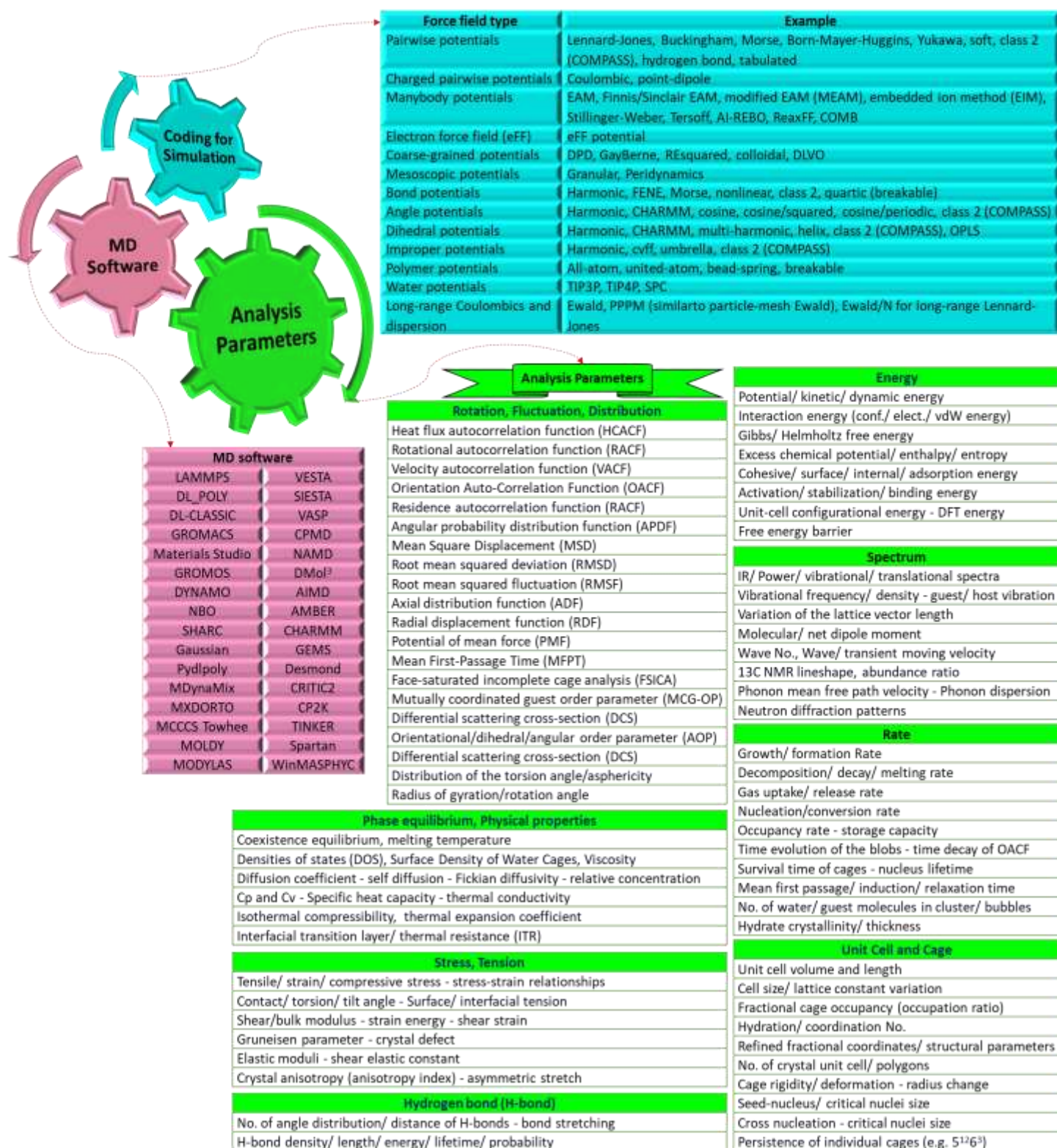


Figure 3: Molecular simulation steps, analysis parameters, and employed software to study gas hydrate systems.

The classical mechanic's methods can be developed based on intermolecular and intramolecular interactions described by empirical potential energy parameters which are the so-called force field (FF). As a fundamental issue underlying all atomistic simulations, Force-field development has drawn considerable attention. A few examples of developed general force fields are MM3, MM4, Dreiding, SHARP, VALBON, UFF, CFF93, AMBER, CHARMM, OPLS, and MMFF. The FFs shown in Figure 3 are also the set of corresponding parameters for a single molecule and the specific force functions for inter/intra-molecular forces. Over decades, several force fields for classical MD simulations have been developed. Therefore, a wide variety of potential models for components has been proposed. As an example, water models to study the characteristics of gas hydrates are: SPC [32–37], SPC/E [38–43], F-SPC [44], SPC/Fw [45,46], TIPS2 [47], TIP3P [48], TIP4P [4,49–59], TIP4P/2005 [60–64], TIP4P-Ew [22,26,65], TIP4P/Ice [27,66–78], TIP4P-FQ [44,79,80], TIP4P/F [81], TIP5P [82,83]; polarizable force fields: AMOEBA [84], COS [85], COS/G2 [86], Stillinger-Weber (SW) model [87–89], mW model [90–96], and KKY potential [97].

These potential approaches can simply be applied to water molecules as employed in many studies due to their computational efficiencies, however, their performances for different systems and operating conditions are dissimilar. For instance, notwithstanding the SPCE/OPLS-UA model that indicated no CH₄ hydrate formation at 230-260 K and 30 MPa, the formation of CH₄ hydrate using the SPC/TSE model was observed [98]. To analyse the accuracy of the coarse-grained mW model, it was compared with TIP4P+OPLS-AA and SPC+OPLS-UA CH₄ models at 250 K and 6 MPa. Based on the results, the mW model was found to be more accurate at recreating structural properties but it has a limitation in identifying small details in the RDFs. This is so that, the utilization of the coarse-grained model depends on what phenomenon is of interest [99]. The coarse-grained method to determine the surface tension at the interface of CH₄ and water, melting temperature, and the hydration number of sI and sII clathrate

hydrate may estimate the CH₄ and water diffusivity, however, it overrates the mobility of guests in solution. Also, it was revealed that the density temperature evolution depends on the ability of a water model in reproducing the solubility of hydrophobic molecules [100].

By comparing CH₄ Hydrate properties using polarizable AMOEBA and COS/G2 force fields with TIP4P and TIP4P-FQ models, it was clarified that AMOEBA and COS/G2 models are better at reproducing the experimental data than TIP4P and TIP4P-FQ models that do not explicitly account for polarization [84]. To simulate the hydrate properties of full occupied CH₄ hydrate at low-temperature ranges, the performance of TIP4P-FQ intramolecular was found to be higher than the other four-site potential models [50,51]. Also, the TIP4P-FQ model outperforms F-SPC, TIP4P, SPC, and SPC/E force fields in predicting the experimental velocity spectrum [44]. Previously, using rigid/non-polarizable SPC/E, TIP4P-Ew, and the rigid/polarizable TIP4P-FQ potentials, thermal conductivity, and diffusivity of CH₄ hydrate were evaluated. Although the predicted thermal expansion using polarizable force fields and isothermal compressibility in comparison with experiments were in agreement, the deviations are substantial at very low temperatures [85]. MD simulations of Metastable sI and sII hydrates showed that in describing the librational region of the spectra, the polarizable force-fields outperform non-polarisable models [101]. Also, to describe differences in the spectroscopy of clathrate hydrates, a combination of anharmonic bond potentials and accurate molecular electrostatics can be applied [102]. The thermal conductivities estimated from MD simulations were found to be independent of pressure. In addition, with decreasing cage occupancy, the simulations predicted a slight increase in thermal conductivity [80]. The three coexistence phases (H-L_w-V) of CO₂ hydrate by examining the performance of TIP4P/Ice, TIP4P/2005, MSM, EPM2, TraPPE, and ZD molecular models showed that the combination of TraPPE and TIP4P/Ice potential models gives a striking accuracy to predict the experiments of CO₂ phase equilibrium up to 200 MPa [103]. Also, TIP4P/Ice and mW models can more

accurately reproduce the three-phase coexistence than SPC, SPC/E, TIP4P/2005, and TIP4P force fields [39]. Previously, heat capacity, thermal expansion coefficient, and compressibility of CH₄+CO₂ hydrate using TIP5P, TIP4P, TIP3P, TIP4P2005, TIP4PIce, TIP4PEw, SPC/E for water, DACNIS-UA, TKMAA, OPLS-UA for CH₄, and rigid three-site TraPPE, EPM2, EPM potentials for CO₂ were simulated. Using the combination of TraPPE CO₂ rigid potential, DACNIS united-atom CH₄ potential, and TIP4P/2005 water models, the most accurate results compared to experimental values were obtained [48].

The assumptions of the vdWP theory using MD analysis determined that changes in the configurational energy by the nonplanar TIP5P model are higher than the planar SPC [83]. The selected potential model can also affect the estimation of hydrate fractional occupancies. For instance, by employing Gubbins (MG) and OPLS potential models, the accommodation of 2 and 5 CH₄ molecules in large cages of sH hydrate at 300 K and 2 GPa were estimated respectively [104]. Recently, it was shown that the stability of CO₂ hydrate is quite relevant to the water molecular interactions in which the order of different FF for this hydrate stability would be TIP4P/Ice > TIP4P/2005 > SPC/E > SPC/Fw [105]. Simulations displayed that the type of water force field quantitatively affects the estimated hydrogen bonding but does not affect the qualitative trends. Also, a stronger electrostatic interaction with the guests and neighbour water molecules was observed for TIP4P/ice potential compared to the SPC/E model [106]. Simulation of the CH₄ hydrate in conditions of oceanic sediments showed that the SPC/E water model requires a shorter simulation time than the TIP4P model [41]. However, the impacts of potential models on CH₄ hydrate dissociation conditions in the other research were found to be negligible. [107]. Although the deviation of MD predictions by TIP4P/Ice and OPLS-UA force fields was approximately 3 K below the experimental values, calculated CH₄ solubility was found in good agreement with continuum scale models [108]. According to an evaluation of the accuracy of TIP4P/Ice, GAFF, and OPLS-AA force fields in predicting the phase equilibrium and enthalpy of TBAB semiclathrate at various temperatures,

it was concluded that compared to the experimental data, the OPLS-AA outperforms the GAFF. Also, TIP4P/Ice model can more accurately describe the interactions of water molecules in TBAB semiclathrate (type B) [109]. MD analysis of hydrogen bonding and guest conformation of 1-propanol hydrate showed that a larger value of the dihedral angle obtained from the single crystal XRD measurement than the value predicted by the MD simulation might be related to the approximate nature of the intramolecular dihedral potential in the force fields [110]. It should be noted that for intermolecular separations higher than 50% of the unit cell length, Lennard-Jones interactions can assume to be zero [35]. [Table S1](#) presents the list of force fields for hydrate formers and promoter/ inhibitor additives employed in different gas hydrate simulations.

Generally, the gas hydrates based on hydrate formers and additives which contribute into the solid phase and solution phase can be classified into pure, binary, and mixed gas hydrates, and the impressions of promoters/ inhibitors/ minerals on gas hydrates. The MD simulations of these hydrate systems are briefly reviewed in the following sub-sections.

2.2 Pure gas hydrates

Many MD studies have focused on determining the crystallization/ dissociation mechanisms and contributing factors that govern the nucleation, dissociation rate, guest-host interactions, molecular mobility, cage occupancy, and cage preference. In this context, it was reported that at the same thermodynamic conditions, dissimilar cage preferences and occupancy of CH₄ and C₂H₆ may result from their differences in size and shape [111]. MD simulations confirmed that the two-steps dissociation model which was introduced by Sloan [1] is reliable for the mechanism of hydrate dissociation [82] in which the dissociation rate directly depends on the fractional occupancy [112]. Also, expanding the water-hydrate interfacial contact area and higher initial temperature can effectively promote hydrate decomposition [113]. It was proved that the thermal-driven breakup of CH₄ hydrate is controlled by the

diffusion of CH₄ molecules from the hydrate surface to the liquid phase; additionally, break-up for empty hydrate was determined about an order of magnitude faster than filled clathrate hydrate [79]. Also, CH₄ in the gas-like fluid transforms from high-density to low-density while reversely occurs for dissolved water in an ice-like phase. This phenomenon is more dominant at an upper melting temperature [40]. Monitoring hydrate dissociations confirmed the formation of nano-bubbles when the mole fraction of dissolved CH₄ in the water phase is at least 0.044 [72]. MD simulations also revealed an approximately similar growth mechanism for CO₂ and CH₄ hydrates [114]; water molecules that are adsorbed and tend to complement the open large cavities induce the rearrangements at the surface of CH₄ clathrate hydrate [97]. Also, the nucleus size and relative positions of the guest molecules have bearing on the control of hydrate nucleation [77]. During the induction times, the formation of structural defects within hydrate lattices can also be observed [64]. The MSD analysis elucidated the anomalous diffusion and anisotropic characteristics of the H₂ guest molecule [115]. Based on MD findings of thermal expansion, the unit cell volume is significantly dependent on temperature. Also, to analyse the self-preservation mechanism of gas hydrates which takes place at the interface, MD simulations have exhibited that the formation of a layer of solid-like water increases the resistance of mass transfer against guest diffusion from hydrate, consequently, inhibiting further dissociation [55]. MD is also highly useful for investigations at equilibrium conditions that cannot be determined by experiments. For instance, analysing H₂ hydrate properties at vigorously high pressure or low temperature (e.g. above several GPa or near 0 Kelvin) cannot be performed by the use of X-ray diffraction patterns but MD verified the existence of different possible structures like (sT' and C0-II) at such harsh conditions [116]. In a way of more accurate modelling, select force fields respective to guests and simulation techniques are highly crucial. Also, between different ensembles, the sequence of accurate clathrate nucleation was found to be NPT > NVT > NVE; however, the crystallinity sequence is exactly reversed [117]. According to the decomposition

of CO₂ and CH₄ hydrates at 180-280 K, it was concluded that hydrate stability using isochoric conditions is lower than that in isobaric conditions [118]. Although remarkable advances in macroscopic measurements have been accomplished, MD simulations as a powerful technique can provide insights into the fundamental mechanisms of gas hydrates at molecular and atomistic levels. In this regard, computational studies of pure gas hydrates would be a reasonable alternative.

2.3 Binary and mixed gas hydrates

To reveal the interfacial behaviors and characteristics of the mixed gas hydrates which play a central role in developing hydrate applications, investigating molecular interactions can aid in accurately determining the characteristics of their processes. Bearing in mind that a fundamental understanding of different aspects of mixed gas hydrates will be critical for many applications e.g. flow assurance, gas storage, and transportation and MD probes can support the macroscopic experimental studies. MD simulations showed the guest-guest interactions which were ignored as the assumption of vdWP theory can contribute to the Helmholtz free energy and subsequently some deviation in predicting hydrate equilibrium conditions [83]. It was also found that the mass transfer, memory effect, and guest molecule chemical potentials are the main controllers of mixed CO₂+CH₄ hydrate in the replacement phenomenon [119]. This phenomenon without structural change occurs first after partial melting of CH₄ hydrate surface and followed by partial collapses of large cages at the interface and entering CO₂ molecules into them [120]. Worth highlighting that at the macro level, the free water was identified as a significant feature during the replacement [121], however, hydrate growth may become a quasi-static equilibrium at the hydrate-liquid interface rather than in the free gas [122]. Interfacial characteristics of brine water and CO₂+CH₄ mixture indicated that the degree of IFT reduction consistent with experimental evidence is directly proportional to the CO₂ concentration [123]. MD simulations can also contribute to molecular level explorations at the early stages of hydrate nucleation and the roles of mixed guest molecules in such

phenomena. Cage composition details for $\text{CH}_4+\text{C}_2\text{H}_6$ and $\text{CH}_4+\text{C}_3\text{H}_8$ hydrate systems indicated that increasing CH_4 composition greatly reduces the appearance of other complete cages whereas grows the formation of standard cages of sI and sII hydrates (i.e., 5^{12} , $5^{12}6^2$, and $5^{12}6^4$ cages). This suggests that the more stable hydrate nuclei should be obtained with an increasing CH_4 percentage. Dissimilar to those mixed systems, $\text{C}_2\text{H}_6+\text{C}_3\text{H}_8$ hydrate exhibited that growing C_3H_8 percentage decreases the formation of sI large cages ($5^{12}6^2$), however, promotes the appearance of sII large cages ($5^{12}6^4$). The results of the average composition of cages over the final 30 ns of nucleating trajectories for the aforesaid mixed hydrates are shown in Figure 4 [124].

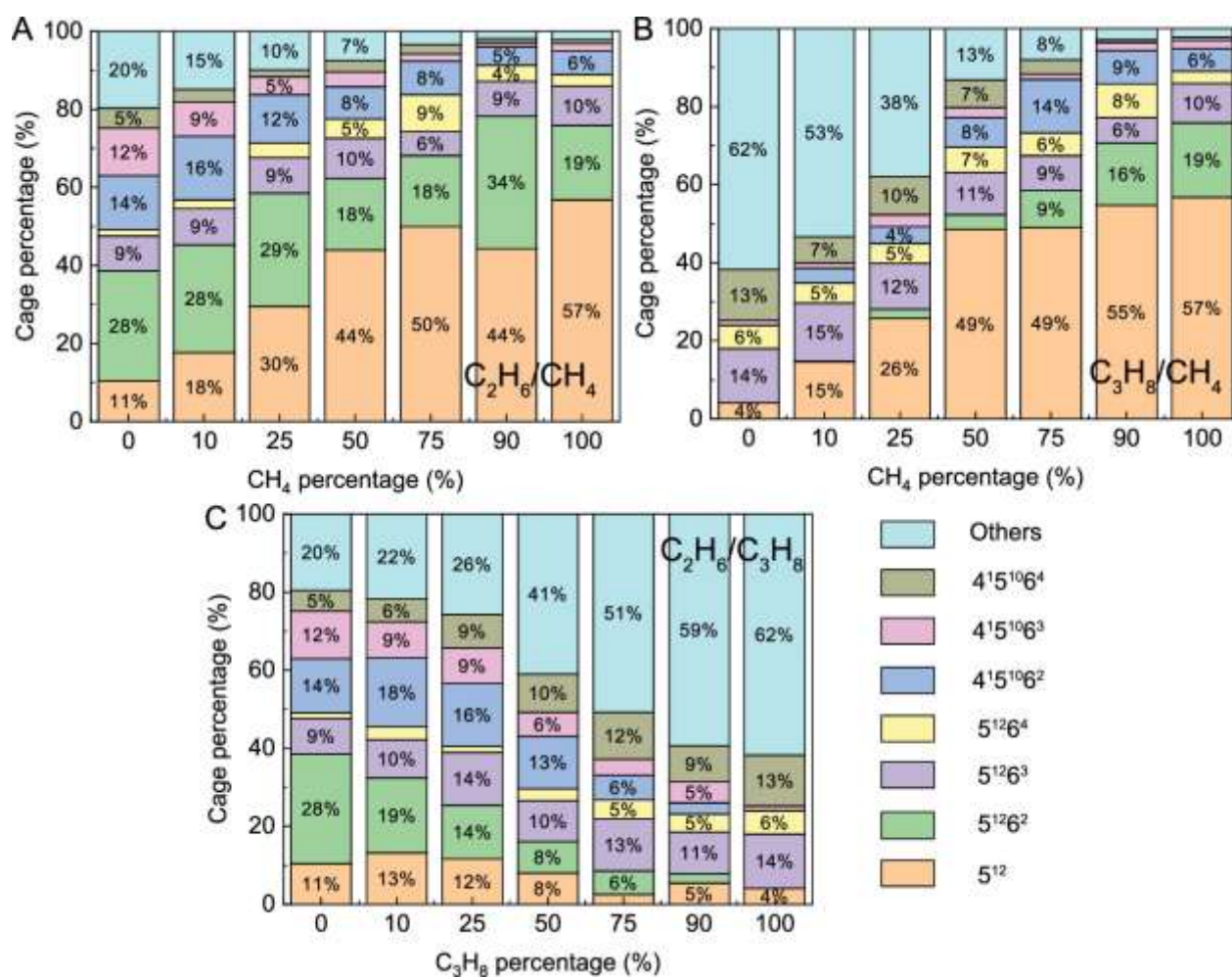


Figure 4: Average composition of cages in (A) $C_2H_6+CH_4$, (B) $C_3H_8+CH_4$, and (C) $C_2H_6+C_3H_8$ mixed hydrates with varying compositions [124].

The number of complete cages during hydrate nucleation in the CH_4+H_2S nano-bubble system revealed that forming a nucleus is normally associated with first the emergence of 5^{12} and then $5^{12}6^2$ cages, followed by both $4^15^{10}6^2$ and $5^{12}6^3$ cages almost simultaneously. Entropy can also play a critical role in the appearance of specific cage types. For instance, since 5^{12} cages are highly symmetric than $4^15^{10}6^m$ and $5^{12}6^n$ cages, they have higher degeneracy and are structurally more consistent with the aqueous phase, resulting in smaller entropic penalty linked with their generation. However, space-filling structures require the appearance of additional cages. Besides, the earlier formation of $5^{12}6^2$ cages because of having higher symmetry than $4^15^{10}6^2$ cages and fewer water molecules than $5^{12}6^3$ cavities is expected to be entropically favorable [125]. For the sake of a weak H_2S-H_2O hydrogen bond, the presence of H_2S in mixed hydrates stabilizes crystal defects in the crystal lattice. These H-bonds are transient with H_2S which typically behaves as the hydrogen bond donor [126].

2.4 Gas hydrates in the presence of promoters

For energy storage and transportation purposes, perhaps CH_4 and H_2 are the most potent candidate. To date, a number of materials/ techniques have been suggested, and a hydrate-based methodology would be the proper option. Since pure H_2 hydrate requires uneconomic conditions, the presence of additives as promoters such as TBAB [127], THF [128], or SF_6 [129] have been suggested. Given that identifying the key controllers of the kinetics and thermodynamics of gas hydrates in contact with promoters for the suggested hydrate-based applications is particularly critical. As a direct and valuable method, MD investigations have manifested various aspects of microscopic contributors. For example, recent MD simulation results of CO_2 double and mixed hydrates showed that the type of large molecular guests in

the large cages plays a major role in the stabilization of the clathrate network. In addition, among different types of sII thermodynamic promoters, cyclopentane, and cyclohexane in comparison with greenhouse F-gases (F-promoters) seem to be more susceptible to maintaining the stability of CO₂ clathrate hydrate [130]. Gas hydrates can also be useful for gas storage aims, however, the common challenge associated with this is to balance the trade-off between thermodynamic efficiency and H₂ storage capacity. Raman spectroscopic measurements revealed that if sII SF₆ hydrate is exposed to H₂ molecules in the vapor phase, as time proceeds, the H₂ concentration in hydrate will hit a peak and then reduce before reaching a stable value [129]. Although this phenomenon cannot be fully explained through spectroscopy experiments, its dynamic process was recently explored through MD simulations in which every large cavity was initially co-occupied by both SF₆ and H₂ molecules. It was observed that by the end of the simulation time as is shown in Figure 5 (left), the neighbouring small cavity became partially broken, and then the H₂ molecule penetrated them. The red structure shows the initially occupied large cage, and the cyan structure is a partially broken small cage containing an H₂ molecule that diffused into it. Also, as Figure 5 (right) shows, over enough time, the periodic clathrate turned into an amorphous structure and H₂ molecules escaped from their original cavities, gathered, and formed clusters which imply that the large cage co-occupancy of SF₆ and H₂ molecules could not be the stable configuration. Besides, infiltrating into small and large cages are two main patterns of H₂ diffusion as displayed in Figure 5 (right). The partially broken hydrogen bond ring coloured in green shows the opening through which the H₂ molecule can move. Although occupancy of large cages by H₂ molecules (once they are filled with SF₆) approaches zero, H₂ molecules have sufficient space to pass through the hexagonal ring without any cage breaking.

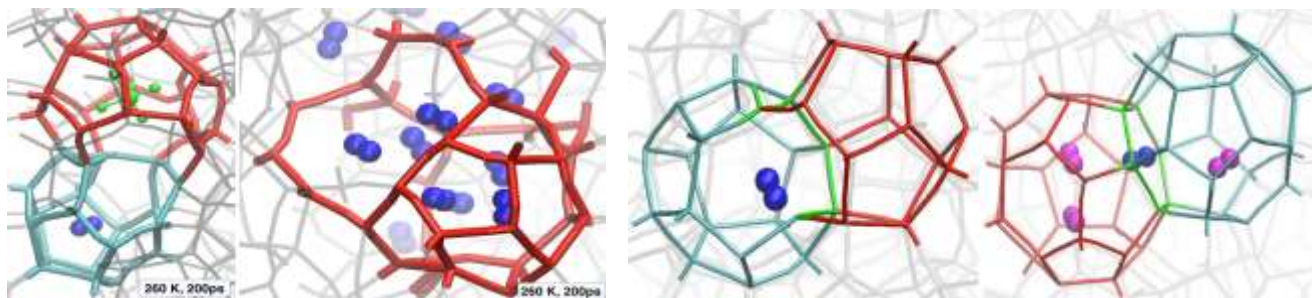


Figure 5: (Left) H_2 penetration into neighboring small cavities. (Right) A cluster of H_2 molecules in the bulk system; diffusion of H_2 molecules from one small cage to another (left) and from one large cage to another (right) [131]. (atoms of H_2 and SF_6 are in blue and green color respectively)

Also, the penetration processes in two special structures of H_2 hydrate (C_0 and sT') determined that dissimilar to the sT' structure, the number of jumps of H_2 molecules among neighbouring channels in the C_0 structure is noticeable. However, due to the growth of H_2 jumps, with changing temperature from 140 to 260 K, a little acceleration of diffusion can be observed [115]. Generally, pressure and temperature have a critical impact on the growth mechanism, cage occupancy, and storage capacity of H_2 +THF hydrate. At 50 MPa and 250 K, the H_2 +THF hydrate growth rate can be 3.6 times higher than the pure THF hydrate which implies the kinetics promotion role of H_2 molecules. The occupancy of small cages can be enhanced by the presence of H_2 molecules while a faster rate of growth can facilitate the formation of H_2 -filled large cages. However, there is a striking resemblance between the growth behaviour of H_2 -filled large cages. However, there is a striking resemblance between the growth behaviour of H_2 +THF and pure THF hydrates which indicates that THF is the major controller of the H_2 +THF hydrate growth [132]. Recent MD investigations showed that the addition of THF+DMSO can help CO_2 molecules diffuse into the CH_4 hydrate more readily than either the case with no additive or utilization of a single THF [133]. It was also manifested that in line with experimental results [134], CO_2 behaviors in the presence of SDS due to very strong distortion of SDS are quite dissimilar to CH_4 interactions with SDS. Both SDS apolar and polar ends lose its shape once it faces CO_2 molecules which becomes SDS incapable to exert any substantial driving activity to promote CO_2 hydrate formation [135]. The exclusion

of SDS and methylene blue (MB) organic molecules from hydrate structure during CO₂ hydrate growth indicated that these molecules do not contribute to the crystalline structure [136]. However, the promotion effects of these molecules are different. For example, the urea surfaces make this molecule operate as a catalyst for the formation of CO₂ hydrate layers [137]. The role of urea in the growth process was found to be an increase in mass transport and catalysation of forming cages at the solid-liquid interface [138]. Recent MD findings for the solution phase including urea and Cu, Fe, Ag metal particles revealed that the mixture of Cu, Fe, and urea (without the inclusion of Ag) possesses desirable promotion effect on the CO₂ clathrate hydrate growth rate [139]. The impact of surfactant at near ambient conditions highlighted that the guest molecules can get absorbed into the aggregates during the aggregation process which provides structural flexibility and enhances the aggregation kinetics [140]. In addition, creating hydrate memory effect by organic molecules (e.g. lecithin) may help hydrate be dissociated more slowly when they are adsorbed on the hydrate surface through their hydrocarbon chains crossing and narrowing the available space for water movement and hydrate [141]. Simulating the marine hydrates in contact with protein also confirmed the role of microorganisms in accelerating marine hydrate formation via an approximation mechanism of enzymatic catalysis [142]. CO₂ hydrate growth in the presence of metal particles showed that the increase of the concentration of the metal particles can accelerate CO₂ hydrate formation but inhibit when their mass fractions were too high which is caused by strong Brownian motion in the solution [143]. The performance of combined promoters may also be more efficient. For instance, the extent of CH₄ in the hydrate phase can be increased due to the addition of both DMSO and THF. The ratio of CH₄ to CO₂ in the existence of THF and THF+DMSO were 1.06 and 1.34 respectively [133]. A combination of absolute thermodynamic and MD modelling for CO₂ interactions used to analyse its consequences for hydrate formation revealed that the generate hydrate nearest to the hematite surface

possesses lower chemical potential, therefore, it would be thermodynamically favoured to adsorb on hematite [144].

2.5 Structure-H of gas hydrates

MD research can be advantageous for determining the micro-scale properties of sH hydrates toward upgrading the hydrate stability and its performances for storage aims. The correlation between the molecular reactivity and stability of H₂+LMGs sH hydrates through analyzing chemical potential and electrophilicity index, electronegativity, and cohesive energy showed that higher hydrate stability corresponds to the larger value of these parameters [145–149]. Studying the stability and chemical activity of 16 promoters on H₂ sH hydrates at 230-270 K highlights that by utilizing alkane components with several heavy atoms less than 7, better stability for H₂ sH hydrates can be obtained. Moreover, small cavities have a higher hydrogen-trapping ability than medium ones [150]. Hence, the role of LMGs in sH hydrate nucleation/ formation mechanisms, as well as molecular-level factors were found to be unequivocally unique. More recently, it was revealed that compared to a pure CH₄ hydrate, 2,2-dimethylbutane (DMB) promotes the formation of CH₄ hydrate by at least 5 times faster than its absence. More specifically, the hydration shell of DMB behaves as a heterogeneous nucleation site for hydrate formation. This occurs due to the longer residence time of water molecules in the first hydration shell. Indeed, the greater the structural order, the slower the dynamics of DMB relative to that of CH₄ [151]. Figure 6 shows the number of complete cages, the time evolution of the potential energy, and the representative trajectory of CH₄+DMB sH hydrate nucleation. As exhibited, from 200-230 ns a sharp increase in the number of CH₄ associated with DMB confirmed the first persistent complete cage of the largest cluster. After that more cages were observed to form near the DMB but at 400 ns, the DMB started to create the hydration shell of CH₄+DMB sH hydrate. In this Figure, carbons and hydrogens of DMB are shown in cyan and purple, the water and CH₄ molecules are represented as blue and orange spheres,

red-dashed lines and red tubes denote the first hydration shells of the associated CH₄ molecules and the cages in the largest cluster. Gray lines, green and cyan spheres are also the liquid water, DMB, and CH₄ molecules respectively.

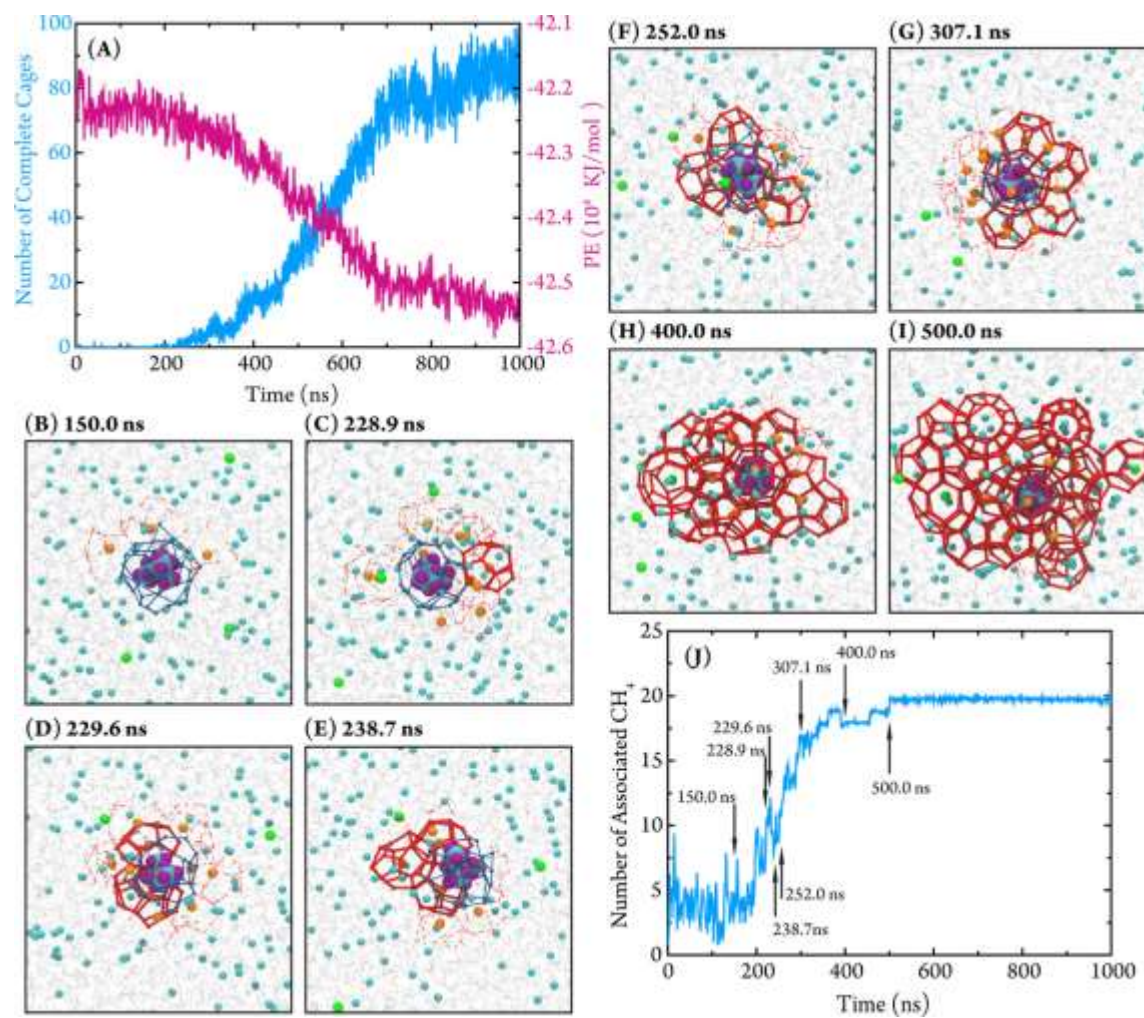


Figure 6: Time evolution of the potential energy, the number of complete cages, and cluster evolution from the trajectory of CH₄+DMB sH hydrate nucleation [151].

Recently, several studies have focused on understanding the effect of guest size and the structural anisotropy which defines the directional dependency of sH hydrate properties [152–154]. In this respect, sH hydrate elastic constants indicated that the type of the LMGs is the main contributor to sH hydrate elastic properties in which LMGs in the large cage of sH hydrate were found to increase its rigidity as

well as resistance to shear deformations. Based on DFT calculations, hexagonal sH hydrate compared to hexagonal ice (Ih) possesses higher resistance and stiffness to principal strains [155,156]. To describe the stiffness of materials, calculated values of Young's modulus as one of the key mechanical properties also demonstrated more resistance of filled sH hydrates to elastic deformation than empty structure. However, compared to sI CH₄ hydrate, sH has lower resistance to shear deformations [157,158]. These differences can be linked to the relatively isotropic sI and sII gas hydrates as compared to the anisotropic characteristics of sH and the role of LMGs. Also, the stability of the overall structure strongly depends on both types of guest molecules (CH₄ and LMGs) [159]. MD investigation focusing on the impact of guest size and conformation on the structure and stability of sH hydrates showed that tilt angle in the cages, structural flexibility, and guest molecule size are identified as key parameters of stability of sH hydrates. Generally, in large cages, molecules with shorter lengths have larger tilt angles. Additionally, the width dimension of the LMG may affect the tilt angle. [153]. Also, the sensitivity of cell parameters and thermodynamic properties of sH hydrates to the temperature is higher than pressure [160]. Structural characterization of sH hydrates (Ar/ CH₄/ N₂/ Xe + neohexene) at the atomistic level and 0 K revealed that in distinction to nearly isotropic clathrate hydrates, anisotropy is a distinguishing feature of sH hydrate. This structure was found to be more resistive to uniform compression and denser as well as having an isothermal bulk modulus in which the type of help gas plays a crucial role [161]. By analysing the free-energy difference to evaluate the thermodynamic stability conditions of sH hydrates, a strong relationship between the σ parameter of Lennard-Jones and free energy was certified. However, using this method, the highest stable structure for sH hydrate cannot be precisely specified [162].

2.6 Gas hydrates in the presence of inhibitors

Safety and economic concerns caused by pipeline plugging have driven the exploration for finding efficient inhibitors that help to prevent the crystallization of clathrate hydrates. Molecular insights into

the decomposition process of gas hydrates under the influence of inhibitors can aid flow assurance experts to find cost-effective and eco-friendly inhibition compounds. In this regard, different MD studies of gas hydrates in the presence of thermodynamic or kinetic organic/ inorganic or mixed inhibitors have been conducted. Also, MD simulations can aid to comprehend the factors that play a pivotal role in controlling thermodynamics and adsorption of kinetic inhibitors of gas hydrates. The interactions between hydrate, water, and inhibitor molecules, cohesion and adhesion of hydrate particles, mobility, binding free energy at interfaces, etc., can be accurately analysed via MD [163]. Evidence suggests hydrate in brine solutions requires a lower dosage of anti-agglomerants (AAs) which keep hydrate particles dispersed in fluids and prevent aggregation [164,165]. In this context, MD simulations showed that salt does not change notably the binding configurations. The free energy of binding with the increase of NaCl concentration from 0 to 10% becomes more negative which can be attributed to the two individual phenomena. The increase of salt content in the aqueous solution decreases the AA solubility, as a result, AA tends to move into the interfacial region. Additionally, an excess of chloride anions near the hydrate surface establishes a negatively charged interfacial layer which disrupts the network of hydrogen bonding close to the hydrate interface [166]. The influence of surfactants on the clathrate wettability indicated that they can bind to both water-clathrate and oil-clathrate interfaces which promote the nucleation for the former and prevent their coalescence with water droplets as well as an agglomeration of clathrate particles through altering the clathrate wettability from water-wet to oil-wet for the latter [167]. It was determined that the contact angle of a water droplet at a dodecane-clathrate interface (Figure 7a) can be $\theta=34 \pm 2^\circ$ and clathrate covered with a monolayer made of intercalated dodecane and dodecanol or with a pure dodecanol monolayer (Figure 7b) at the dodecane-clathrate interface can be the same which is consistent with the water wetting angles more than 150° reported in experiments. Also, Figure 7c shows that being exposed to vapor instead of alkane, the contact angle of dodecanol-covered clathrate can be decreased from 180°

to $88 \pm 2^\circ$. This indicates that the van der Waals interactions among alkane and water stabilize the water droplet higher than vapor acts. For a water droplet, even-if surfactants partially cover the clathrate surface, high contact angles will be obtained. Concur with the experiments, the existence of a surfactant monolayer at the oil-clathrate interface by changing the hydrophobicity of the hydrate surface alters the wettability of the clathrate surface to oil-wet as is shown in Figure 7d. Hence as MD simulations show, the coalescence of a bare water droplet cannot be made by the surfactants but the kinetic barrier by their presence will be increased [167]. MD also determined that although the inhibition effect of light oil (toluene, and iron) on the CH_4 hydrate growth can be observed, the presence of asphaltenes enhances the hydrate growth rate [168]. Investigation of the surfactant layer impacts at the water-oil-clathrate interface on the melting temperature of sII and sI clarified that the factors which contribute to preventing the clathrate agglomeration are the density of the interfacial layer, the length of molecules of surfactant, and the binding affinity and intensity of surfactant molecules to the surface of the clathrate hydrates [167].

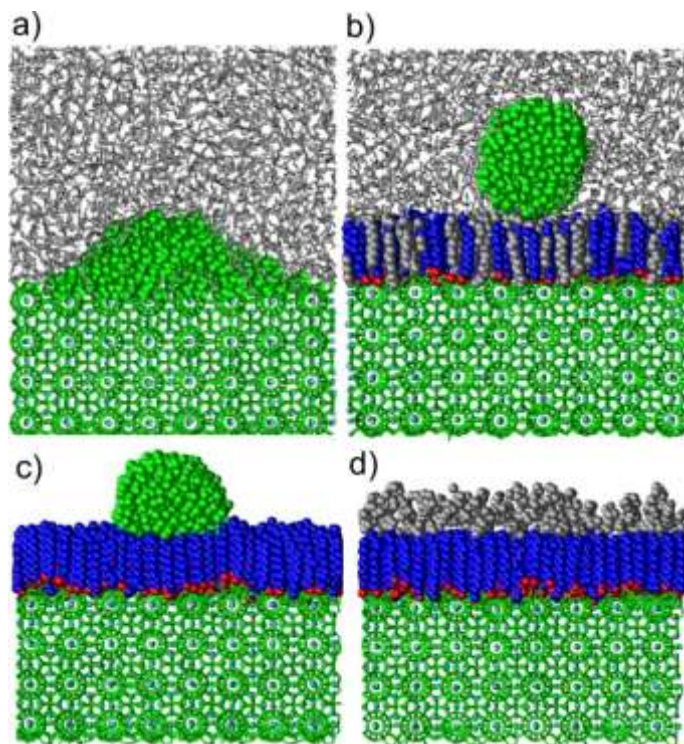


Figure 7: The change of clathrate surface from water-wet to oil-wet by surfactants. (a) Water wets the bare hydrate surface, forming a nonwetting droplet on the surfactant-covered clathrate surface. (b) The clathrate surface is covered with the dodecane-dodecanol interfacial film which repels the water droplet in the alkane oil and behaves as a super-hydrophobic surface. (c) water at the vapor-clathrate interface in the existence of a dodecanol monolayer possesses a less contact angle compared to oil. (d) Dodecane fully wets the clathrate surface [167].

MD also reveals that the inhibition mechanism of organic and inorganic inhibitors is not identical. As an example, in the initial stage, the CH₄ hydrate dissociation rate in the NaCl and methanol solutions are slower and faster than that in pure water respectively. However, in the final stage, the NaCl solution finishes the dissociation process earlier than that in pure water. Indeed, methanol in contact with CH₄ hydrate lowers the free energy of supersaturation and facilitates the CH₄ bubble formation but NaCl due to the strengthening of the hydrophobic interactions can boost the bubble formation more conveniently than pure water [169]. Examining the effect of salt NaCl on the aggregation of CH₄+C₃H₈ hydrates showed that with increasing NaCl concentration in solution, the solubility of the Anti-agglomerate in the aqueous solution decreases, and the binding free energy becomes more negative [166]. Previously, the structural effect of TBAPS kinetic inhibitor on the interface of CH₄+C₃H₈ hydrate was also confirmed, however, the sulfonate group and the hydrophobic end of the molecule make it act as a bi-functional additive [170]. CH₄ hydrate formation with kinetic inhibitor indicated that the surface of the crystals can be blocked by hydrophobic inhibitors at the initial stage while hydrophilic ones impose the inhibition effects by disrupting the water structural network [171]. By making a comparison between CH₄ hydrate in the absence and presence of 5% of methanol, it can be implied that although the existence of methanol reduces the equilibrium stability conditions of CH₄ hydrate, it can simultaneously boost the diffusion of CH₄ molecules at the interface by at least 40% [172]. CH₄ hydrate in the presence of antifreeze proteins (AFPs) showed that the structural aspects of the protein and hydrate surface can contribute to the binding. Also, nonbinding sites are not required for the operation of the protein as a hydrate inhibitor [173]. In

addition, simulations of hydrate growth of CH₄ in the existence of type I and III antifreeze proteins (AFPs) determined that a type I AFP binding on the hydrate surface can be stabilized by the formation of hydrogen bonds between water molecules and anchoring pendant groups [174]. Moreover, the exchange of hydroxyl groups and amide by mutagenesis may change the hydrogen bonding capability and the side chain shape which implies that the hydrogen bonds are not directly responsible for the activity of AFP III antifreeze [175]. MD simulations revealed that inhibitor molecules e.g. PEO/ PVP/ PVCap/ VIMA that have been found to exhibit more inhibition strength experimentally also possess higher free energies of binding [176]. Investigation of the mechanism of CH₄ hydrate growth with combined ethylene oxide (EO) and PVCap showed that the amide H-bonds do not contribute to the adsorption, however, dissimilar to antifreeze proteins, PVCap adsorbs on various crystallographic planes of clathrate hydrates [177]. Analysing the effects of amino acids (glycine/ alanine/ serine/ proline) on CH₄ hydrate formation at 250 K and 15 MPa revealed that among the simulated amino acids, serine has the highest inhibition influence on the hydrate growth whereas alanine and proline have the lowest. Besides, the more the concentration, the less the hydrate growth rate. H-bonds between water molecules and amino acids can mostly damage the cavities [178]. Also, CH₄ hydrate in the presence of 8 different oligomer inhibitors (PVPs/ amino acids) highlighted that except for asparagine, other inhibitors are absorbed into the interface of CH₄-water and suppress the growth of nanoclusters [179]. According to the results of CH₄ hydrate formation with PVP/ PVCap/ PDMAEMA, it was indicated that PVP increases the induction time for the formation of the first hydrate nuclei. Also, the synergistic impact of PVP and PVCap were found to be more efficient than PVCap alone [180] but the existence of PMAEMA reduced the temperature at which sustained hydrate growth can be observed. Hydrate content and the number of H-bonds also can be decreased by increasing the temperature [181]. A molecular study of CH₄/ C₃H₈ hydrate + PVP/ PVCap determined that PVCap kinetically outperforms PVP and such kinetic inhibitors without coming into direct contact

with hydrate surfaces can trigger the dissolution of the hydrate [182]. The properties of kinetic hydrate inhibitors (KHIs) on CH₄ hydrate growth at 260 K and 80 MPa showed that during the adsorption process, the heterocyclic ring of PVCaps intensively tends to approach the interface of hydrate-liquid [183]. By conducting molecular insights into the kinetic inhibitor adsorption on the surface of CH₄ hydrate, although the contribution of the adsorption affinity of hydrogen bonding between water molecules and the amide group of the inhibitor was not observed, a monomer of the KHI (PVCap) was found to be vigorously adsorbed on the surface of CH₄ hydrate [184]. GNs in the aqueous phase can also form H-bonds with water at the hydrate surface which reduces the water activity to obstruct hydrate growth [185].

2.7 Gas hydrates in the presence of minerals

For gas hydrates in marine sediments, the zone hydrate formation is an area of sedimentary rock or unconsolidated clay saturated with gas and water in which clay minerals include common constituents of various grain particles in gas hydrate-bearing sediments such as chlorite, kaolinite, illite, montmorillonite, etc. [186,187]. The intercalation of CO₂ in minerals within brine aquifers mainly causes significant alteration in the spacing between the mineral layers [188]. The interactions among surrounding sediment grains and hydrates determine the physical properties of the hydrate-sediment matrix. It has been experimentally uncovered that the appearance of the solid surfaces could notably influence hydrate formation [189]. In some studies due to the reduced water activity in confined space, the stronger inhibition compared to the promotion effects has been reported [190,191]. In contrast, experiments revealed CH₄ hydrate formation for conditions milder than that required in the bulk can be promoted by the confinement effects of porous materials [192–195]. Due to the experimental challenges associated with quantifying hydrate nucleation and growth, MD simulations were recently applied for examining the CH₄ hydrate in hydroxylated silica pores. Based on the results, at relatively milder conditions (at simulation pressures below 10 MPa), confinement impacts can surprisingly induce CH₄ hydrate growth.

In fact, at constant temperature and in the bulk system, the external pressure governs CH₄ hydrate growth whereas in the confinement zone and at pressures less than the minimum expected experimental CH₄ hydrate phase equilibria, it grows regardless of the applied external pressure. Figure 8 displays the initial configuration and lateral growth of the hydrate slab at the water-hydrate interface in which complete cages adjacent to the silica surfaces were replaced by a thin water layer [196]. The formation of a convex-shaped CH₄ slab in the first 1 μs is followed by spontaneous lateral hydrate growth for the sake of CH₄ nanobubble formation over 6 μs.

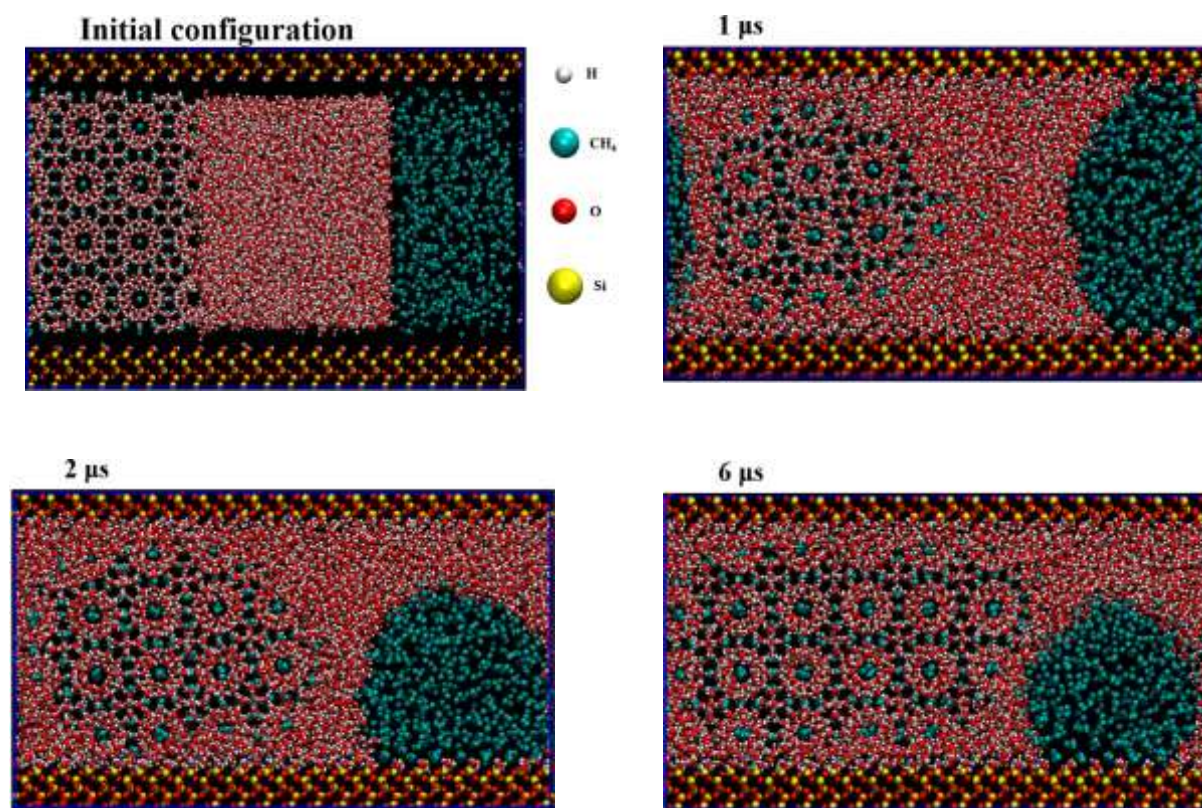


Figure 8: The initial configuration; (1μs) The transition of CH₄ slab into a convex-shaped CH₄ bubble, (2μs) a hemisphere CH₄ nanobubble formation, (6μs) The final snapshot of the simulation. [196].

Compared to pure water, bentonite can enhance the hydrate formation of a gas mixture rather than pure CH₄. However, by applying MD simulations and neutron scattering experiments, it was found that hydrate

formation is insensitive to the addition of impurity particles. Although mineral particles cannot be the most source of heterogeneous nucleation sites, they may do so indirectly if they have a notable impact on gas uptake into solution [197]. Somewhat more recently, MD manifested that kaolinite and illite minerals do not greatly influence the CO₂ hydrate phase equilibrium whereas montmorillonite owing to its interlayer cations can shift the equilibrium curve toward lower stability conditions [198]. Besides, the montmorillonite surface by providing nucleation sites can enhance CH₄ hydrate nucleation in which the montmorillonite interlayer produces stable CH₄ hydrate with H₂O and a basal montmorillonite d-spacing of 24 Å [199]. Also, the increase of basal d-spacing may improve the diffusion coefficients for guest molecules [200]. Since experimental techniques are associated with some limitations of the spatial and temporal resolution to obtain direct evidence of the gas hydrate nucleation and growth behaviour in contact with minerals, MD simulations have therefore emerged as a useful technique for such discoveries. Simulations of the hydrated silica-water interface (CH₄+CO₂+SiO₂) and the contact angle determined that the density of water in the layers far from the surface is less than the layers adjacent to the silica. Similarly, at the interfaces of liquid and silica, the concentrations of the gases are more than regions in the gas phase [201]. Evaluations of the reaction rate of CO₂ hydrate with Na₂CO₃ showed that the carbonation of the Na₂CO₃ is not mainly affected by the kinetics of the reaction at the surface. In addition, the reaction rate of Na₂CO₃ is considerably less than K₂CO₃ [202]. Under an external electric field, the migration of NaCl solution in the calcium silicate hydrate (C–S–H) nanochannel showed that bond stability can be weakened by an external electric field. Additionally, the diffusion coefficient of ions confined in pores is around three orders of magnitude less than the mobility of ions in the bulk solution [203]. It was also reported that the cosine oscillation electric field can induce and play positive roles in both dissociation and growth of CH₄ hydrate [71].

2.8 Semi-clathrate gas hydrates

Aqueous solutions of semiclathrate hydrate formers would be an option for hydrate-based gas separation or cold storage aims. The ionic clathrate hydrates can form from ionic guests such as (TBAB) which have two distinct types of D cages occupied by gas species e.g., CO₂, N₂, and CH₄ [204–208]. Although the cage occupancy and anisotropic angular distribution of D_A and D_B cavities are dissimilar, they have nearly equal volumes. In this regard, MD simulations revealed that CO₂ and CH₄ molecules prefer to occupy the elongated D_A cavity and the regular quasi-spherical D_B cavity respectively. Although semiclathrate hydrate properties e.g. density, enthalpy, and equilibrium formation conditions have been successfully predicted by the MD approach, the number of such investigations is quite a few. By analyzing the rotation angle, for CO₂ molecules in two cavities at 100-250 K and the lattice structure of CO₂+TBAB, it was revealed that the lattice vibrations of bromide and tetra-n-butylammonium groups about their equilibrium positions are small and they are kept in place by electrostatic and van der Waals interactions with the neighbouring water molecules [209]. With the utilization of crystal XRD and MD simulations, CH₄ trapped in the two dodecahedral cavities of TBAB semiclathrate hydrate reported that cage occupancies for CO₂ and CH₄ can be 0.87 and 0.17 in D_A and 0.49 and 0.99 in D_B respectively. [208].

3 Hydrate phenomena and properties

Alongside the abovementioned MD simulations of gas hydrate systems, the other characteristics such as structural, energetic, mechanical, physicochemical properties, and dynamical behaviour of gas hydrates have always been the subject of MD investigations. As exhibited in Figure 9, there are different gas hydrate characteristics and phenomena which can be investigated through MD methods at the molecular level. In this section, the MD simulations of phenomena and properties of gas clathrate hydrates such as those indicated in Figure 9 are reviewed.

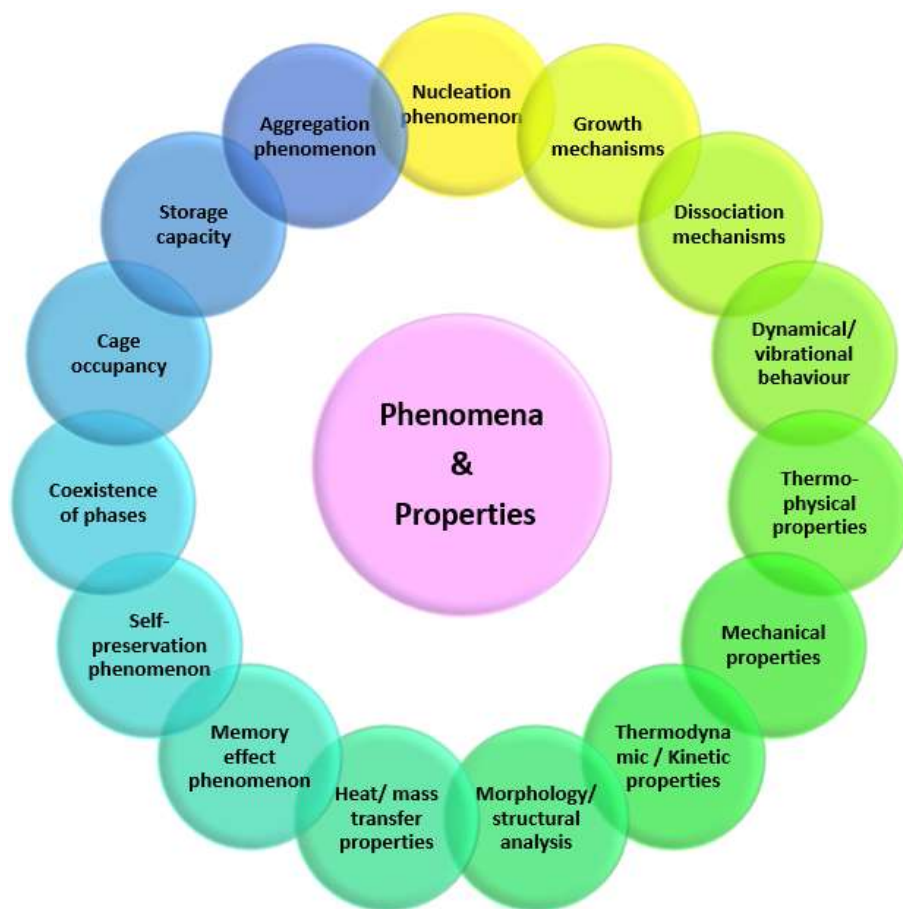


Figure 9: Studied phenomena and properties of hydrates via MD simulations.

To study the aforementioned phenomena/ properties of gas hydrates, there are several terms such as the dimension of the simulation box, type of ensemble in the final/ production phase of simulations, and simulation time MD simulations that need to be properly selected. For example, to obtain a high-ordered hydrate, an NPT ensemble and to observe a critical nucleus as large as possible, an NVE ensemble would be more proper to implement respectively [117]. The nucleation of H_2S hydrate under constant energy conditions (NVE) showed that in the early stages of the nucleation, unlike NPT or NVT which ignored the heat generated by hydrate formation, the NVE simulations can closely mimic experimental conditions [210]. After conducting the simulations, the final step will be the calculations of analysis parameters.

Table 1 summarizes some examples of MD simulations studied on different gas clathrate hydrate systems.

Table 1: Molecular simulation studies on different hydrate systems.

System	Investigated phenomenon/ properties	Analysis parameters	P (MPa)	T (K)	Hydrate supercell/ Simulation box (nm ³)	Production run time (ns)	Ref.
CO ₂	Relationship between hydrate occupancy and dissociation	<ul style="list-style-type: none"> ▪Density ▪Energy ▪Release rate ▪Dissociation of empty/ filled cages ▪Interface velocity 	3	270- 290	4×4×4 (10.8×4.8×4.8)	50 ns in <i>NPT</i>	[112]
	NEMD, EMD simulations of the thermal-driven breakup of hydrate	<ul style="list-style-type: none"> ▪Fluctuations autocorrelation function (ACF) ▪No. of hydrate-like CO₂ molecules 	20	260- 320	4×4×4 (9.6×4.8×4.8)	10 ns in <i>NPT</i>	[211]
	Three coexistence phases	<ul style="list-style-type: none"> ▪Cage occupancy ▪P-T prediction ▪potential energy ▪Density ▪Cross interaction parameters 	2-500	260- 295	2×2×2 (2.4×2.4×7.2)	400 ns in <i>NPT</i>	[103]
	Homogeneous nucleation mechanism	<ul style="list-style-type: none"> ▪No. of cages ▪Growth rate ▪Nucleus size ▪Mutually coordinated guest (MCG) 	50	250- 273	(4.8×4.8×4.8)	1 μs in <i>NPT</i>	[212]
CH ₄	Surface tension at the interface of CH ₄ , water, hydrate	<ul style="list-style-type: none"> ▪Density ▪Enthalpy ▪Interaction parameters 	10	300	4×4×4 (9.6×4.8×4.8)	800 ns in <i>NPT</i>	[100]

	▪RDF					
Cross-nucleation and equilibrium conditions of the empty and filled lattice	<ul style="list-style-type: none"> ▪No. of cages ▪Surface density of cages ▪Persistence of individual $5^{12}6^3$ cages 	10	240-320	<ul style="list-style-type: none"> $6 \times 6 \times 6$ $(14.4 \times 7.2 \times 7.2)$ 	400 ns in <i>NPT</i>	[93]
Self-preservation mechanism	<ul style="list-style-type: none"> ▪Density ▪Melting rate ▪Time evolution of total energy 	1-20	200-275	<ul style="list-style-type: none"> $4 \times 4 \times 4$ $(9.6 \times 4.8 \times 4.8)$ 	100 ns in <i>NVT/E</i>	[55]
Relationship between grain size, mechanical instability, and fracture behavior	<ul style="list-style-type: none"> ▪Cohesive energy distributions ▪Stress-strain relationships 	10-50	203-283	<ul style="list-style-type: none"> $8 \times 8 \times 8$ 	10 ns in modified <i>NPT</i>	[73]
Mechanical properties of monocrystalline hydrate, Young's modulus, and strain rate	<ul style="list-style-type: none"> ▪Tensile stress ▪Compressive stress 	10	203-283	<ul style="list-style-type: none"> $8 \times 8 \times 8$ $(9.6 \times 9.6 \times 9.6)$ 	1 ns in modified <i>NPT</i>	[89]
Interfacial free energy, and stress at the crystal-liquid interface	<ul style="list-style-type: none"> ▪Density ▪Potential energy ▪Interfacial tension ▪Excess enthalpy/ entropy ▪RDF 	30	271	<ul style="list-style-type: none"> $4 \times 4 \times 10$ $(4.8 \times 4.8 \times 20.0)$ 	20 ns in <i>NP_{NAT}*</i>	[46]
Hydrate growth based on non-equilibrium thermodynamics	<ul style="list-style-type: none"> ▪Potential energy ▪No. of guest molecules ▪RDF ▪hydrate growth rate 	5-25	240-270	<ul style="list-style-type: none"> $3 \times 3 \times 3$ $(3.6 \times 3.6 \times 11.0)$ 	10 ns in <i>NPT</i>	[213]
Mechanical failure of monocrystalline hydrate	<ul style="list-style-type: none"> ▪Tensile stress 	10-20	120-325	<ul style="list-style-type: none"> $4 \times 4 \times 12$ 	100 ns in <i>NPT</i>	[91]
Clathrate hydrate nucleation in the presence of a growing ice front	<ul style="list-style-type: none"> ▪Density ▪Interfacial transition layer ▪Potential energy 	500	250-270	<ul style="list-style-type: none"> $8 \times 8 \times 8$ $(9.1 \times 10.6 \times 36.0)$ 	10 ns in <i>NPH</i>	[214]

		▪Competing growth					
	Thermo-physical properties of hydrate using AIMD simulations	▪Thermal expansion ▪Compressibility ▪Total energy ▪Heat capacity ▪Elastic moduli	1-400	210-323	1×1×1	63 ps in <i>NVT</i>	[215]
H₂	Special structures of H ₂ hydrate (C ₀ and sT')	▪MSD	200-1000	140-260	-	15 ns in <i>NPT</i>	[115]
	Vibrational properties, and structural changes	▪Enthalpy ▪Lattice constant ▪P-T prediction ▪Density ▪RDF	3-100 GPa	300	-	15 ps in <i>NVT</i>	[216]
	Inter-cage hopping in II clathrate hydrate	▪Diffusion coefficient ▪Activation energy ▪Cage occupancy ▪Inter-cage hopping	150	80-240 K	2×2×2 (3.4×3.4×3.4)	50 ns in <i>NVT</i>	[217]
H₂/ deuterium (D₂)	Molecular scape from hydrate phase	▪No. of guests ▪Activation energy ▪Cage occupancy ▪Binding free energy	200	150-180	2×2×2 (3.4×3.4×3.4)	1 μs in <i>NPT</i>	[218]
H₂S	Rapid nucleation of hydrate	▪Order parameter ▪Nucleation rate	50	230-265	3×3×6	80 ns in <i>NVE</i>	[210]
C₂H₆/ C₃H₈	Decomposition and encapsulation energies	▪RDF ▪Dissociation Enthalpy	0.1	273	3×3×3 (sI) 2×2×2 (sII)	1 ns in <i>NPT</i>	[219]
CO₂+CH₄	Replacement phenomenon	▪Cage-guest Distance Coordination number	3.2-6.0	260-280	2×2×4	5 ns in <i>NPT</i>	[120]
	Replacement and co-growth of CO ₂ and CH ₄ hydrates	▪No. of adsorbed guests ▪Cage composition ▪No. of cages	2-10	250-275	3×3×3	1 μs in <i>NPT</i>	[220]

	Heat capacity, thermal expansion coefficient, and compressibility	<ul style="list-style-type: none"> ▪ Lattice parameter ▪ Density ▪ Compressibility ▪ Thermal expansion ▪ Specific heat capacity 	10-100	271	$3 \times 3 \times 3$ $(3.6 \times 3.6 \times 3.6)$	3 ns in <i>NPT</i>	[48]
CH₄+CO₂+N₂	Effect of N ₂ on the process of CO ₂ /CH ₄ replacement	<ul style="list-style-type: none"> ▪ No. of guest molecules in phases ▪ No. of rings ▪ Density 	6	280	$3 \times 2 \times 2$ $(6.5 \times 2.4 \times 2.4)$	2 μs in <i>NPT</i>	[221]
CH₄+C₂H₆+C₃H₈	Interfacial tension and behavior of mixed hydrocarbons and water	<ul style="list-style-type: none"> ▪ Density ▪ Angular distribution ▪ Interfacial tension ▪ No. of adsorbed molecules 	10-20	275-298	$(3.2 \times 3.2 \times 10)$	5 ns in <i>NVT</i>	[222]
CH₄+SDS+CAPB	Impact of surfactant on hydrate formation	<ul style="list-style-type: none"> ▪ No. of guest and host molecules ▪ Distribution of asphericity, acylindricity ▪ ACF 	0.1-5	275-298	-	3 μs in <i>NPT</i>	[140]
H₂ + LMGS	Storage capacity of sH hydrate	<ul style="list-style-type: none"> ▪ Order parameter ▪ Potential energy ▪ Density ▪ No. of guests ▪ Storage capacity 	70-110	230-260	$(5.0 \times 4.2 \times 3.0)$	2.3 μs in <i>NPT</i>	[223]
HFC-41/ CH₄+pinacolone	Anisotropic expansion of the sH hydrate lattice	<ul style="list-style-type: none"> ▪ Lattice constant ▪ RDF ▪ MSD ▪ OACF 	0.1	90-180	$3 \times 3 \times 3$	25 ps in <i>NPT</i>	[224]
CH₄+NaCl/KCl/ CaCl₂	Hydrate dissociation, kinetic energy, and transport parameters	<ul style="list-style-type: none"> ▪ RDF ▪ MSD ▪ Diffusion coefficient ▪ potential energy 	-	273	$2 \times 2 \times 3$	600 ps in <i>NVT</i>	[225]

CH₄+Ethanol+	Hydrate growth in the	▪Order parameter	0.1	100-	2×2×2	500 ps in <i>NPT</i>	[226]
1-Propanol/	presence of inhibitors	▪Gas uptake		250			
2-Propanol		▪RDF					
CH₄+	Hydrate growth in the	▪RMSD	15	250	7×3×2	30 ns in <i>NPT</i>	[174]
Antifreeze	existence of antifreeze	▪RMSF			(8.4×3.6×9.8)		
proteins (AFPs)	proteins	▪No. of H-bond					
		▪Order parameter					
CH₄+ SiO₂	Thermal conductivity of	▪Thermal conductivity	1	253-	2×2×3	1 ns in <i>NVE</i>	[227]
	hydrate in porous media	▪Potential energy		273			
CH₄+clay/	Effects of impurity	▪Density	18	278-	-	100 ns in <i>NPT</i>	[197]
kaolinite/ silica	nanoparticles on hydrate	▪Adsorption energy		298			
nanoparticles	nucleation						
CH₄+ Na-MMT	Formation and	▪Order parameter	50	250-	10×6×1	3 μs in <i>NPT</i>	[228]
+ fatty acids	dissociation of hydrate in	▪No. of H-bond		303			
	porous media	▪No. of cages					
		▪RDF					
		▪Diffusion coefficient					

*: isothermal–isobaric–isointerface area (NP_NAT) ensemble

3.1 Hydrate nucleation and growth

Understanding the phenomena of clathrate hydrates can promote hydrate-based applications [229]. A good illustration of this is a large number of discussions of the hydrate nucleation phenomena [230,231]. Generally, order parameters to analyse and trace the hydrate nucleation are utilized. Analyzing 7 classes of order parameters with a total of 33 individual variants for the nucleation mechanism of hydrate of THF as a water-soluble guest determined that the approximations of order parameters that provide the transition state (TS) and reaction coordinate, based upon water structure are more appropriate than those which are based on guest structure [232]. Also, it was reported that survival probability and mean first-passage time (MFPT) methods can accurately calculate both the CH₄ hydrate nucleation rate and critical

nucleus size [70]. Using a coarse-grained method for self-assembled monolayers (SAMs), the impacts of hydrophobic and hydrophilic surfaces on THF hydrate nucleation showed that with a homogeneous mechanism in CH₃-SAM and OH-SAM, the nucleation rate in OH-SAM due to its higher bulk guest concentration can be faster [233]. To study the structure of clathrate-ice interfaces, conducting four different clathrate and ice nucleation methods revealed that the interfacial transition layer between clathrate and ice is always disordered and could assist in the heterogeneous nucleation of clathrates from ice [214]. A combination method of coarse-grained mW model and the forward flux sampling (FFS) showed that in the vicinity of the water-CH₄ interface, hydrate nucleation occurs and then the transition from amorphous to a crystalline structure accelerates its rate [90]. Compelling evidence suggested that hydrates can nucleate through multiple pathways in which direct formation to the globally stable crystalline phase is one of them [234]. The face-saturated incomplete cage analysis (FSICA) method can also be used to identify all possible face-saturated cage compositions in a system [37]. Via examining mutually coordinated guest order parameter (MCG-OP), it was found that both host and guest structuring are crucial to accurately describe the hydrate nucleation [67]. Since MCG-OP considers the effects of both guest and solvent molecules, it can be compatible with defective or hitherto undetected cavity types [235].

Figure 10 displays a crystallization pathway of clathrate hydrates. As is shown, the crystallization mechanism can be divided into three consecutive stages. The first step starts with the generation of blobs and the half-cages in which solvent-separated guest molecules and dilute solution are in equilibrium. The next steps are the organization of clathrate cages leading to the amorphous structure and finally the formation of amorphous maturation [236]. Previously, the nucleation mechanisms of sI and sII clathrate hydrate were also investigated by employing the coarse-grained model which demonstrated that the current process occurs first for the small cages followed by large ones [236].

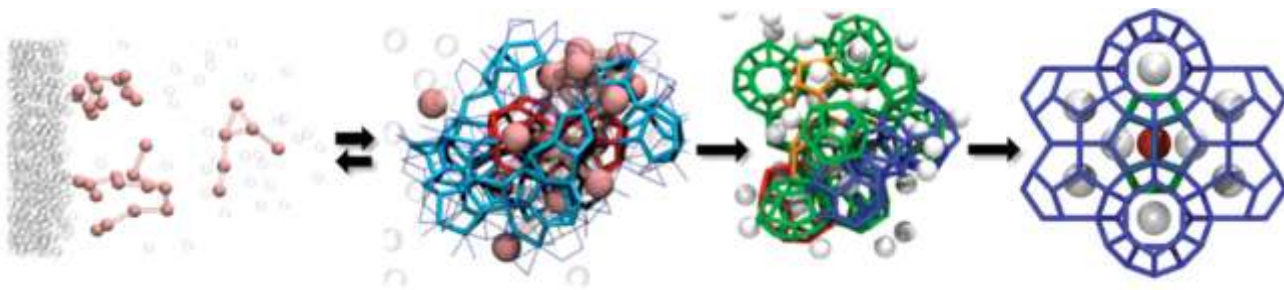


Figure 10: Multi-step blob mechanism of clathrate hydrate crystallization pathway [236].

The nucleation path in the other study was found to occur in four stages: an increase in the concentration of solvated CH_4 via diffusion, formation of an unstructured cluster, a decrease in water content of the nucleus, and finally reordering process of solvated CH_4 , and water molecules (consistent with the blob hypothesis) [96]. Hence, the hydrate growth rate was found to be due to CH_4 solubility and diffusivity in a liquid phase and CH_4 adsorption by incomplete water cages at the solid-liquid interface [66,237]. The effect of hydroxylated silica on CO_2 kinetic hydrate formation revealed that the process of CO_2 hydrate nucleation takes place in 3 stages: the formation of an ice-like layer; intermediate structure and motif layer; and finally, the nucleation seeds can form [238]. With the use of massively parallel molecular dynamics for the formation of CH_4 hydrate, two distinct steps were recognized for this process: first, the concentration of solvated CH_4 clusters increases through penetration from the water- CH_4 interface. Then, consistent with the blob hypothesis, the process of re-ordering solvated water and CH_4 molecules takes place [94]. To probe hydrate nucleation using forward and backward flux sampling it was demonstrated that the half-cage order parameter describes well the reaction coordinates of hydrate nucleation. Results also suggested that besides the two-step-like mechanism, there exist multiple active transition pathways for hydrate nucleation [95]. Hence, the clathrate nucleation of hydrophobic guests comparable to CO_2 and CH_4 has suggested that destabilizing the formation of blobs could be a good strategy to apply for the presentation of the nucleation stage [239]. By conducting a two-step method for clathrate structural recognition to identify the nuclei and analyze the relative position of cages, it was observed that 80% of

the enclathrated CH₄ molecules at the transition state are contained in sII clathrate-like crystallite. To determine the relative propensity to form sI versus sII motifs, dissipation of the nucleation energy is also essential [240]. It was revealed that CH₄ subcritical clusters can be formed at the initial stage but a higher time is required to aggregate and give the hydrate nucleus which indicates that the aggregation stage is a vital controller of hydrate nucleation [241]. MD simulations also showed the phenomenological similarities between crystallization and protein folding processes [125]. Moreover, high-level kinetic similarities between pure and corresponding mixed hydrate cannot be a reliable basis for predicting the composition of early-stage mixed hydrate nuclei [242]. Somewhat recently, the rugged funnel-shaped potential-energy landscapes and the consistency of structural biased dynamics associated with hydrate nucleation were also studied [125,243]. For the heterogeneous nucleation process, a three-step such as induce, promote, and nucleate mechanisms was suggested to occur. The heterogeneous nucleation of CH₄ hydrate from an aqueous CH₄ solution can be facilitated by the interface of hexagonal ice [244]. Nucleation of gas hydrate with high CH₄ super-saturation without necessarily allowing enough time for the structure to anneal may enable water cavities to form more quickly [36]. Nucleation of CH₄ hydrate indicated that the stability of the initial formation of sII in comparison with sI would be higher [245]. In addition, the equilibrium conditions of the empty and filled lattice of sI and sII CH₄ hydrates showed that the cross-nucleation from sI to sII or vice versa needs the formation of 5¹²6³ which is not native to both structures but plays a crucial role [93]. Previously, homogeneous and heterogeneous nucleation of CH₄ hydrate using the six-site model for water molecules were analysed that indicated at the early stage of hydrate formation as amorphous intermediate hydrate structures, a set of seven cage types (5¹², 5¹²6², 5¹²6³, 5¹²6⁴, 4¹5¹⁰6², 4¹5¹⁰6³, 4¹5¹⁰6⁴) can be formed in the nucleated solids [78]. Probing the C₂H₆ hydrate nucleation showed that the full cavity most likely to form first is 4¹5¹⁰6² which then transforms into both 5¹² and 5¹²6² cavities [111].

Explorations of CO₂+CH₄ hydrate nucleation highlighted that stability of CH₄+CO₂ nanobubbles in the water phase and the nucleation rate may be affected by the difference in hydrophobicity between CO₂ and CH₄ molecules. In addition, the temporary formation of metastable cages such as 4¹⁵10⁶², 4¹⁵10⁶³, and 4¹⁵10⁶⁴ was verified [246].

The role of operating conditions is also determinative. The melting temperature and clathrate growth at 260 K specified that a 37% decrease in the rate of CO₂ hydrate growth with increasing pressure from 3 to 100 MPa can be observed while the pressure increase can facilitate the growth rate of CH₄ hydrate [114]. In addition, increasing the pressure is less effective than lowering the temperature in promoting hydrate nucleation because the latter induces more water cages to form while the former makes them less prevalent [62]. Additionally, the crystalline structure may grow faster once the sub-cooling is reduced. Also, at moderate temperatures, transient coexistence of crystalline sI and sII may take place [68].

Homogeneous nucleation mechanism of CO₂ hydrate through transition path sampling at 260-273 K revealed that amorphous structures with 4¹⁵10⁶² cages at low temperature are predominantly formed which is the most abundant, however, increasing temperature above 265 K can change the 4¹⁵10⁶² to 5¹²6² cages and help the formation of sI crystalline [212]. In addition, during the nucleation stage, 4¹⁵10⁶² cavities were found to be the most popular type in the amorphous [247]. By simulating the amorphization of THF hydrate from 1.1 to 1.2 GPa and amorphous densification between 130 K and 170 K, it was found that repulsive interactions between guest and water molecules result in holding a cage structure around the guest [248]. Exploration of the amorphous CO₂ hydrate in a water-CO₂-silica three-phase system revealed that the nucleation and growth of the nucleus occur mostly near the three-phase contact line. Also, the SiO₂ surfaces act as a stabilizer to prolong the lifetime of hydrate cavities. For large cavities, the translational formation of the cage is formed before the rotational structure. However, these occur at

the same time for small cavities [249]. Simulation for sI and sII empty hydrate lattices at 150-50 MPa to evaluate their meta-stability, and growth mechanism revealed that the empty lattice of sII hydrate can be stable at 275 K and 130 MPa. In addition, the empty sI lattice can nucleate guest-free sII with superior stability. Also, more stability of empty lattice of water clathrates than ice at < 130 MPa and higher stability than liquid at < 275 K and negative pressure were confirmed [250]. Based on H₂+CH₄ hydrate simulations, the fastest growth rate can be observed at 250 K and 50 MPa. Additionally, at the constant temperature, the more the pressure results in the higher the crystal growth. Also, the temperature above 240 K is directly and inversely proportional to cage occupancies of CH₄ and H₂ molecules respectively [251]. The growth rate and storage capacity of sII H₂+THF double hydrate at 230-250 K and 10-110 MPa specified that the growth rate of the THF hydrate at 250 K and 50 MPa can be enhanced by the existence of H₂ but THF molecules are the main controller of the growth process. Moreover, the increase of pressure leads to the multi-occupancy of large cages, e.g. triple H₂ molecules can occupy the cages [132]. By analyzing CO₂ storage capacity at 289 K and 2.5 MPa, it was confirmed that, unlike pure CO₂ hydrate, CO₂+THF hydrate can be fully stable as displayed in Figure 11. Furthermore, the growth evolution showed that THF significantly boosts CO₂ diffusion at the hydrate-liquid interface [252].

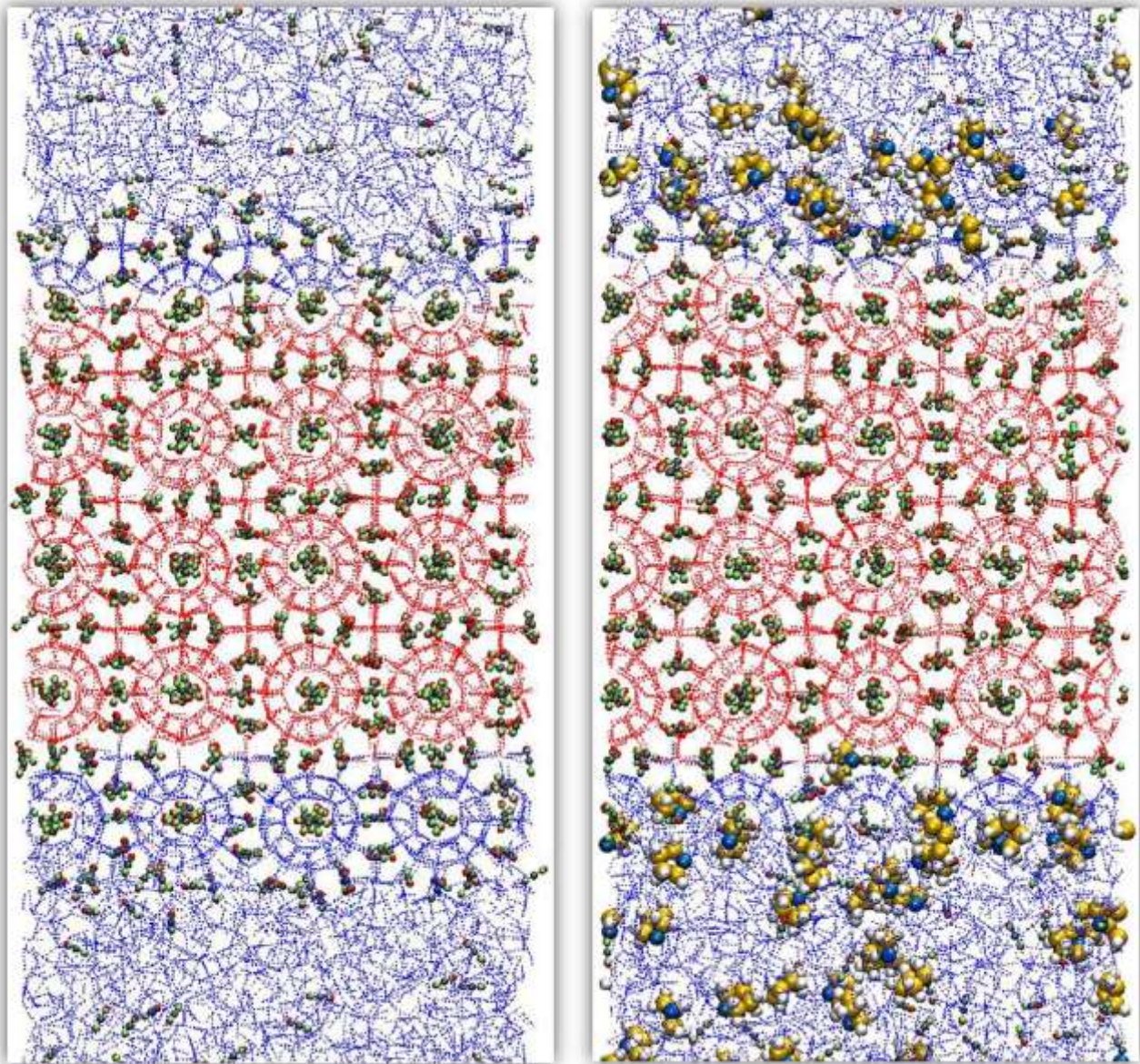


Figure 11: Final trajectories of the hydrate systems including CO₂ (left), and CO₂+THF (right) in the aqueous phase [252].

Interestingly, investigation of CH₄ hydrates in the presence of cyclic organic compounds (COC) such as cyclohexane, cyclobutanone, and tetrahydropyran revealed that replacing COC molecules with CH₄ results in reducing the unit cell stability. Also, the interface of heterogeneous crystal growth of CH₄ hydrate indicated a strong affinity for CH₄ molecules. In addition, the growth rate for hydrate crystals

was found to be four times higher than ice [253]. Analyzing the impact of high solvated CH₄ concentration on the nucleation of CH₄ hydrate revealed that a very rapid increase of nucleation rate with solvated CH₄ concentration can be observed which proved that even beyond bulk super-saturation, CH₄ molar fraction is a key that triggers the homogeneous nucleation of clathrate [88]. Analysing the solid-liquid water transition in the presence of a low concentration of CH₄ molecules as a hydrophobic guest determined that hydrate growth is dependent on the rate of empty growth lattice and filling of the cages by guest diffusion. Hence, CH₄-encaged molecules may act as a catalyst [32,254]. Commonly, thermodynamic properties of CH₄ hydrate such as entropy, enthalpy, and kinetic energy during the process of hydrate formation are in accordance with the organized water molecules in the crystal [255]. Based on the non-equilibrium thermodynamics it was also certified that the process of clathrate growth starts by dissolving CH₄ molecules to the interface and then arranging water molecules [213]. MD simulations for spontaneous nucleation and growth of CH₄ hydrate showed that without an energetically unfavourable interface, uncommon 5¹²6³ cavities facilitate the structural coexistence of two dominant hydrate types (sI and sII) [77]. Also, a simulation of the effect of CH₄ adsorption on the lifetime of a dodecahedral water cluster (DWC) showed that the adsorbed CH₄ molecules can prolong the DWC lifetime, thereby being more adapted for the DWC [256]. The DWC itself vigorously adsorbs CH₄ molecules. As an inherent driving force, shell radii between 6.2 and 8.8 Å is the strong net attraction in DWC which may control the hydrate formation [257]. It was revealed that the stability of the hydrate nucleus on the surface deems to be dependent on the affinity for guest molecules. Decomposition of the hydrate nucleus can also be induced by a strong affinity for guest molecules. [258]. The controllers of cross-nucleation between 4 recognized polymorphs of Frank-Kasper clathrate specified that the sequence of intrinsic growth rate between crystalline polymorphs is sII > TS > HS-I > sI. Among non-identical polymorphs, the formation of an interfacial transition layer mostly prevents cross-nucleation to the quicker growing polymorph.

However, an increase in super-cooling overcomes this barrier [259]. The simulations of hydrate nucleation at a CH₄-water Interface indicated that although the nucleated structure was incompatible with the two most common bulk crystal structures, it contains structural elements of both of them [36]. For metastable sI and sII polymorphs of empty hydrate lattices, it was found that the detailed shape of the librational spectra, and the differences between clathrate spectra and the ice, arise from dynamical correlations between each molecule and its local area [260]. Previously, super-saturation, dynamic and thermodynamic properties of hydrate nuclei were studied using a coarse-grained mW model. It was found that under realistic conditions of formation in industry and nature, homogeneous nucleation of clathrate hydrates does not contribute to their crystallization [92]. The convergence of the nucleation rate is also dependent on the spatial distribution of the spontaneously formed hydrate seeds [90]. Probing the effects of ice on CH₄ hydrate nucleation showed that items such as high CH₄ concentration near the ice surface, H-bond between the ice surface and hydrate lattices, absorption of released heat from hydrate formation by ice, and finally the preference of cages to occur in the vicinity of ice surface rather than in bulk solution can promote the nucleation of CH₄ hydrate [261]. Also, in the direction of fluid–fluid (water-C₃H₈) phase separation, high-density fluctuations may promote hydrate nucleation [262]. Investigation of the crystal growth of water-soluble hydrate formers showed that the growth rate of THF hydrate is an order of magnitude less than that of EO and H₂S hydrates. However, the surface trapping effect which leads to the slow growth rate of THF hydrate was not observed for EO hydrate [263]. It should be noted that the growth behaviour and induction time for the gas mixture or in the presence of additives are dissimilar. The sequence of induction time for pure guest hydrate nucleation at a fixed solution composition and the temperature was found to be C₂H₆ < C₃H₈ < CH₄ < H₂S. Additionally, there is a strong nonlinear correlation between induction time and guest composition of mixed hydrates [124,264]. The effect of single-walled carbon nanotubes (SWCNTs) on selective separation as well as capturing CO₂ molecules

from the mixture of CO₂+H₂ exhibited that Quasi-one-dimensional (Q1D) hydrates can be an eligible approach aiming at H₂ purification from syngas which may offer a safe and clean way to store a large-quantity of H₂ [265,266]. It was found that CH₄+THF mixed hydrate has a lower induction time than both pure CH₄ and THF hydrates. Also, at the initial stage of crystal growth, 5¹², 5¹²6², 5¹²6³, and 5¹²6⁴ cages were observed but in the later stages, a possibility of structural change from sI or sII was confirmed [267]. Insights into the effects of hydrophobic solid surfaces on the formation of gas hydrate determined that at hydrophobic surfaces, the tendency of getting more local ordering of water and interfacial gas enrichment (IGE) gives rise to the promotion of hydrate formation [65]. According to MD simulations to investigate the formation of hydrate at the gas-liquid-metal interface, it was found that the metal surface remarkably accelerates the water conversion to hydrate by facilitating the dissolution of CH₄ in THF solution. Increasing the liquid-gas interface curvature can result in a decrease in the energy barrier by 15% [268]. Based on the influence of amino acids on CH₄ hydrate growth at 10 MPa and 270 K, the presence of L-histidine in the system was found to markedly boosts CH₄ hydrate growth kinetics [269]. Also, analyzing the possible formation of CH₄ clathrate attributed to the capability of single SDS (form dimer to solvation sphere) revealed that the folding of SDS in water solutions is the crucial dynamical step to create incipient small cages when the SDS micelle cannot be formed [270]. The presence of hydrophobic nanoparticles can also influence CH₄ hydrate crystallization. With the addition of nanoparticles, a trend of hydrate nucleation on the side of the aqueous phase occurs but crystallization on the solid substrate may not be observed [271]. It was found that nanoparticles can facilitate CO₂ dissolution by improving the CO₂ migration from the bulk of the solution to the interfaces, however, it may physically block the CO₂ migration at high concentrations [272]. The spatial distributions and orientations of sH hydrates of CH₄ + TBME/ NH/ MCH showed that dissimilar to TBME, the hydrophobic interaction of MCH and NH with water restricts the host-guest contact, leading to lower initial kinetics of hydrate formation. The attraction

between water and MTBE may also perturb the formation of a hydrate cavity and curb the CH₄ occupancy [273]. The effects of 2,2-Dimethylbutane (DMB) on nucleation of CH₄ hydrate proved that DMB serves as a nucleation site that promotes the formation of CH₄ hydrate so that the presence of a pretty larger hydrocarbon in low concentration notably affects hydrate nucleation [151]. The influence of electrolyte solutions (MgCl₂/NaCl) on hydrate formation of CO₂ in terms of ions mobility, density profile, and cage content is also unique. Although the inhibition behaviour of MgCl₂ and NaCl on CO₂ hydrate growth can be observed, the inhibition effects of MgCl₂ compared to NaCl are higher [274]. This may be attributed to their inhibition mechanisms at the initial stages. The kinetic structure and thermodynamic properties of CH₄ hydrate from the aqueous phase in the existence of NaCl solution revealed that the hydrate growth rate with the addition of 2 mole% NaCl can be decreased by 30% to 50%. Also, the presence of ions affects the CH₄ cage occupancy [275]. According to the Hofmeister series, the influence of non-ionic and ionic salts (NaCl/ NH₄Cl/ guanidinium chloride (GdmCl)/ methanol) on the growth of CH₄ hydrate was explored. At low concentrations e.g. 1 wt.%, the promoting effect of a few monovalent salts was the highest for GdmCl while at high concentrations e.g. 10 wt.%, methanol showed the slowest hydrate formation kinetics [276]. The growth of CH₄ Hydrate with ethanol+1-propanol/ 2-propanol confirmed that the kinetics of hydrate formation for pure water is faster than the existence of inhibitors [277,278]. CH₄ hydrate growth +Na-MMT+ leonardite humic acid (LHA) showed that at a high concentration of LHA when the self-aggregation takes place, LHA acts as kinetic inhibition for the hydrate formation on clay minerals [279]. To explore the nucleation of CH₄ in the presence of two different kinetic inhibitors, it was revealed that the inhibition impact of KHIs (PVP-A/ PVP) on CH₄ hydrate at great sub-cooling decreases significantly and unexpected promotion behaviour can be observed. Under such conditions, KHIs decrease the mass transfer resistance, leading to the nucleation of CH₄ hydrate [280]. According to the simulation results of CH₄ hydrate in contact with seawater (water with 3.5 wt% NaCl), the inhibition

effects of electrolyte ions on nucleation, and orientation of water molecules were certified. Also, ions facilitate the mass transfer of CH₄ but suppress simultaneously the penetration of guest molecules [281]. The investigation of CO₂ hydrate near silica surfaces revealed that the nucleation of CO₂ hydrate tends to occur on relatively less hydrophilic surfaces more easily. Also, the structuring of molecules induced by amorphous solid surfaces is less ordered than that by crystalline surfaces at the initial hydrate growth [282]. The CH₄ hydrate sandwiched by hydroxylated silica nanopores showed that the growth of CH₄ hydrate in the pore centre is more than on the surfaces where a thin film of water exists. In addition, at pressures lower than that required for the growth of CH₄ clathrate in the bulk, CH₄ hydrate forms in the nano-pores [196]. The simulation of CH₄ hydrate nucleation between hydrophobic graphite and hydrophilic silica surfaces also showed that due to the adsorption of CH₄ molecules by graphite surface to form a nanobubble and induce hydrate-like water ordering by graphite near the surface, hydrate nucleation does not occur. In contrast, silanol groups on silica form strong H-bonds which can stabilize the incipient hydrate and facilitate the formation of CH₄ hydrate [283]. Also, the effects of impurity nanoparticles (clay/ kaolinite/ silica nanoparticles) on CH₄ hydrate nucleation highlighted that in a similar fashion to heterogeneous ice nucleation, impurity particles for water-soluble guest molecules act as a promoter but CH₄ hydrate formation is mostly insensitive to the existence of impurity particles [197]. CH₄ hydrate nucleation in the presence of porous sediments showed that the dissolved CH₄ molecules migrate to the clay surface where hydroxylated edge sites of clay can facilitate the hydrate nucleation [284]. However, the external surface of clay mineral effects on CH₄ hydrate formation demonstrated that unlike the neutral-charge layer, CH₄ hydrate nucleation for clay minerals with a negatively charged layer can occur in the bulk-like region but away from the clay mineral surfaces (water-mineral interface) [285]. According to the simulations of the influence of silica and clay nanoparticles on CH₄ hydrate formation, notwithstanding the hydrate nucleation of hydrophilic molecules like CO₂ and THF which can be

promoted by adding impurity particles, the nucleation of CH₄ hydrate is not sensitive to these materials [286]. The impact of organo-minerals such as adsorbed zwitterionic glycine on the sodium montmorillonite surface (Na-MMT) during CO₂ hydrate nucleation showed that employing organo-mineral complexes increases the interface area of mineral and water to accelerate the nucleation and crystal growth stages of CO₂ hydrate [287].

3.2 Hydrate stability and dissociation

To use the hydrate applications, comprehending hydrate stability and dissociation mechanisms would be the critical objective. Previous MD studies of the relationship between hydrate occupancy and dissociation rate and interface velocity showed that identical total occupancy may result in different dissociation behaviour [112]. Also, for pressure ranges up to 500 MPa it was demonstrated that decay temperature is directly dependent on the cage occupancy [60]. MD simulation of the hydrate cluster dissociation elucidated that the kinetics of atomic dissolution is 5 times faster than crystal hydrate dissolution [42]. Additionally, the kinetic rate of hydrate dissociation was found to be remarkably higher than that in the hydrate formation [288]. Generally, smaller guest molecules than CH₄ may have the earlier hydrate dissociation. However, quick decomposition can be observed for those molecules larger than the cage diameter of sI hydrate [289]. According to the MD results of non-equilibrium adiabatic CH₄ hydrate dissociation, releasing large amounts of CH₄ near a surface increases the formation of bubbles and subsequently the rate of mass transfer [43]. MD also suggests that the dissociation process may occur in four successive stages in which the dissociation process equally undergoes small cages first and next to large cavities [290]. The first step of dissociation is the diffusive behaviour of water and cell size increase which leads to fracture of the cages. This is mainly followed by the escape of guest molecules from broken hydrate cages and then aggregating together [82]. Although the resistance of heat and mass transfer during the dissociation increases, it mostly reduces the rate of CH₄ hydrate decomposition [113].

Dissociation of partially occupied hydrates is somewhat faster than those of fully occupied hydrates [291]. Also, during the initial homogeneous stages of melting gas hydrate, the aggregation and migration of CH₄ molecules are critical [64] and until the threshold of bubble formation, hydrate exists as a metastable superheated solid [63]. Simulations of gas hydrate dissociation in sediments determined that the dissociation occurs layer by layer in a shrinking core manner. In addition, the released CH₄ molecules aggregate and subsequently evolve into nanobubbles [292]. Relative to the case where the hydrate is in contact with silica, the presence of a water layer between the silica surface and the hydrate phase increases the hydrate stability [293].

The Non-equilibrium MD (NEMD) simulations methodology to study the acoustic-propagation properties of CH₄ clathrate hydrate, and reproduce the P-wave and S-wave velocities in the elastic-response regime of sI and sII CH₄ hydrate can be accurately employed [53]. NEMD simulations for CH₄ clathrate hydrate dissociation indicated that the dissociation rate of hydrate surrounded by (50% CH₄+50% H₂O) and 100% CH₄ are nearly 30% and 55% lower than 100% pure water [79]. Analyzing thermal-driven CH₄ hydrate breakup at the water-hydrate interface proved that before the threshold of hydrate dissociation, the fluctuation-dissipation theory is valid and can properly describe the nature of the non-equilibrium [45]. NEMD and EMD simulations of the thermal-driven breakup of CO₂ hydrate at 300 K to 320 K demonstrated that Onsager's hypothesis (about the composition-dependence of corresponding decomposition rates above the melting points) is applicable for an initial period of hydrate dissociation [211]. Also, fluctuation-dissipation at the interface plays a critical role [294]. Thermal-driven break-up of C₃H₈ hydrate interfaced with liquid water also specified that the Arrhenius equation can predict the dissociation rate of C₃H₈ hydrate satisfactorily [295]. According to the dissociation of H₂+C₃H₈ hydrate by NEMD and EMD simulations using pairwise potentials, it was found that different surface-cavity terminations lead to substantial differences in initial break-up rates [296]. NEMD

simulations in a range of externally applied electromagnetic fields showed that below a certain intensity threshold, electromagnetic fields cannot bring about structural distortion or dissociation effect on bulk clathrate [297]. However, once a CH₄ molecule escapes from a distorted cavity, it is not possible to re-enter them even in the absence of static fields [298]. In the presence and absence of an electric field, release and uptake of sII neon hydrate showed that activation energies for uptake and release of neon in the absence of an electric field were 14.9 and 16.4 kJ/mol which indicated a good agreement with the experimental measurements, however, the release value in the existence of an electric field was declined to 6.5 kJ/mol [299]. By simulating the role of the magnetic field on the formation of CH₄ hydrate in the existence of micro-organisms (proteins), some evidence of oriented magnetic fields on the hydrate-formation kinetics by a prototypical aromatic peptide was confirmed [300]. Based on exploring the dissociation of N₂ hydrate within SWCNT and under the axial electric field, the electric field was found to change the orientations of water dipoles which results in altering the diffusion coefficient and hydrogen-bonding network of the water molecules [301]. Figure 12 shows the NEMD simulation of C₃H₈ hydrate in contact with liquid under the electric field. Analysis of this system revealed that an electric field under 0.7 V nm⁻¹ does not lead to dissociation of pre-existing bulk clathrates but field strengths more than that result in significant differences in the initial dissociation rates. In addition, the dissociation rates were observed to be strongly dependent on temperature [302]. It was also revealed that lower frequency and higher intensity may facilitate C₃H₈ hydrate dissociation [303]. Based on the results of an external electric field to simulate the dissociation of CH₄ hydrate using the non-polarizable models, the presence of an electric field leads to the formation of an ice Ih-type structure while in continuous simulation without the external field, the ice-like structures become disordered. This may result in the separated gas and liquid phases [76].

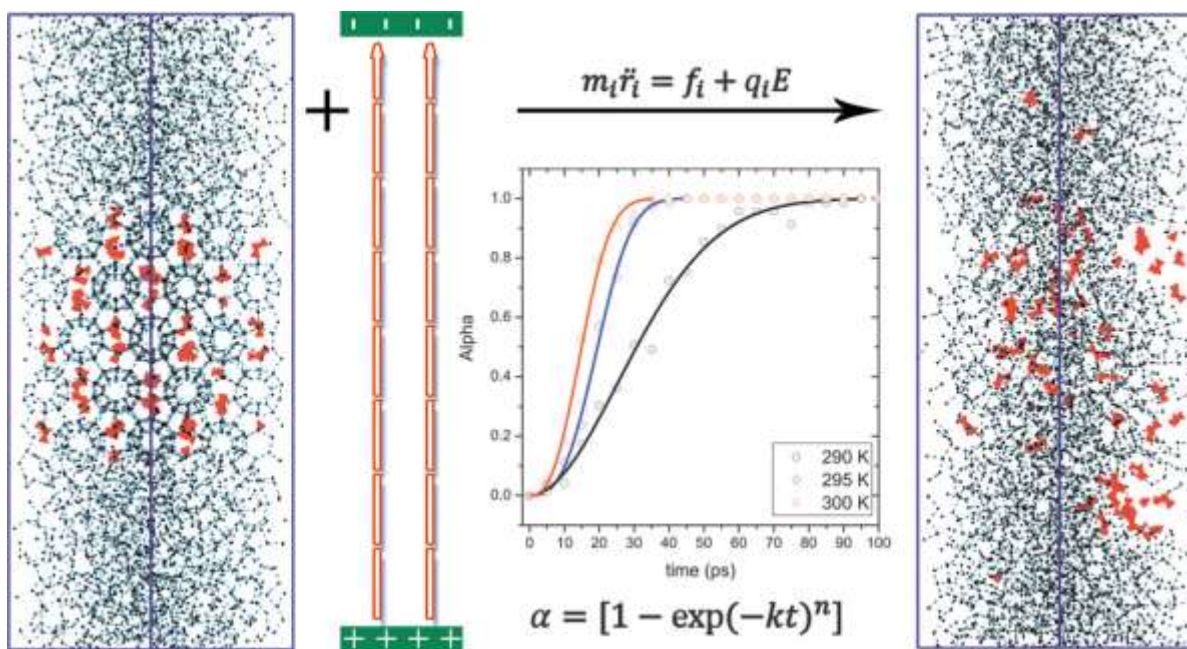


Figure 12: NEMD of both electric field-driven breakup of planar C_3H_8 hydrate interfaces with liquid water in electric field intensity [302].

Based on MD monitoring of the evolution of the CH_4 hydrate dissociation in the inclusion of two water reservoirs, it was found that the released CH_4 molecules at the initial steps reach the gas phase so that increase the gas pressure on the hydrate phase. As the hydrate dissociates, CH_4 molecules aggregate and form nano-bubbles [304]. Moreover, the slow diffusion of CH_4 molecules out of the liquid phase results in the agglomeration of CH_4 molecules and forms quasi-spherical bubbles with a radius of 11 Å [72]. Concerning the simulation of CO_2 and CH_4 bubble formation during the dissociation, it was found that when the CH_4 and CO_2 molecules occupy the small and large cages respectively, the most stable structure can be attained. In addition, the size of formed bubbles during the dissociation for each guest molecule is most likely dependent on solubility conditions of CH_4 and CO_2 in water [305]. The impact of the grain boundary structures on CO_2 hydrate at 220-310 K revealed that the stability of CO_2 hydrate somewhat above the bulk melting temperature can remain stable which confirms the relevance of thermal stability of polycrystalline hydrate to the guest's type and the grain boundaries [306]. Simulating the

decomposition and encapsulation energies of sI and sII highlighted that the encapsulation energies of guests may stabilize the cavities of sI and sII hydrate but the larger molecules give higher encapsulation energies [219].

The process and behaviour of hydrate dissociation are vigorously relevant to the encaged guest type. By analyzing the diffusion barriers, it was found that the residence of specific cavities and the overall occupancy markedly affect the dissociation of CO₂ hydrate. Moreover, unlike CH₄ and CO₂ molecules, small guests e.g. H₂ molecules due to the little penetration barrier can diffuse into the liquid phase during the early stages of the breakup of cavities [307]. MD simulations of H₂+THF hydrate during decomposition revealed that the encaged THF molecules in large cavities increase the resistance of the diffusion behaviour of H₂ molecules, however, THF serves as a strong stabilizer [308]. According to the NH₃ hydrate MD analysis, this molecule within temperatures up to 240 K gives more stable host-guest configurations than CH₄ [309]. The dissociation conditions and structural change of krypton in the existence of large molecular guests showed that the cell size of krypton increases with increasing temperature which results in clathrate distortion, small bubble formation, and krypton aggregation in the aqueous phase [310].

By evaluating CH₄ hydrate dissociation, kinetic energy, and transport parameters with the utilization of inorganic salts, it was revealed that under the same concentration, the sequence of ion's coordination number (CN) is Na⁺>Ca²⁺>K⁺ while CaCl₂>KCl>NaCl is the order of diffusion coefficients [225]. It was also shown that the generation of CH₄ bubbles in dense NaCl solutions near the hydrate interface accelerates the CH₄ hydrate dissociation [311]. However, methanol and NaCl through dissimilar mechanisms facilitate the formation of bubbles. In addition, the ions in the solution enhance the hydrophobic interactions and cause non-uniform distribution of dissolved CH₄ molecules [169]. With respect to the MD results of the effects of alcohols on the decomposition of CH₄ hydrate, it was found

that the chain length and the hydrogen number of alcohols are inversely and directly proportional to the CH₄ hydrate decomposition respectively [312]. MD investigation of CH₄ hydrate decomposition with the addition of methanol indicated that up to 10 MPa and temperatures above 280 K, the effect of pressure would not be tangible. However, the combination of increasing temperature, reducing pressure, and the addition of methanol notably increase the rate of hydrate dissociation [313]. The inhibition capability of alcohols on gas hydrates was found to be disparate [314]. By probing the C₃H₈ hydrate dissociation mechanism with methanol it was found that the hydroxyl and methyl groups of methanol create H-bonds with water molecules that destroy the original H-bonds of water molecules in the hydrate. Also, engaging methanol molecules in small cages may enhance C₃H₈ diffusion behaviours and shorten the decomposition time of C₃H₈ hydrate [315].

3.3 Hydrate cage occupancy and storage capacity

The fraction of gas adsorption into the cages is of great importance for hydrate-based gas separation and gas storage. Results of cage rigidity and the maximum/ optimum cage occupancy for various types of cages using quantum calculations (MP2, M05-2X, and DFT-D) revealed that the maximum and the optimum number of CO₂ molecules that can occupy the cavities are: one and one for small cages (5¹²) of all clathrate hydrates, two and one for sI large cages (5¹²6²), two and two for sII large cages (5¹²6⁴), two and one for sH medium cages (4³5⁶6³) and seven and five for sH large cages (5¹²6⁸) [316]. Intermolecular potentials using *ab initio* quantum mechanical to determine the reference energy/ chemical potential of sI CO₂ hydrate indicated that the filled fraction for small and large cavities can be around 32%~51% and 98% respectively [317]. Interestingly, by decreasing the fractional cage occupancy from 95% to 85%, the decomposition rate was found to be increased by 30% [318]. However, the size of guest molecules not only affects the fractional occupancy but also changes the unit cell value [61].

Based on free-energy barriers and profiles of H₂ hydrate, the energy barriers dramatically decline with increasing the occupancy of small cages by H₂ molecules [319]. In addition, the free-energy barriers for H₂ molecules from 1 to 5 molecules per large cavity showed a linear decrease for 1 to 3 molecules but become larger for 4 molecules [320,321]. MD simulations indicated that for pressures above 400 MPa, a small number of the large cavities can fill with five H₂ molecules [322]. However, the maximum cage occupancies of H₂ molecules in sI and sII large cages (5¹²6⁴ and 5¹²6²) and small cages (5¹²) were determined as eight, six, and two while in the optimum case were two, two, and one [323]. Also, the negative interaction energy of sI H₂ hydrate up to 50% large cage occupancy revealed that H₂ hydrate at 150 K and 10 MPa can be metastable which would be competent for H₂ storage [324]. MD simulations for sH hydrate storage capacity highlighted that the optimum number of guest molecules in the large cavity for CH₄, C₂H₆, C₃H₈, n-butane, and n-pentane were determined at 4, 3, 2, 2, 1 respectively [325]. In addition, hydrate systems with small empty cages would have higher stability than that with large empty cages. Also, hydrate stability in small cages was found to be less sensitive to CO₂ molecules compared to CH₄ [326]. Somewhat more recently, the effect of cage occupancy, pressure, and the temperature was investigated to reveal the thermos-physical properties of CH₄ hydrate in the existence of methanol. Cage stability is directly dependent on temperature and inversely proportional to the pressure and cage occupancy. Studying 100% to 75% fractional occupancy at a certain condition indicated that the lower the cage occupancy, the higher the diffusion coefficient [327]. According to MD analysis, the rate of fractional occupancy in the small cavities of sI hydrate at the pressure ranges below 1 MPa was found to be insignificant but gradually increases with elevating the pressure [328]. Recently, the intra-cage behaviour of guest molecules in doubly occupied large cavities of sII hydrates at 100 K was probed by AIMD simulations in which the qualitative consistency of tetrahedral sites with the neutron scattering classical diffusion findings was confirmed [329].

MD investigation of the occupancy and growth of binary H₂+THF clathrate hydrate at 50 MPa and 304-333 K determined that with increasing the super-cooling, more large cages can be filled with H₂ molecules but it does not affect the small cavities [330]. Investigation of the structural stability of sII H₂ hydrate determined that the increasing temperature reduces the optimum occupancy of large cages [331]. Previously, migration of H-radicals and energy barriers calculated at the MP2 level showed that the compatible calculated H-radical migration rates with the actual migration rates [332]. According to the diffusive properties of inter-cage H₂ migration in H₂ and H₂+THF hydrates at 5 MPa and 200-260 K it was found that H₂ migration does not occur. Also, the diffusivities of H₂ in H₂+THF hydrate are an order of magnitude lower than that of pure H₂ hydrate [333]. Dynamical cage behaviour and H₂ migration in both H₂ and H₂+THF hydrates at 200-250 K showed that cage hopping events can be facilitated by temporary openings of small-cage faces with the reformation and relaxation of stabilizing H-bonds [334]. Also, the inter-cage hopping in sII clathrate hydrate showed that although the small cavities remained with only one guest molecule, the large cage occupancy with two and three H₂ molecules appeared to be the most stable. The activation energy for guest diffusion was found to be the lowest and the highest for the 4 and 1 occupancy models respectively [217]. Figure 13 exhibits the hopping and exchange for CH₄ molecules between the bubble and large cages. As is shown, initially 2 CH₄ molecules occupied cages A and C. Then, the guest molecules hop into the neighbour cages (e.g. cage B). However, the exchange of CH₄ molecules between the bubble and the 5¹²6³ cavities is different. During guest hopping between the hydrate cage and bubble, a water molecule in the 5¹²6³ cavities was replaced by another water molecule. In addition, the diffusion rate during the hopping process of CH₄ hydrate growth was estimated to be in the order of 10⁻⁹ to 10⁻⁸ m²/s, which is 3~4 order of magnitude faster than that during hopping between one-occupied and empty hydrate cages [335].

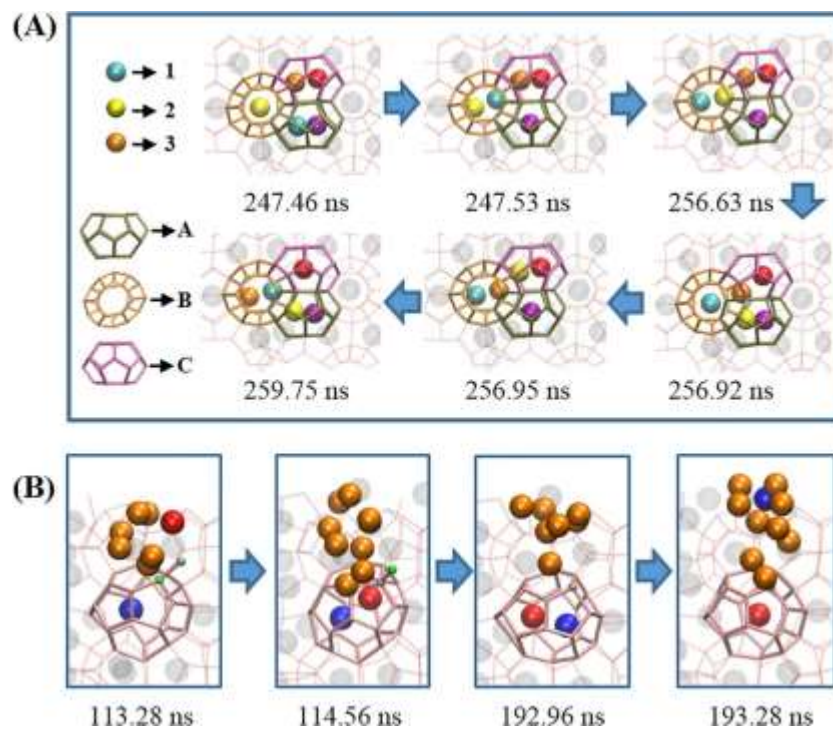


Figure 13: The hopping of CH₄ molecules: (A), the hopping and replacement for CH₄ molecules between three large cages; (B), the exchange of CH₄ molecules between the bubble and large cage [335].

MD analysis of H₂+SF₆ hydrate revealed that SF₆ molecules inhibit the diffusion of H₂ molecules and tend to occupy close to 100% of large cages [131]. Newly, the storage capacity of H₂ in the H₂+THF and H₂+MCH hydrates as a function of temperature and pressure was evaluated. Also, the small cage occupancy is directly proportional to the thermal expansion coefficient and unit cell volume [336]. It was revealed that the cage occupancy is directly and inversely proportional to the pressure and temperature respectively. H₂ storage capacity (wt.%) can be reduced by increasing the molecular weight of the promoter [337]. NEMD simulations to study the H₂ storage in C₃H₈ hydrate revealed that the diffusion coefficient at 273 K is approximately 1.5 times higher than that at 260 K. Based on the experimental and theoretical estimations, H₂ content in C₃H₈ hydrate can be stored by 1.04 wt% and 1.13 wt% respectively [338]. It was estimated that for the pressure ranges over 270 MPa, double occupancy prevails with the

single occupancy of Argon molecules in large cages [339]. MD also specified that under moderate conditions, no double occupancy occurred for the small sI, sII, sH, or medium sH cages but multiple occupancies for the large cages can be observed [340]. Previously, based on MD and vdW-P theory, simulations were performed to study the CH₄ content and occupancy in various hydrates. It was proved that multi CH₄ occupancies in large cages of sH hydrate are intensively dependent on pressure and temperature. Based on MD analysis of the storage capacity of H₂+THF sII hydrate and H₂+MCP sH hydrate at 274K and up to 500 MPa, it was estimated that large cages at high pressure can be filled with up to 8 H₂ molecules whereas small and medium cages can be occupied by just a single molecule. Moreover, the capacity of pure H₂, double sII, and sH hydrates were estimated at 3.6 wt %, 1.05 wt%, and 1.4 wt% respectively [341]. To determine H₂ storage in sH hydrate at 230-260 K and 70-110 MPa, MD simulations determined that the pressure effects at 250-260 K on the H₂ storage capacity are not substantial, however, temperatures below 240 K can at least double the storage capacity of H₂ molecules. Also, the main diffusion barriers are found to be the presence of small cages on the boundary layer and the scape of H₂ molecules due to the low kinetic energy [223]. Simulations of the CO₂ storage in sH hydrate at 100 K and ambient pressure, 273 K and 10 MPa, and also 300 K and 500 MPa revealed that although a single CO₂ molecule occupies the small and medium cages, occupancy of 5 and 3 molecules in large cages at the low and high temperatures are the most favoured [342]. Cage occupancy of double hydrates is also found to be slightly dependent on the type of large guests. For example, by simulating CO₄ and different large guests, it was determined that the most and least CH₄ storage capacity can be 11.9% and 9.6% for tetrahydropyran and cyclohexane respectively [343]. However, conducting simulations on the dependency of small guest cage occupancy and LMGs showed no relevance between CH₄ occupancy in small cages and the LMGs but the stability of sH CH₄ hydrate requires more than 40% cage occupancies of small and medium cavities by encaged molecules [344]. MD simulations of the sH

H₂+MTBE hydrate at 10-200 MPa revealed that the configurational energy of the unit cell increases when MTBE molecules are replaced by H₂ molecules in the large cavities. Also, the volume and energy of the clathrate at the lower temperature are not sensitive to the number of H₂ guests in the large cavities [345]. In addition, hydrate guest occupancies in interstitial sites using DFT and MD simulations showed that by occupying interstitial sites, H₂ can be incorporated within H₂+ tert-butylamine hydrate crystal structures [346]. Also, CH₄ may be able to replace MTBE in large cavities. In the absence of large guests, theoretically pure CH₄ hydrate requires a pressure higher than 0.5 GPa to form sH clathrate [347]. The simulated double hydrates of CH₄ and LMGs at 278 K and up to 1 GPa to investigate the storage capacity determined that the lattice constant can be expanded by increasing the size of LMGs and temperature while it is inversely proportional to the operating pressure. In addition, the fractional occupancy of small cavities by CH₄ is entirely related to the type of LMGs. [348]. By computing Quantum free-energy rates of diffusion of H₂ molecules at 8 to 200 K it was shown that at temperatures lower than 25 K, the quantum rate is greater than the classical rate whereas it inversely occurs at above 25 K [81]. Theoretically, MD simulations demonstrated that the guest-free sIII clathrate can overtake sII and sH clathrate and emerge at negative pressure e.g. -583 MPa and 0 K or -341 MPa and 300 K [349].

3.4 Hydrate guest role

Understanding the structural properties of clathrate hydrates such as metastable clathrate crystals, guest role and size, interstitial defects, structural configurations, and vibrational analysis can be worthwhile either to promote or prevent clathrate formation which has been highly sought after. The elucidation of these specifications has been to some extent uncovered via MD explorations. The anomalous shift in the stretching vibration frequencies of free and guest molecules in small and large cages is not the same. The dynamics and molecular environment of guest molecules can be reflected by the changes in molecular vibrations [350]. Also, the water-guest attraction regulates the nucleus pathway in which weak attraction

along with the poor molecule mixing in the interface layer hinders the nucleus from growing in the water phase. In such systems, along with the interface, the hydrate grows but develops toward the gas phase whereas this shift does not occur for the strong attraction [351]. To predict the ^{13}C NMR powder lineshapes of the guests, classical MD simulations were performed at the 77-250 K temperature range. A limited range of motion of C_2H_4 molecules in the cages at low temperature was observed while with increasing temperature, guest molecules gain greater rotational freedom [352]. Also, estimating the ^{13}C NMR lineshape of CO_2 at a low temperature is less accurate with experiments [353]. MD simulations showed that the thermal conductivities and speed of sound for CO_2 and Xe hydrates are lower than the empty lattice or CH_4 hydrate which points out the importance of host-guest coupling [260]. However, the overall nucleation mechanism for all guests was found to be similar and multiple competing channels form the nuclei. The size of guest molecules mostly determines the structure of the nuclei rather than ordering the stable or metastable hydrate crystals or the cage composition [236]. The stability of the hydrate nucleus in the presence of the solid surface to investigate the effect of its affinity indicated that the hydrate nucleus can be stabilized by the slab with a weak affinity for guest molecules which may be correlated with the ordered water structure on the solid surface [258]. As Figure 14 (left) shows, the free energy of the guest molecules directly corresponds to affinities. Figure 14 (right) exhibits the initial configuration in that grey plate, cyan balls, silver dots, and red sticks represent the solid slab, guest molecules, liquid water, and the largest hydrate cluster respectively. Based on this Figure, the decline of interaction among the guest molecules and the solid slab weakens the free energy for the guest molecules adsorbed on the slab which is nearly equivalent to that of guest molecules adsorbed on the hydrate nucleus surface for $\varepsilon_{si} = 0.15 \text{ kcal}\cdot\text{mol}^{-1}$. This demonstrates the existence of competitive adsorption behaviour among the slab and the hydrate which notably decreases the hydrate nucleus size, leading to the nucleus decomposition. Also, a stronger affinity slab has a shorter hydrate nucleation lifetime.

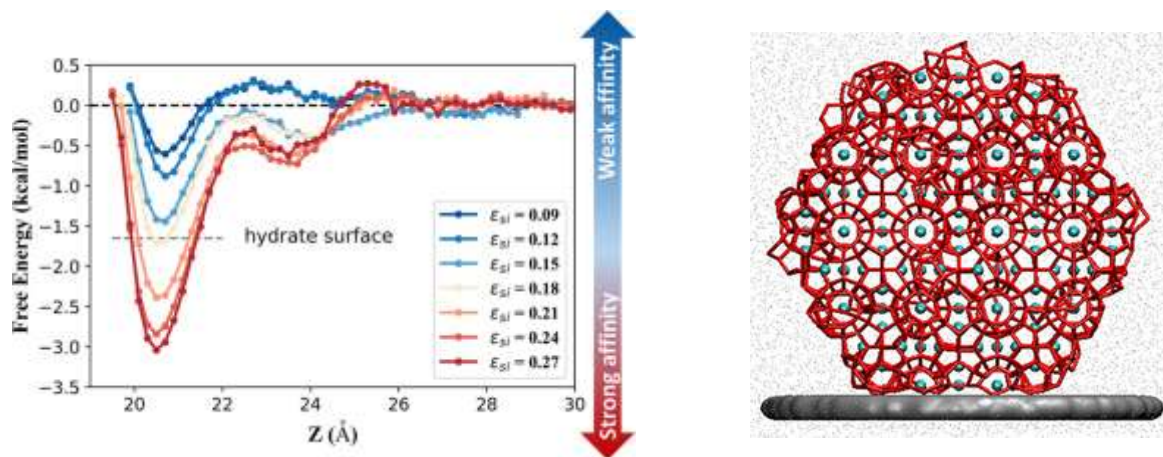


Figure 14: The effect of different guest molecule affinities ($\epsilon_{si} = 0.27\text{--}0.09 \text{ kcal}\cdot\text{mol}^{-1}$) on the free energy of guest molecules adsorbed onto the slabs [258].

Although the experimental investigations provide important insights into the hydrate structural analysis and role of guest molecules, some details are not still well-understood whereas MD simulations can systematically be employed to explore such gaps. For example, the effects of guest molecules on hydrate growth showed that the attraction of guest-water molecules may control the nucleus growth rate whereas the size of guest molecules may probably determine the hydrate structure [351].

3.5 Thermo-physical and mechanical properties of gas hydrate

Understanding the mechanical characteristics is of importance for utilizing and predicting the stability of gas hydrate formation. By employing *ab initio* density functional theory (DFT) thermo-physical properties of CH₄ hydrate can be accurately calculated [215]. Structural, energetic, and mechanical properties of CO₂ hydrate by calculating DFT and MD methodology showed that although the cage distortions are mostly isotropic, a loss in the ideal symmetry of the empty structure is mainly due to the guest molecules [354]. MD indicates the surface tension and intermolecular active forces may decrease with increasing temperature and pressure in the system [355]. Moreover, simulating CH₄ hydrate under compression and tension clarified that the tensile stress is lower than maximum compressive stress under

the same conditions [49]. Tensile stress and strain rate of calcium silicate hydrate at 300 K demonstrated that the dynamic tensile of calcium silicate hydrate directly depends on water content but it is an inverse function of strain rates [356]. Also, the shear modulus of CO₂ and CH₄ hydrate directly and inversely depends on temperature [357]. According to the MD results of nucleation of CH₄ hydrate sandwiched by elastic silica, the nucleation was found to be dependent on the elasticity of silica in which the weak elastic provides a minimum induction time but in the case of extremely weak elastic, the process of nucleation due to the fluctuation of the layers cannot easily take place [358]. Through MD simulations of acoustic and elastic properties e.g. anisotropy factor, bulk modulus, shear modulus of CH₄, C₂H₆, C₃H₈, i-butane, and empty hydrates, it was found that the repulsive impact of the guest molecule upon tension and compression can weaken and strengthen the structure of hydrate lattice respectively. The sensitivity of the large cavity to pressure is also more than the small cavity, however, H-bond and O-H bond lengths exhibit the opposite behaviour [158]. MD CH₄ hydrate simulations of strain-stress, shear strain, and elastic moduli to investigate the ideal strength under shear deformation showed that CH₄ hydrate has no dominant slip system but displays brittle behaviours in terms of its strength [359]. By evaluating the hydrate elasticity and by analyzing the parameters e.g. binding energy, shear elastic constant, and anisotropy index, it was revealed that due to a cage-like structure, gas hydrate frameworks are very isotropic. The presence of highly symmetric proton configurations is one of the reasons for the higher anisotropy of ice *Ih* [360]. The relationship between grain size, mechanical instability, and fracture behaviour of polycrystalline and monocrystalline CH₄ hydrates highlighted that CH₄ hydrate is highly sensible to cage occupancy changes and environmental conditions. In monocrystalline CH₄ hydrate, dislocation-free brittle failure was observed which showed that upon depressurization process, the polycrystalline CH₄ hydrate can be destabilized by strain-induced [73]. The determination of mechanical properties of different gas hydrates indicated that under a uniaxial mechanical load, gas hydrate's stability

is critically affected by the polarity, shape, and size of the guest molecules [361]. According to MD insights, the impact of defects on the mechanical property of sI CH₄ hydrate was also found to be substantial. Additionally, considerable fluctuations for F₃ order parameters just before hydrate structure failure were observed. Also, at a deletion rate of 9%, the mechanical property was decreased [52]. Previously, to analyse the heterogeneous crystal growth of hydrate, a number of defects were observed. Typically, independent of the applied pressure and temperature, around 20% of hydrate cages were unoccupied while some large cages trapped two CH₄ molecules [59]. Explored Heterogeneous crystal growth of H₂S hydrate demonstrated that the growth rate of H₂S hydrate is higher than CH₄ hydrate, however, in the newly formed crystal of H₂S hydrate, a relatively low level of defects was observed [362]. Also, increased gas concentration reduces the induction time for H₂S hydrate nucleation while the homogenous nucleation process features the amorphous initial formation more specifically at high super-saturations [363]. Analysis of the dynamic and structural nature of water is also confined in the quasi-two-dimensional pores which is the main binding phase in the cement. It was shown that the defective silicate chains and the interlayer calcium can render a hydrophilic interaction among the C-S-H and confined water. In this regard, the dynamical behaviour of the confined water as a glassy material at an intermediate range up to 4.2 Å was evidenced [364]. To investigate the steady-state heterogeneous crystal growth during the sI crystal formation, an unexpected kind of structural defect (consisting of 5¹²6³ cages) was previously confirmed. Additionally, an in-situ transformation of sI to sK was found to be possible albeit under prevailing operating conditions [58]. It was also revealed that guest-host H-bond leads to the formation of Bjerrum L-defects in the clathrate phase where 2 adjacent water molecules have no covalently bonded hydrogen atom between them. By conducting this simulation, it was determined that an activation barrier for the THF-water defect formation is about 8.3 kJ/mol [365]. NEMD simulations revealed that diffusion of released CH₄ from the hydrate surface during decomposition is not homogenous

and the solution phase does not necessarily remain isothermal [38]. The examination of the mechanical failure of monocrystalline CH₄ hydrate revealed that the failure may take place in two phases: gradual crack growth and quick crack propagation [91]. To investigate the mechanical properties of monocrystalline CH₄ hydrate and its intrinsic differences from ice, the impacts of guest occupancy, crystal-orientation temperature, and strain rate at 263 K and 10 MPa highlighted that although the influence of crystal orientation is not significant on monocrystalline CH₄ hydrate, mechanical strength greatly depends on temperature, strain rate, and large cage occupancy [89]. The mechanical stress–strain curves of CH₄ hydrate under three different strain rates and the directional deformation are presented in Figure 15.

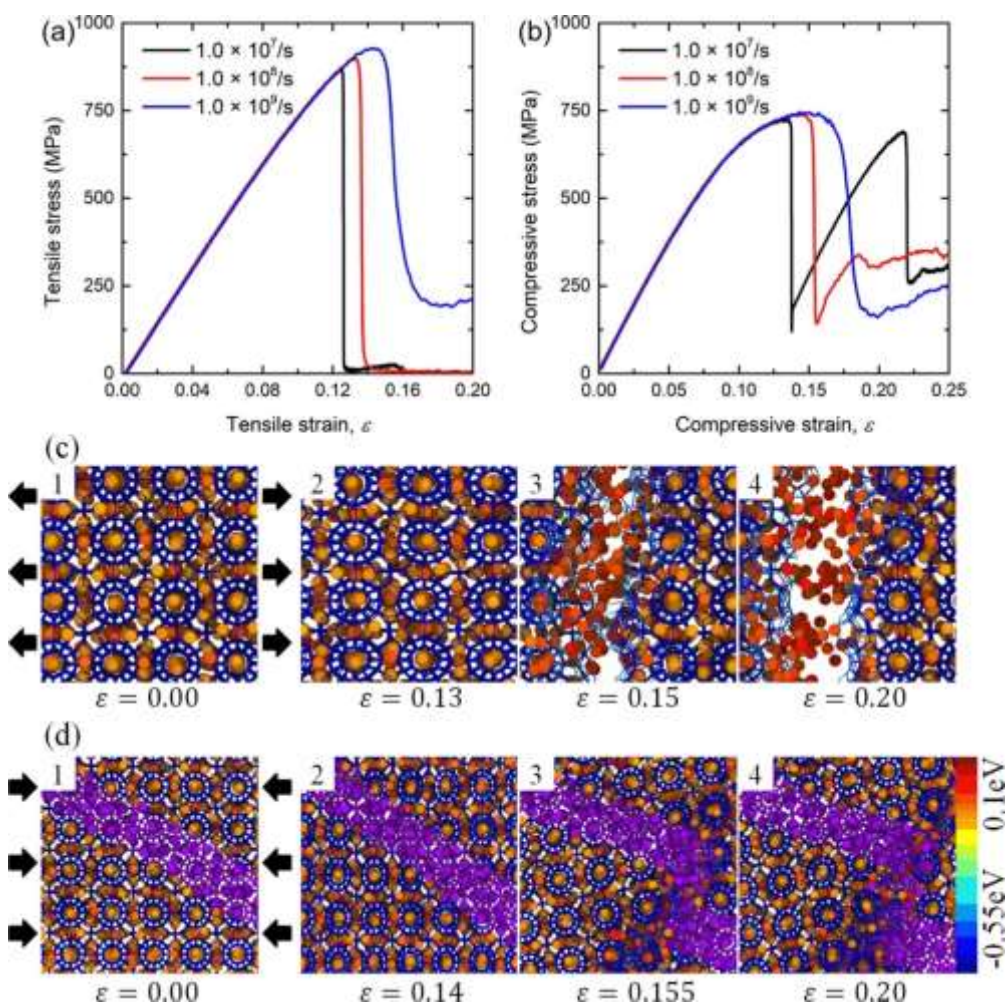


Figure 15: Mechanical properties of CH₄ hydrate for various engineering strain rates. (a) Tensile stress versus strain curves. (b) Compressive stress versus strain. (c) Tensile deformation test conditions. (d) Compressive deformation test conditions [89]. (The color of molecule particles is based on the potential energy. The purple molecules indicate the uniformly structural changes during deformation)

By evaluating the mechanics of CH₄ hydrate at the liquid water-hydrate interface, it was observed that in the region close to the interface, CH₄ molecules due to dense hydrogen-bonded water molecules cannot be soluble. However, the growth of local density and surface excess of CH₄ near the surface indicate anomalous CH₄ adsorption to the CH₄-water interface [366]. Previously, the free energies and thermal expansivity of clathrates showed that the large expansivity of xenon hydrate stems from guest molecules and a difference in oxygen atom arrangement between ice and hydrate plays a minor role [367]. By employing rigid geometry and adopting the non-polarizable model in NEMD simulations it was shown that except at low temperatures (50 K), although the estimated thermal conductivity values are relatively insensitive to the cage occupancy, a small number of defects (1%) in the water lattice can result in a large reduction (10%) in thermal conductivity [86]. The effect of electrostatics methods on the prediction of CH₄ hydrate thermal conductivity showed that the estimated thermal conductivities by non-periodic techniques would be more in agreement with the experiments, however, electrostatic treatments affect thermal conductivity [368]. Thermal conductivity may also depend on both rigidities of the framework and guest-host interactions but the lower thermal conductivity compared to ice *Ih* is because of differences in crystal structures [369]. Studying the mechanisms for thermal conduction of CH₄ hydrate revealed that the guest-host interactions and crystal structure can contribute to glasslike temperature dependence so that they can lower the hydrate thermal conductivity relative to ice *Ih* [370]. Based on MD analysis of the thermal conductivity of CH₄ hydrate in porous media, it was determined that with increasing temperature, CH₄ hydrate-SiO₂ thermal conductivity enhances more quickly near the freezing point. Also, at a certain temperature, smaller pore sizes have a larger effect on increasing hydrate thermal

conductivity which is several times less than the porous media [227]. Mechanisms for thermal conduction in H₂ hydrate demonstrated that dissimilar to CH₄, a single occupation of the small cavity by H₂ results in more harmonic energy transfer and less resonant scattering [371]. The thermal properties of sI CH₄ + LMGs hydrates also showed that at high-temperature ranges, guest molecules with strong host-guest interactions have more thermal expansion but increasing the size of LMGs subsequently decreases the C_p, and C_v of hydrate [372]. A correction technique to determine isothermal expansion coefficients and lattice parameters at 287 K and 10-100 MPa highlighted that the isothermal expansion coefficient and specific heat capacity of C₃H₈ hydrate are somewhat higher than that in THF hydrate. Furthermore, by comparing C₃H₈ + CO₂ and C₃H₈ + CH₄ binary hydrates, it was revealed that the small cages occupied by CO₂ molecules dissimilar to CH₄ can result in increasing compressibility and expansion coefficient but decreasing the heat capacity [373]. According to structure-mechanical properties of CH₄/ Xe/ N₂/ H₂/ Ar + neohexane sH hydrates using *ab initio* atomic simulations, the type of help gas was identified as the main contributor to the shear, elastic and anisotropic characteristics of sH hydrate. Also, the brittleness of filled sH hydrate was found to be higher than empty sH hydrate which can be attributed to the role of guests inside cavities. However, wave velocities of ice *Ih* were determined less than sH hydrates [156]. Based on the simulated values of the lattice constants using anisotropic site-site potential models of the N₂+ neohexane hydrate, the same tendency as obtained by the powder X-ray diffraction was observed [374]. Moreover, studying the impact of help gas on the crystal lattice of Ar/ Kr/ CH₄+ neohexane (NH) at different temperatures determined that the order of their crystal lattice as CH₄ > Kr > Ar may be the evidence of the effects and importance of small guest molecular sizes [152]. It was also indicated that the lattice constants and system average volume for xenon hydrate in the presence of cyclopropane and propane are directly related to the guest molecular sizes and pressure [375]. Recently, the lattice expansion of CH₄ sII hydrate in the existence of linear and cyclic guest molecules (cyclobutanol and

butyraldehyde) with the identical formula of C_4H_8O was studied to figure out the effect of bonding characteristics on hydrate thermal expansion. Based on the results, less lattice expansion in the system consisting of cyclobutanol in comparison with butyraldehyde was observed but the promotion impact of cyclobutanol on hydrate equilibrium was found to be higher [376].

Simulation of the thin liquid film at the hydrate/ CH_4 gas interface showed that the thin liquid film makes substantial contributions to the interfacial properties [69]. The interfacial free energy, excess enthalpy, and stress at the crystal-liquid interface by employing a normal pressure-cross-sectional area (NP_{NAT}) ensemble elucidated that interfacial tension is directly and reversely proportional to temperature and pressure respectively [46]. Based on molecular analysis of the interfacial mechanics and thermodynamics at the interface of liquid water and natural gas, it was found that for the pressure range of less than 50 MPa, the increase of pressure and temperature reduces the interfacial tension [377]. Interfacial tension and CH_4 hydrate morphologies at the liquid water-hydrate interface showed that the nucleation may take place in the ranked order of film-shaped, cap-shaped, lens-shaped, and homogeneous [378]. By examination of the interfacial tension and behaviour of the single and mixture of CH_4 , C_2H_6 , and C_3H_8 hydrates, equilibrium molecular dynamics and hydrate formation at the water-gas interface were observed. At the interface, the surface adsorption of gas molecules enhances the gas local concentration but reduces the interfacial tension, however, both of these factors boost the formation of gas clathrate hydrate [222]. To determine the mechanical and key structural properties such as bulk modulus and equilibrium lattice volume for C_3H_8 , $C_2H_6 + C_3H_8$, $CH_4 + C_3H_8$, and empty hydrates, it was determined that the equilibrium lattice volumes are directly proportional to the guest size in double hydrates which may experience the greater increase than pure hydrates [379]. Based on MD analysis of H_2 and D_2 molecular escape from the hydrate phase using the same force-field parameters at 150 to 195 K, it was also revealed that the stronger collisions between the cage and guest molecules may increase the

likelihood of cage distortion which results in the subsequent escape of D₂ molecules. Also, the leakage rates of H₂ were observed to be lower than those for D₂ [218]. In addition, the in-slab translational order parameters of the CH₄ hydrate surface in contact with the gas phase determined that the molecules of water in the adlayers and slab surfaces can be periodically arranged while the internal slabs can remain completely rigid [97].

3.6 Dynamical and vibrational behaviour

According to MD results, it was confirmed that the power spectra would be a reliable technique to evaluate the vibrational behaviour of guests in the hydrate phase [380]. Previously, the idea of a resonant scattering mechanism for the host-guest phonon interactions in clathrate hydrates was supported by MD simulations [381]. The dynamical behaviour of ice and sI clathrate hydrate showed that the phonon densities in both states are similar which represents the resemblance of their infrared spectra. Also, the phonon density enhancement of selective vibrational modes can be induced by guest species [34]. Phonon scattering off the rotational and vibrational motions of guests was also found to be responsible for the lower thermal conductivity of hydrates relative to ice [382]. According to the exploration of lattice vibrations in clathrate hydrates for Xe at 220 K, vibrational couplings lead to energy exchange between the guest vibrations and the host lattice which results in anomalous glasslike behaviour in the thermal conductivity [383]. Examining anomalous thermal conductivity of clathrate hydrates using EMD, and NEMD simulations revealed that the resonance scattering model is the most likely mechanism of anomalous temperature dependence and low thermal conductivity [384]. Evaluating energetic dynamical and H-bond vibrational properties in sI and sII hydrates demonstrated that H-bond energies are the largest in magnitude in structure I, followed by sII hydrate and then ice *Ih*. However, unlike ice and sII hydrate, energy transfer via H-bonds in sI occurs at higher frequencies [74]. MD calculations for the vibrational spectral band of molecular H₂ trapped in hydrate indicated that H₂ molecules in large cavities provide a

high-frequency peak than those in small cages [385]. With the combination of *ab initio* and classical MD simulations to analyse the host-guest H-bonding in alcohol clathrate hydrates it was found that due to the strong host-guest H-bonding, vibrational spectra of alcohol O–H bonds consistent with experimental Raman spectra display large frequency shifts. Also, conducting similar investigations of dynamical properties of H₂S revealed that a preferred orientation of the dipole-vector exists at 150 K whereas there is no preferred orientation at 300 K [386]. Interestingly, some alcohol molecules can fit into 5¹²6⁴ cages of clathrate hydrates [387]. The mechanical and vibrational features of CH₄/ Xe/ CO₂+ neohexane sH hydrates showed that vibrational frequencies are dependent on interatomic distances of hydrate and pressure. Also, the relations between interatomic distances, H-bond strength, and vibrational frequency shifts, affected by guest-host interactions and guest type were previously certified [388]. Based on the stretching vibrations of sI CH₄ hydrate investigated with *ab initio* density functional theory, although the consistency between calculated vibrational frequencies and experiments was confirmed, the computed H-C stretching vibrations are less than in the free molecules [350]. In addition, *ab initio* results to investigate the vibrational spectra of sH CH₄ hydrate indicated that asymmetric and symmetric stretching vibrational frequencies of CH₄ molecules are higher in small cages of sH hydrate than in medium cavities, however, the H-C bond length in medium cavities is slightly longer [389]. Also, asymmetric and symmetric stretching vibrational frequencies of CH₄ molecules are higher in small cavities than in large ones of sI hydrate [390]. Based on the infrared spectra of sII C₃H₈, i-butane, CH₄ + C₂H₆, CH₄ + C₃H₈, and empty hydrates using DFT, it was determined that the calculated vibrational frequencies of the guest and host molecules can give comparable results to experimental data which can be useful to detect the presence of gas hydrates [391]. Analysing the properties of multiple occupied N₂ clathrate hydrates showed that there is a large vibrational host-guest coupling for both double and single occupancies, however, the frequency range is broader in the former case [392]. To evaluate the mobility of water and

CO₂ molecules in the hydrate phase, it was elucidated that negligible composition change leads to significant impacts on the mobility of CO₂ molecules. For example, altering 3% cage occupancy may result in two orders of magnitude change in the diffusion coefficient [393]. Molecular H₂ mobility within the sII clathrate at 200-265 K demonstrated that on the time scale of the simulations, H₂ molecular mobility in the small cavities is not substantial but can be more clathrate stable with the presence of THF [394]. By analysing apparent high mobility and transport of interstitial H₂O defects in CH₄ hydrate crystal it was concluded that H₂O molecules are somewhat high mobile entities within a gas hydrate. Also, the presence of empty cavities facilitates the pathways for the H₂O molecular transportation between them [395]. In addition, the hydration shell nucleation of guest molecules becomes more ordered with increasing the concentration of guest molecules, resulting in a decline of entropy and guest mobility [396]. The results of first-principles DFT and MD to determine the thermal properties of gas hydrates at very low temperatures confirmed the negative thermal expansion which is similar to the ice [397]. The structural and dynamical properties of CO₂, CH₄, and Xe hydrates indicated that by elevating the temperature, the lattice expansion of CO₂ hydrate is larger than that in Xe and CH₄ hydrates. Also, the thermal conductivities of both Xe and CO₂ are less than that in CH₄ or even empty lattice [398]. With respect to the analysis of structural and energetic properties of the sI and sII CO hydrate it was revealed that increasing the content of CO molecules in the large cages can stabilize the sII but destabilize sI hydrate [399]. Using the rotational autocorrelation function (RACF) to study the host and guest rotational dynamics, it was highlighted that altering the rotational motion of both water and guest molecules affect the proportion of them in mixed CO₂+CH₄ hydrate [400]. Dynamical and energetic properties of H₂+THF hydrate through EMD simulations determined that the van der Waals component with the surrounding water molecules in the constituent cavities is the largest contribution to the interaction energy of both guests [401]. Using first-principles DFT to study the vibrational, structural, and mechanical properties of

THF and THF+Xe hydrates, it was found that compared to THF hydrate, THF+Xe hydrate has a larger OH stretching frequency, lattice volume, hydrogen bond length but lower Poisson ratio, density, compressional wave velocity, and hydrogen bond stretching frequency [402].

3.7 Coexistence of phases

The phase equilibrium of clathrate hydrates more specifically those at thermodynamically difficult to conduct the experiments in a laboratory can be studied using MD simulations. For example, the thermodynamic stability of sII neon hydrate at 480 MPa and 260 K with employing MD simulations was confirmed [403]. Previously, by employing MDs, a new phase between equilibrium conditions of sT' and C₀-II structures for H₂ hydrate was explored which helped to clarify the experimental puzzle of H₂ hydrates [116]. In addition, at ambient temperature and 3-130 GPa, the phase diagram of the C₂ structure of H₂ hydrate was determined [216]. However, the results of MD to determine the three-phase coexistence of H₂ hydrate at 90-400 MPa were found to underestimate the experiments by approximately 25 K [404]. The MD predictions of the three-phase coexistence of binary CH₄+CO₂ hydrates at 40, 100, and 200 MPa compared to experimental data indicated that the deviation of simulations up to 100 MPa is not significant [405]. MD insights into the stability of CO hydrate at 17.3 MPa and 243 K revealed that CO tends to form sI rather than sII clathrate, however, with increasing pressure and temperature to 10 MPa and 252 K, it would be more prone to generate sII clathrate hydrate [406]. In addition, *ab initio* intermolecular potentials at the MP2 level to calculate energy surfaces of CH₄-H₂O and CO₂-H₂O showed that the angle-dependent approach improves the prediction of the CO₂ and CH₄ hydrates phase equilibrium [407]. MD simulation of three-phase equilibrium suggested that compared to NPT simulation, NVT has two benefits: first, there is no need to control the pressure. Second, NVT reduces the number of time steps in simulations. Therefore, it is more suitable to study the phase coexistence of gas hydrates [75]. Moreover,

the equilibrium condition can be reached more accurately when the temperature and total energy become constant [408].

3.8 Gas exchange phenomenon

Evidence suggests the gas exchange occurs via a transient co-occupation of CO₂ and CH₄ in cages [409,410]. Both MD simulations as well as Raman spectroscopic confirmed that the process of CO₂/CH₄ replacement is the breakup of the cage, the escape of CH₄, and cage occupation by CO₂ molecules [411]. It was also elucidated that during gas exchange phenomena, replacement takes place at both small and large cages of CH₄ hydrate without changing the structure and with the partial collapse in which the hydrate surface is partially melted so that the interface becomes active [120]. MD evidence indicated that the increase of temperature from 250 to 270 K accelerates the kinetics of the CO₂/CH₄ replacement by at least 1.5 times, also, a swap of the guest molecules without a breakup of cages was confirmed [410]. Previously, for CO₂ capture and storage analysis, Gibbs free energy calculations for CO₂ clathrate hydrates in the presence of H₂S, CH₄, N₂, and SO₂ were performed. Based on the results, the negative values of ΔG for SO₂ and H₂S impurities compared to CO₂ showed more stability in the hydrate phase. Although at lower concentrations, these impurities act as promoters, large amounts of them decline the CO₂ capture and storage capacity [412]. According to the CH₄+CO₂ hydrate formation, the concentration of CO₂ plays a key role in the kinetics of CH₄+CO₂ hydrate formation. However, increasing CO₂ concentration in the aqueous phase cannot give faster growth [413]. It was also revealed that during the gas replacement process, CO₂ molecules in mixed bubbles mostly surround the CH₄ molecules so that they influence the process of gas exchange specifically at the initial stage [414]. MD investigations of the replacement or co-growth of CO₂ and CH₄ hydrates indicated that CH₄ hydrate in the presence of CO₂ gas is more stable than with CO₂ solution. It was also estimated that nearly 20% of the dissociated CH₄ hydrates can be replaced by CO₂ and most likely CO₂+CH₄ mixed hydrates can be formed [220].

Additionally, the guest behaviour in a porous environment would be dissimilar. MD exploration of the transport properties of CH₄ and CO₂ hydrates in Na-montmorillonite clay determined that increasing CH₄ and CO₂ molecules in Na-montmorillonite interlayers probably result in a decrease in their self-diffusion coefficients [200]. Also, the formation and dissociation of CH₄ hydrate in the clay pore with fatty acids showed that the existence of fatty acids slightly accelerates the breakdown of CH₄ hydrate in the heterogeneous sediment [228]. Moreover, the amorphous layer formation was detected as a barrier against mass transfer which results in a slower rate of CO₂/CH₄ replacement because as time proceeds, the CO₂ amorphous layer forms on the CH₄ hydrate surfaces [119].

Another barrier is that since CH₄ relative to CO₂ possesses a smaller size, it can be more stable in small cages [415]. Also, N₂ guests can be used as a carrier gas because it does not compete directly with CO₂ during CH₄ substitution. In addition, the substitution of CH₄ in the small cavities with N₂ has positive free energy [416]. The effects of N₂ on the process of CO₂/CH₄ replacement showed that the addition of N₂ aids CO₂ penetrate into all CH₄ hydrate cages on a broader scale. However, this diffusion is sensitive to the ratio of CO₂ to N₂ [221]. The CH₄ replacement by flue gas in the hydrate phase in the presence of SO₂, H₂S, N₂O/ NO, and CS₂ revealed that N₂O, SO₂, and CS₂ molecules tend to occupy the large cavities of sII and sI hydrates while NO, and H₂S have no preference to occupy small or large cavities. It was also confirmed that CS₂, N₂O, H₂S, and SO₂ can replace CH₄ in the hydrate phase and help the process of gas exchange [417].

3.9 Memory effect phenomenon

Another feature of gas hydrate is the observation of a memory effect. In this phenomenon, the resulting solution from the decomposed hydrate is able to form a hydrate more readily with a shorter induction time than a fresh solution. Since a number of studies to test a memory effect have failed to observe this

phenomenon, it may be concluded that the memory effect does occur for just specific hydrate systems. Generally, the most popular model for the memory effect is the residual structure hypothesis which assumes the dissociation of gas hydrates leads to the formation of some “hydrate melt”. Therefore, residual structures that retain some structural features of the hydrate phase would persist in the liquid water phase for a long time after the decomposition. For example, guest molecules with associated pentagonal rings of hydrogen-bonded water molecules provide nucleation sites for the second formation of gas hydrate when they cooled again [7]. MD simulations of CH₄ hydrate pointed out that the lifetime of a hydrogen bond in water is of the order of picoseconds [418]. It was also assumed that the freezing of a bulk hydrate-forming solution occurs similarly to the freezing of bulk water in which the nucleation is mostly heterogeneous and rarely homogeneous in bulk solution [419]. MD investigations of the properties of melting temperature considering the memory effect suggested that the coordination atoms of oxygen in water are an important factor in the memory effect but water molecules near the interface of water-CH₄ have fewer memory properties [40]. Moreover, MD modelling of two-step nucleation and memory effect in CH₄ hydrate clarified that the areas locally richer by CH₄ molecules enhance the apparent nucleation rate more readily. Also, the memory of the crystal for fast re-crystallization was found to be insignificant [4].

3.10 Self-preservation phenomenon

Self-preservation would be an advantageous property for the transportation and storage of gas hydrates. Quasi-harmonic lattice dynamics (QLD) model considering guest-guest interactions of multiple H₂ occupancies to predict self-preservation and the thermodynamic properties of sII H₂ hydrate specified that the pressure in the H₂ hydrate is more than that in the ice phase but the hydrogen bonds between the ice and hydrate do not allow hydrate to be destroyed, so that, the hydrate phase under heating remains stable [420]. The self-preservation mechanism using the combination of the NVT and NVE method for

CH₄ hydrate also proved that the coupling of heat and mass transfer resistances is the driving mechanism for self-preservation impact [55]. According to MD simulations, the water self-diffusion coefficient can be changed by altering the temperature and guest concentration [57]. Based on the calculated order parameter, it was found that the order parameter value for ice, hydrate and liquid water phases by partially heating is different but it can be changed at the interfaces of the phases [54]. Through MD exploration of the self-preservation of hydrate decomposition using NVT/E to represent different levels of heat transfer resistance, it was observed that heat transfer resistance facilitates the formation of the solid-like layer which inhibits further hydrate dissociation. Also, the increase in pressure and particularly the decrease in temperature enhances self-preservation [55]. Moreover using MD algorithms, the number and type of cavities in the amorphous and clathrate (sI, sII, HS-I) coexist with ice during CH₄ hydrate self-preservation can be monitored [421]. Recently, the sources of the THF hydrate anomalous preservation surrounded by CP hydrate were evaluated outside the stability conditions. The melting temperature of uncoated THF hydrate was determined to be 270 K whereas THF hydrate encapsulated with CP hydrate could not be dissociated up to 290 K. This phenomenon also indicates the transformation of the THF hydrate from heterogeneous to homogeneous mechanism. As Figure 16 (a) shows, the potential energy of the coated THF hydrate for the 50 ns at 280 K can remain stable. Figure 16 (b) indicates the partial breakup of coating layers of THF hydrate at 290 K which results in the liquid layer formation on top of the hydrate phase. With increasing temperature to 300 K, the potential energy suddenly elevates due to the dissociation of all CP hydrate followed by the breakdown of THF hydrate as the inner layer [422].

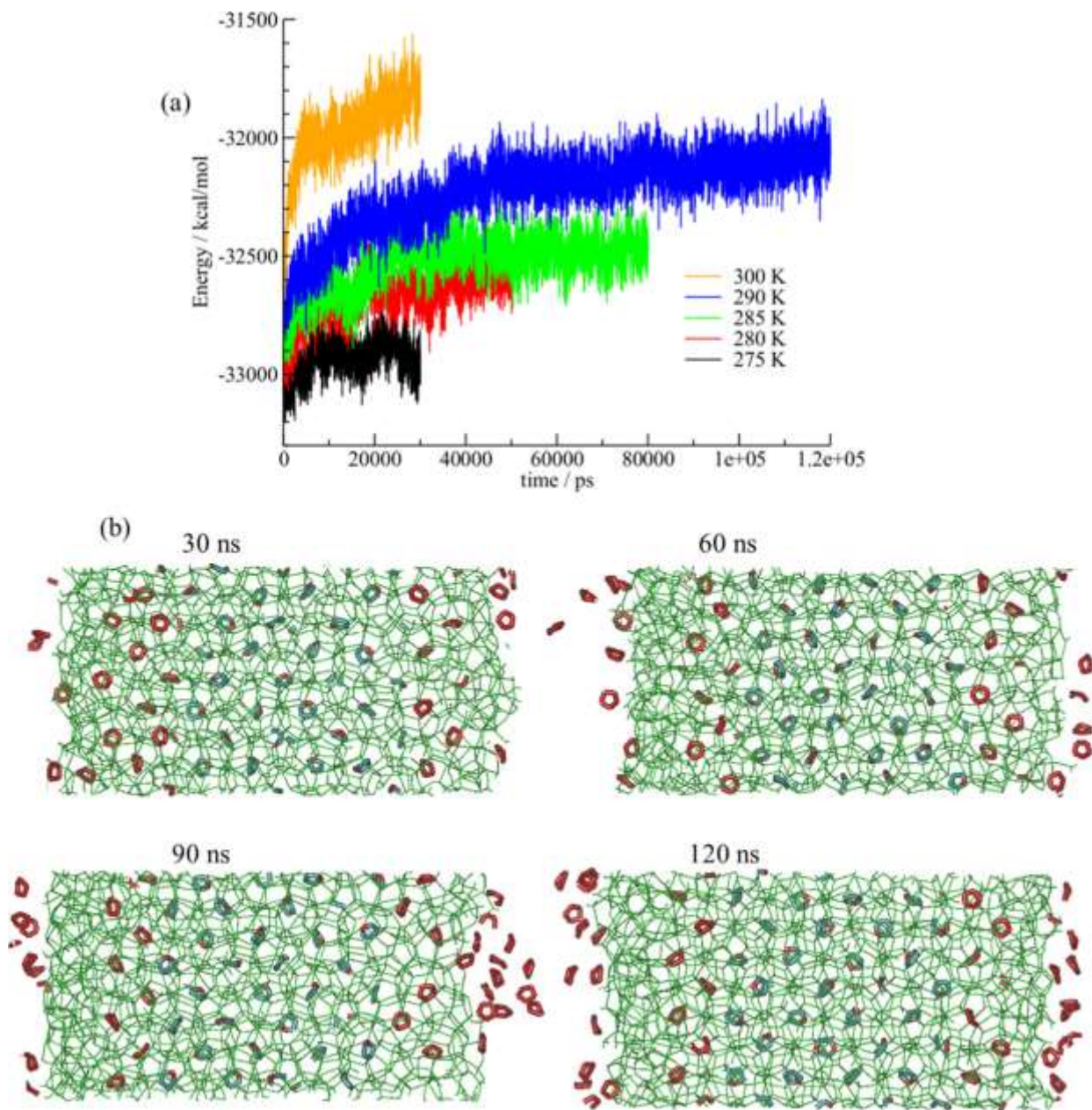


Figure 16: (a), Time evolution of the potential energy of THF hydrate coated by CP hydrate at different temperatures; (b) Snapshots of THF-CP hydrate decomposition at 290 K [422]

In the self-preservation phenomenon, the effects of H-bonds are also critical. In this regard, the influence of host-guest H-bond on the properties of H₂S/ Xe + pinacolone/ MTBE sH hydrates showed that MD simulation results are consistent with the observations of X-ray crystallographic at the temperature range

from 100 K to 250 K [423]. Based on MD simulations of the halogen bonding in BrCl, mixed Cl₂+Br₂, Cl₂, and Br₂ clathrate hydrates, the short and strong interactions between water molecules and bromine atoms in the sI clathrate hydrates were observed [424]. In addition, halogen bonding in Cl₂ and Br₂ clathrate hydrates studied via both MD simulations and DFT showed that the obtained halogen-water distances were compatible with values observed in X-ray diffraction [425]. The anomalous halogen bonding interactions between Cl₂ and Br₂ with water indicated that a combination of the dihalogen interaction with oxygen lone electron pairs may result in the halogen bonding non-bonded interactions between water in the clathrate and guest molecules [426]. To explore the hydrogen bonding in binary sI hydrate at 183-263 K, it was found that although gas molecules in the small cavity do not themselves form hydrogen bonds with water, the THF molecules affect the occurrence of hydrogen bonds. Also, nearest neighbour guest-guest interactions were found to influence the stability and structure of the clathrate network [427]. Hydrogen bonding study of pure and binary CO₂+THF hydrates indicated that a small percentage of hydrogen bond formation between water and THF occurs, however, it cannot be seen through single-crystal X-ray diffraction at low temperatures. Interestingly, the presence of hydrogen bonding guests can increase CO₂ migration and boosts the hydrate formation kinetics [428]. MD simulations of the role of the medium and small cage guests on the hydrogen bonding of the large cage guests with the hydrate framework water molecules exhibited that the presence of CH₃F enhances the H-bonding probability of the TBME with the water molecules [429]. Exploring the hydrogen bonding in binary sI hydrate at 100-250 K revealed that ethanol molecules by forming a long-lived proton- accepting and donating hydrogen bonds with water molecules, support the general cage integrity of clathrate hydrate [430]. The hydroxyl groups of inhibitors e.g. ethanol/ 1-propanol/ 2-propanol may act as both proton donors and proton acceptors. Also, the probability of H-bond between hydroxyl atoms with water molecules was found to be considerable [226]. By conducting the MD exploration of the microscopic

properties of the HFC-41/ CH₄+ pinacolone hydrate, although the anisotropic expansion of the sH hydrate lattice for sH HFC-41+pinacolone hydrate was not observed, weak H-bonding of the water and pinacolone molecules was detected [224]. Recently, hydrogen bonding analysis for CH₄/ HFC-32 +N-methylpiperidine (NMP) hydrate determined that for HFC-32+NMP sH hydrate which was more stable than CH₄+NMP sH hydrate, NMP molecules with water form H-bonds whereas similar hydrogen bonding for CH₄+NMP system was not observed [431].

4 Proportions of MD gas hydrate investigations and future research guidelines

Figure 17 exhibits the proportion of different MD studies that have been conducted on gas hydrates in the literature. As is shown, semiclathrate hydrates possess the least share of these investigations whereas over half of MD studies have been performed on pure gas hydrates. It should be kept in mind that to perceive fundamental microscopic characteristics of gas hydrates, most MD studies have been performed on CH₄ hydrate, therefore that is why its percentage is near one-third of the total explorations. In addition, the proportions of insight into the specifications of other pure, binary, and mixed gas hydrates, the effects of minerals, thermodynamic/ kinetic promoters, and inhibitors have been reasonably close together.

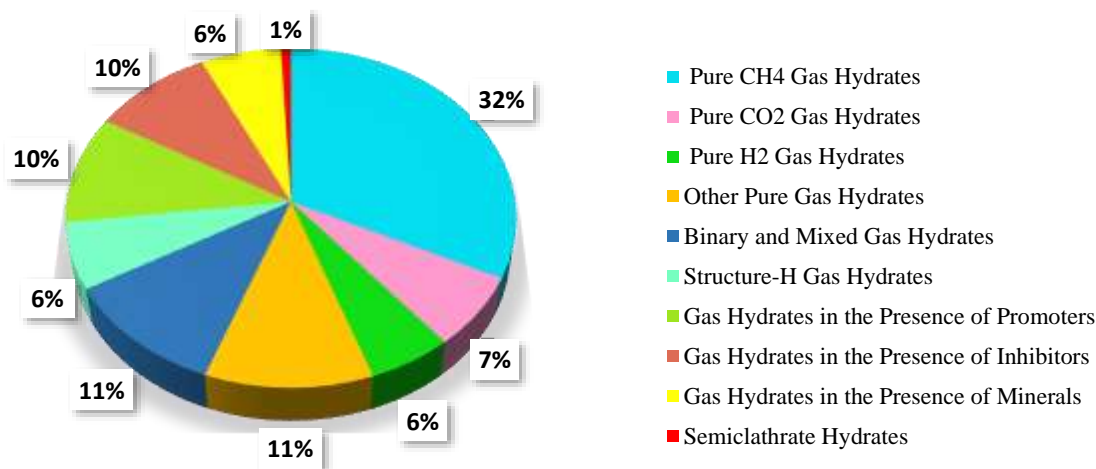


Figure 17: The proportion of MD simulation studies on gas hydrates.

In this work, different phenomena and properties of gas hydrates were overviewed. The summary list of MD computational studies conducted on clathrate (sI/II/H) hydrates ranging from pure/binary/mixed components, or in the presence of promoters/ inhibitors/ minerals is presented in [Table 2](#).

Table 2: Investigated phenomena for various hydrate systems using MD simulations.

Component	Dynamical and Nucleation/ aggregation	Growth	Dissociation	Thermo-physical <small>mechanistic</small>	Mechanical properties	Thermodynamic / Kinetic properties	Morphology/ structural analysis (stabilization)	Heat and mass transfer properties	Memory effect	Self-preservation	Coexistence of phases/ Phase equilibrium	Cage occupancy/ Storage capacity	Ref.
Pure hydrates	CH ₄	*	*	*	*	*	*	*	*	*	*	*	[4,23,38–46,49,26,50–59,27,60–64,66–70,32,71–80,33,82,84,86,88–94,34,95–100,108,113,117,213,35,214,215,234,235,237,244,245,253–255,36,256,257,261,290,291,297,298,304,318,355,37,366,368–370,378,380,421,432]
	H ₂	*	*	*	*	*	*	*	*	*	*	*	[40,115,420,116,216,218,294,307,331,385,404]
	CO ₂	*	*	*	*	*	*	*	*	*	*	*	[25,212,221,247,328,361,428]
	C ₂ H ₄ , C ₂ H ₆ , C ₃ H ₈	*	*	*	*	*	*	*	*	*	*	*	[21,111,219,262,295,302,303,352]
	Acid gas (H ₂ S, CO, CS, CS ₂ , SO ₂ , N ₂ O)	*	*	*	*	*	*	*	*	*	*	*	[210,240,353,361–363,386,399,406]
	Noble gas (Ar, Kr, Xe, Ne, N ₂)	*	*	*	*	*	*	*	*	*	*	*	[163,299,398,403,339,340,367,381,383,384,392,396]
	Liquid hydrocarbon (EO, CP, THF, THP, C ₃ H ₉ NO, C ₂ H ₄ O, CH ₂ O, C ₃ H ₆ O ₂ , C ₃ H ₈ O)	*	*	*	*	*	*	*	*	*	*	*	[106,110,232,233,248,250,263,365,382,397]
Gas hydrates + KHP/THP	CH ₄ / H ₂ / Xe + THP (C ₃ H ₈ / THF/ TBME/ C ₂ H ₆ / TMO/ EO/ FA/ CB/ CP/ ISA/ TBAB/ SF ₆ / Asphaltene)	*	*	*	*	*	*	*	*	*	*	*	[109,131,269–271,300,308,330,333,334,336,337,132,347,372,375,387,394,401,402,141,142,168,208,252,267,268]
	CH ₄ / H ₂ / N ₂ + KHP (SDS/ CAPB/ SWCNT/ Protein/ lecithin)	*	*	*	*	*	*	*	*	*	*	*	[140,185,266,301]
Gas hydrates + KHI/THI	CH ₄ + C ₂ H ₆ / C ₃ H ₈ / C ₄ H ₁₀ + THI (Salt: NaCl/ NH ₄ Cl / KCl/ CaCl ₂) / (Alcohol: (methanol/ Ethanol/ 1-Propanol/ 2-Propanol/ Glycerol)	*	*	*	*	*	*	*	*	*	*	*	[166,169,277,281,311–315,327,170–172,182,225,226,275,276]
	CH ₄ + KHI (Na-MMT/LHA/ Amino acids/ PVPs/ PVCap/ PDMAEMA/ PEO / VIMA/ AFPs)	*	*	*	*	*	*	*	*	*	*	*	[167,173,183,184,279,280,174–181]

Mixed gas hydrates	CH ₄ + C ₂ H ₆ / C ₃ H ₈ / Acid gas (H ₂ S/ CO NO/ N ₂ O/ SO ₂ / CS ₂)	*	*	*	*	*	*	*	*	*	*	[102,124,417,125,222,242,243,373,377,408,412]
	H ₂ + CH ₄ / C ₂ H ₆ / C ₃ H ₈			*	*	*				*		[251,265,296,338]
	Cl ₂ + BrCl/ Br ₂ / NH ₃	*				*		*				[309,422,424–426,429]
sH hydrates	H ₂ / CH ₄ / N ₂ / Ar/ Kr/ Xe/ H ₂ S + LMGs	*	*	*	*	*		*	*	*	*	[150,152,346,374,423,156,160,161,223,224,310,341,345]
Hydrate + minerals	CH ₄ + NMP/ SiO ₂ / Clay/ Graphite / Kaolinite	*	*	*	*	*	*	*	*	*	*	[29,196,292,293,356,358,364,431,197,203,227,228,283–286]

The performed MD frameworks to analyse the phenomena and characteristics associated with hydrate-based applications either for gas separation or utilization can complete or at least support the experimental measurements. Despite a number of MD simulations of gas hydrates reported in the literature, conducting MD simulations to investigate unexplored gas hydrate characteristics can also complete the present findings at the macroscopic level. Hence, we aim to highlight some suggestions for future research below:

- Additional MD investigations on the combined effects of various inhibitors and promoters such as KHI+THI and KHP+THP on clathrate hydrates to identify the relationship between hydrate phenomena and the presence of additives can be performed. These simulations may help to comprehend the characteristics connected to the performance parameters of hydrate-based gas separation, CO₂ hydrate utilization, or other hydrate-relevant applications.
- More MD simulations to understand some gas hydrate phenomena such as memory effect and self-preservation at different environmental or operating conditions need to be carried out.
- In spite of several experimental suggestions to utilize the semiclathrate hydrates in different processes of hydrate-based applications e.g. secondary refrigeration and air conditioning aims, the least proportion of MD simulations is for this type of gas hydrates. Therefore, a manifold of

MD explorations to reveal the molecular mechanisms of semiclathrate hydrate promoters can be conducted.

- Only a few studies focused on clathrate hydrate structural transitions from sI to sII or sH and the coexistence of the phases based on the guest molecular sizes, shapes, and concentrations at the molecular scale. Hence, simulations considering different thermophysical properties during the process of structural transitions need to be performed.
- Heterogeneous nucleation MD simulations with fast mass transfer in the hydrate formation, and analyzing the role of nanobubbles during the nucleation phase can be considered for further research.
- To date, although some MD simulations in porous media have been performed, studying the systems simultaneously including mixed minerals e.g. kaolinite, quartz, montmorillonite, and kaolinite can help to assess the effects of these components on gas hydrate phenomena.
- In order to evaluate the influences of permeability and wettability on hydrate-bearing sediments, more MD studies should be carried out.
- Discrepancies between MD simulations and experimental outcomes in terms of consistency with the real condition can be assessed by providing simulations of the phenomena at larger scales. Therefore, a better understanding of the mutual relationship among mechanisms can be achieved.

Nomenclature

HBGS	Hydrate-Based Gas Separation	NPE	Nonyl Phenol Ethoxylates
HBCC	Hydrate-Based Carbon Capture	LAE	Lauryl Alcohol Ethoxylates
NGH	Natural Gas Hydrate	SDS	Sodium Dodecyl Sulfate
InSAR	Interferometric Synthetic Aperture Radar	SL	Sulfonated Lignin
MRI	Magnetic Resonance Imaging	SHS	Sodium Hexadecyl Sulfate

¹³ CNMR	Carbon-13 Nuclear Magnetic Resonance	STS	Sodium Tetradecyl Sulfate
COC	Cyclic Organic Compounds	SDBS	Sodium Dodecyl Benzene Sulfonate
CN	Coordination Number	DMSO	Di-Methyl Sulf-Oxide
GHSZ	Gas Hydrate Stability Zone	TMS	Tetra-Methylene Sulfone
CSMHYD	Colorado School of Mines Hydrate	SWNT	Single-Walled carbon Nano-Tube
T/KHI	Thermodynamic/Kinetic Hydrate Inhibitor	MWCNT	Multi-Walled Carbon Nano-Tube
T/KHP	Thermodynamic/Kinetic Hydrate Promoter	Na-MMT	Sodium Mont-Morillonite
LMGs	Large Molecule Guests	SW-CNTs	Single-Walled Carbon Nano-Tubes
AIMD	<i>Ab Initio</i> Molecular Dynamics	SAMs	Self-Assembled Monolayers
TMO	Tri-Methylene-Oxide	TMS	Tetra-Methylene Sulfone
FA	FormAldehyde	EO	Ethylene Oxide
CB	CycloButane	LHA	Leonardite Humic Acid
MWCNT	Multi-Walled Carbon Nano-Tube	ACF	Auto-Correlation of the Fluctuations
THT	Tetrahydrothiophene	AOP	Angular Order Parameter
THF	TetraHydroFuran	APDF	Angular probability distribution function
PMF	Potential of Mean Force	CGMC	Grand Canonical Monte Carlo
OACF	Orientation Auto-Correlation Function	DFT	Density Functional Theory
QLD	Quasi-harmonic Lattice Dynamics	DWC	Dodecahedral Water Cluster
TCF	Time Correlation Function	FSICA	Face-Saturated Incomplete Cage Analysis
VACF	Velocity Auto-Correlation Function	HCACF	Flux AutoCorrelation Function
RDF	Radial Displacement Function	MCG-OP	Mutually Coordinated Guest Order Parameter
RACF	Rotational Auto-Correlation Function	MCG	Mutually Coordinated Guest
RPMD	Ring Polymer Molecular Dynamics	MSD	Mean Square Displacement
RMSF	Root Mean Square Fluctuation	NEMD	Non-Equilibrium Molecular Dynamics
RPMD	Ring Polymer Molecular Dynamics	MFPT	Mean First-Passage Time

Conflict of interest

The authors declare no conflicts of interest.

Acknowledgment

The first author would like to thank Curtin University Malaysia for providing Curtin Malaysia Postgraduate Scholarship (CMPRS) and necessary resources and financial support for this work. Also, we would like to thank Dr. Saman Alavi for their valuable suggestions and comments on this work.

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